

Functional inequalities and the symmetry properties of the extremal functions

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IN COLLABORATION WITH

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Attainability and value of best constants in functional inequalities

Suppose that we know the existence of a constant C such that in some space X , for all $u \in X$,

$$F(Du, u, x) \leq C G(D^2u, Du, u, x) .$$

Functional inequalities play an important role in obtaining **a priori estimates** for solutions of PDEs, in analyzing the **long time behavior** of solutions of evolution problems, in describing the **blow-up profile** for finite time blow-up phenomena, etc

Knowing the exact value of C or the properties of the optimal solutions (when they exist) can be crucial in some critical cases.

Important questions :

- Is C attained in X ?
- If yes, how do the optimal functions u look like?

If we know **a priori** that the optimal solutions have some symmetry properties, for instance, that they are radially symmetric, then it might be easier to compute the value of C .

Caffarelli-Kohn-Nirenberg (CKN) inequalities

$$\left(\int_{\mathbb{R}^d} \frac{|u|^p}{|x|^{\frac{b}{p}}} dx \right)^{2/p} \leq C_{a,b} \int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{\frac{2}{a}}} dx \quad \forall u \in \mathcal{D}_{a,b}$$

with $a \leq b \leq a+1$ if $d \geq 3$, $a < b \leq a+1$ if $d = 2$, and $a \neq \frac{d-2}{2}$

$$p = \frac{2d}{d-2+2(b-a)}$$

$$\mathcal{D}_{a,b} := \left\{ |x|^{-b} u \in L^p(\mathbb{R}^d, dx) : |x|^{-a} |\nabla u| \in L^2(\mathbb{R}^d, dx) \right\}$$

$$b-a \rightarrow 0 \iff p \rightarrow \frac{2d}{d-2}$$

$$b-(a+1) \rightarrow 0 \iff p \rightarrow 2_+$$

$$\frac{1}{C_{a,b}} = \inf_{\mathcal{D}_{a,b}} \frac{\int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{\frac{2}{a}}} dx}{\left(\int_{\mathbb{R}^d} \frac{|u|^p}{|x|^{\frac{b}{p}}} dx \right)^{2/p}}$$

The symmetry issue

$$\left(\int_{\mathbb{R}^d} \frac{|u|^p}{|x|^{bp}} dx \right)^{2/p} \leq C_{a,b} \int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{2a}} dx \quad \forall u \in \mathcal{D}_{a,b}$$

$C_{a,b}$ = best constant for general functions u

$C_{a,b}^*$ = best constant for radially symmetric functions u

$$C_{a,b}^* \leq C_{a,b}$$

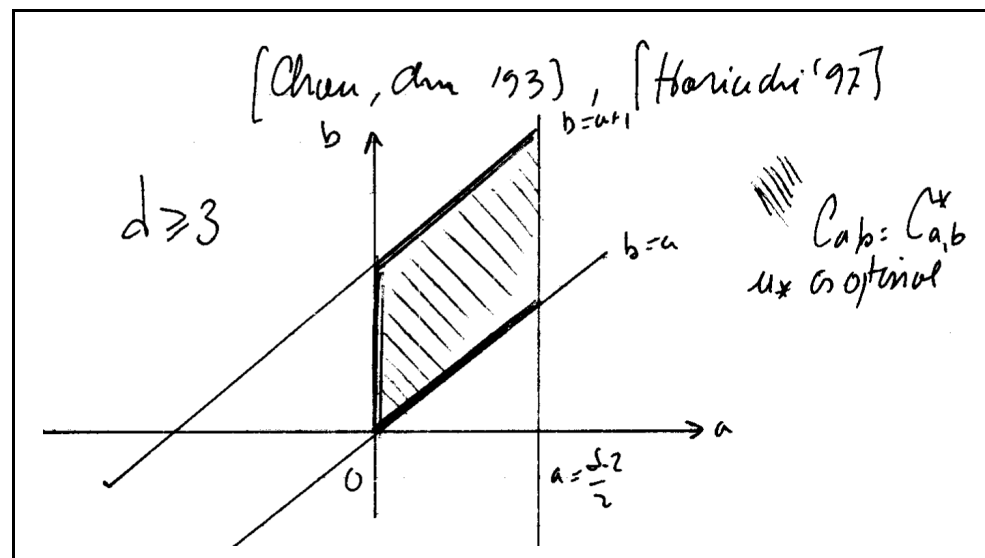
Up to scalar multiplication and dilation, the optimal radial function is

$$u_{a,b}^*(x) = \left(1 + |x|^{-\frac{2a(1+a-b)}{b-a}} \right)^{-\frac{b-a}{1+a-b}}$$

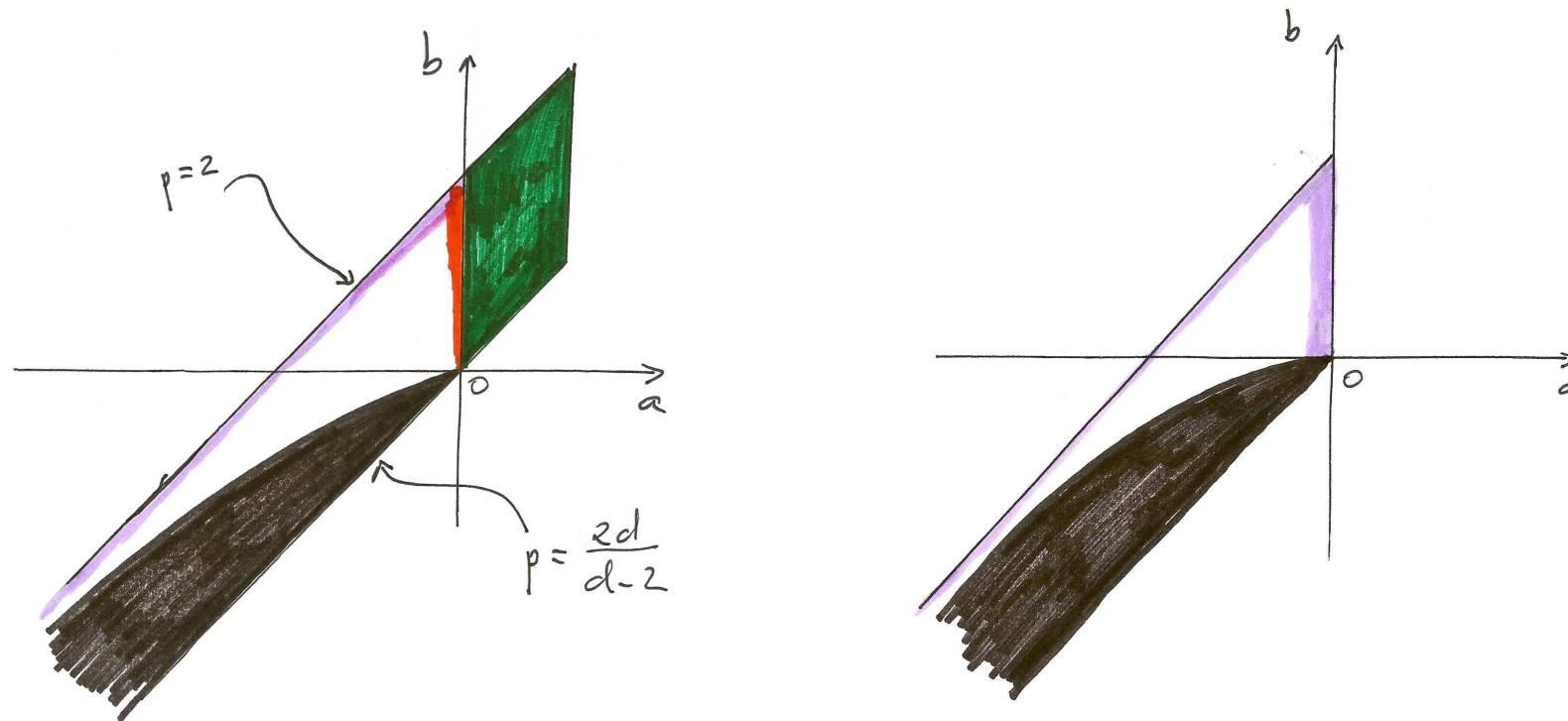
Questions: is optimality (equality) achieved ? do we have $u_{a,b} = u_{a,b}^*$?

Known results (Aubin, Talenti, Lieb, Chou-Chu, Lions, Catrina-Wang, ...)

- Existence inside the half-strip $a < b < a + 1$, $a < \frac{d-2}{2}$
- Symmetry (and existence) in the zone $a \leq b < a + 1$, $0 < a < \frac{d-2}{2}$
- Nonexistence for $a < 0$ and $b = a$ or $b = a + 1$.



Symmetry and symmetry breaking



SYMMETRY BREAKING: Catrina-Wang, Felli-Schneider.

Aubin, Talenti, Horiuchi, Lieb, Chou-Chu,...

Lin, Wang; Dolbeault, E., Tarantello ($d=2$)

Dolbeault, E., Loss, Tarantello

$$b^{FS}(a) = \frac{d(d-2-2a)}{2\sqrt{(d-2-2a)^2 + 4(d-1)}} - \frac{1}{2}(d-2-2a)$$

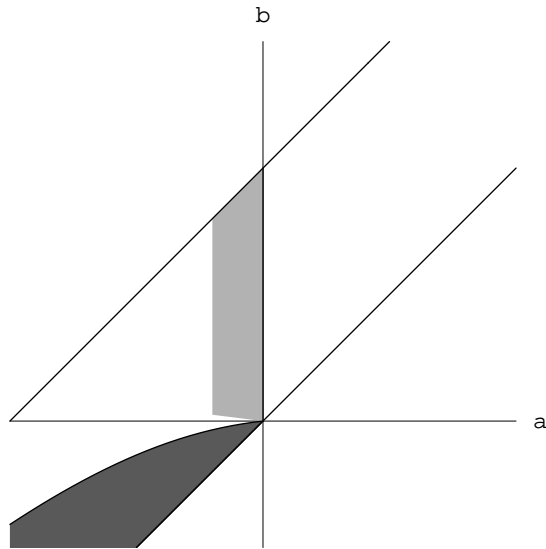
Relation with Onofri's inequality ($d = 2$) (Dolbeault, E., Tarantello)

A generalized Onofri inequality:

On \mathbb{R}^2 , consider $d\mu_\alpha = \frac{\alpha+1}{\pi} \frac{|x|^{2\alpha} dx}{(1+|x|^2(\alpha+1))^2}$ with $\alpha > -1$

$$\log \left(\int_{\mathbb{R}^2} e^u d\mu_\alpha \right) - \int_{\mathbb{R}^2} u d\mu_\alpha \leq \frac{1}{16\pi(\alpha+1)} \|\nabla u\|_{L^2(\mathbb{R}^2, dx)}^2$$

- For $d = 2$, radial symmetry holds if $-\eta < a < 0$ and $-\varepsilon(\eta) a \leq b < a + 1$



Emden-Fowler transformation and the cylinder $\mathcal{C} = \mathbb{R} \times S^{d-1}$

$$t = \log |x|, \quad \omega = \frac{x}{|x|} \in S^{d-1}, \quad v(t, \omega) = |x|^{-a} u(x), \quad \Lambda = \frac{1}{4} (d - 2 - 2a)^2$$

- Caffarelli-Kohn-Nirenberg inequalities rewritten on the cylinder become standard interpolation inequalities of Gagliardo-Nirenberg type

$$\|v\|_{L^p(\mathcal{C})}^2 \leq C_{\Lambda, p} \left[\|\nabla v\|_{L^2(\mathcal{C})}^2 + \Lambda \|v\|_{L^2(\mathcal{C})}^2 \right]$$

$$\mathcal{E}_\Lambda[v] := \|\nabla v\|_{L^2(\mathcal{C})}^2 + \Lambda \|v\|_{L^2(\mathcal{C})}^2$$

$$C_{\Lambda, p}^{-1} := C_{a, b}^{-1} = \inf \left\{ \mathcal{E}_\Lambda(v) : \|v\|_{L^p(\mathcal{C})}^2 = 1 \right\}$$

$$a < \frac{d-2}{2} \implies \Lambda > 0, \quad a < 0 \implies \Lambda > \frac{1}{4} (d-2)^2$$

Strategy of Catrina-Wang (or Felli-Schneider) for the symmetry breaking:

$$\mathcal{E}''(v_{\Lambda,p}^*)(v) = (|a|p - (1 + |a|^2) + \cdots) \|v\|_{L^2}^2 < 0$$

when v is orthogonal to all functions which do not depend on the angular variables ω .

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This is done by studying the properties of the spectrum of the operator

$$-\Delta + \Lambda - (p-1) |v_{\Lambda,p}^*|^{p-2}$$

in $H^1(\mathcal{C})$ and then using the decomposition in spherical harmonics.

Auxiliary results for the symmetry's proof

Normalize the extremal functions so that they satisfy $-\Delta v + \Lambda v = v^{p-1}$, that is

$$\int_{\mathcal{C}} |\nabla v|^2 dx + \Lambda \int_{\mathcal{C}} |v|^2 dx = \int_{\mathcal{C}} |v|^p dx$$

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Lemma. Let $d \geq 2$, $p \in (2, 2^*)$. For all $\Lambda \neq 0$, we have

$$\|v_{\Lambda,p}\|_{L^p(\mathcal{C})}^p = (C_{\Lambda,p})^{-\frac{p}{p-2}} \leq (C_{\Lambda,p}^*)^{-\frac{p}{p-2}} = \|v_{\Lambda,p}^*\|_{L^p(\mathcal{C})}^p = 4 |S^{d-1}| (2 \Lambda p)^{\frac{p}{p-2}} \frac{c_p}{2 p \sqrt{\Lambda}}$$

where c_p is a function increasing in p such that

$$\lim_{p \rightarrow 2_+} 2^{\frac{2p}{p-2}} \sqrt{p-2} c_p = \sqrt{2\pi}.$$

Proof of of the symmetry result for b near $a + 1$ $d \geq 3$

By contradiction. Suppose the existence of $(\Lambda_n)_{n \in \mathbb{N}}$ and $(p_n)_{n \in \mathbb{N}}$ s. t.

$$\lim_{n \rightarrow +\infty} \Lambda_n = \Lambda \geq (d-2)^2/4, \quad \lim_{n \rightarrow +\infty} p_n = 2_+,$$

and such that the corresponding global minimum $v_n := v_{\Lambda_n, p_n}$ satisfies

$$\mathcal{F}_{\Lambda, p}[v_{\Lambda_n, p_n}] < \mathcal{F}_{\Lambda, p}[v_{\Lambda_n, p_n}^*], \quad -\Delta_y v_n + \Lambda_n v_n = v_n^{p-1} \quad \text{in } \mathcal{C}.$$

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Define $c_n^2 = (\Lambda_n p_n)^{-\frac{p_n}{p_n-2}} 2^{\frac{p_n}{p_n-2}} \sqrt{p_n - 2}$ et $W_n := c_n v_n$.

$$-\Delta W_n + \Lambda_n W_n = c_n^{2-p_n} W_n^{p_n-1} \quad \text{in } \mathcal{C}, \quad \int_{\mathcal{C}} |\nabla W_n|^2 dy + \Lambda_n \int_{\mathcal{C}} W_n^2 dy = c_n^2 \int_{\mathcal{C}} v_n^{p_n} dy.$$

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We have : $\lim_{n \rightarrow +\infty} c_n^{2-p_n} = \Lambda$ and $\limsup_{n \rightarrow +\infty} c_n^2 \int_{\mathcal{C}} v_n^{p_n} dy \leq |S^{d-1}| \sqrt{2\pi/\Lambda}$.

Hence $(W_n)_{n \in \mathbb{N}}$ is bounded in $H^1(\mathcal{C})$

Elliptic estimates and Harnack's inequality imply that $W_n \rightarrow W$ and

$$-\Delta W + \Lambda W = \Lambda W.$$

Hence $W_n \rightarrow W \equiv 0$.

End of the proof of the symmetry result for b near $a + 1$

Differentiating with respect to ω the equation

$$-\Delta W_n + \Lambda_n W_n = c_n^{2-p_n} W_n^{p_n-1} \text{ dans } \mathcal{C},$$

we obtain:

$$-\Delta \chi_n + \Lambda_n \chi_n = (p_n - 1) c_n^{2-p_n} W_n^{p_n-2} \chi_n \quad , \quad \chi_n := \nabla_\omega v_n := (\sin \omega)^{2-d} \frac{\partial}{\partial \omega} \left((\sin \omega)^{d-2} v_n \right)$$

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$$0 = \int_{\mathcal{C}} |\nabla \chi_n|^2 dy + \Lambda_n \int_{\mathcal{C}} |\chi_n|^2 dy - (p_n - 1) c_n^{2-p_n} \int_{\mathcal{C}} W_n^{p_n-2} |\chi_n|^2 dy .$$

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Since $\int_{S^{d-1}} \nabla_\omega W_n(t, \omega) d\omega = 0$, an expansion of χ_n in spherical harmonics shows that

$$\int_{\mathcal{C}} |\nabla \chi_n|^2 dy \geq (d - 1) \int_{\mathcal{C}} |\chi_n|^2 dy .$$

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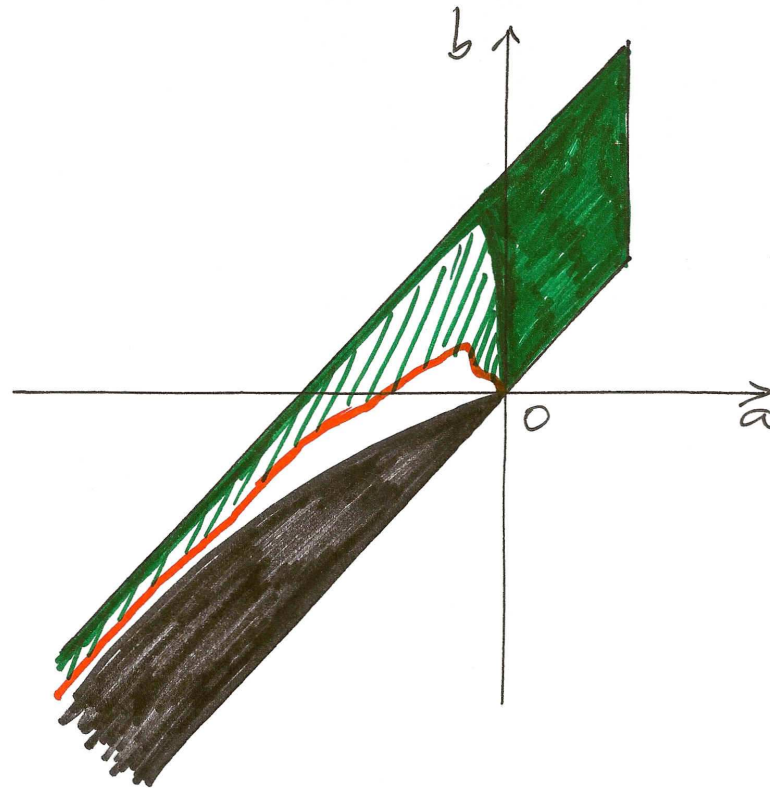
Since the functions W_n are bounded from above by $W_n(0, \omega_0)$, we obtain

$$0 \geq \left(d - 1 + \Lambda_n - (p_n - 1) c_n^{2-p_n} W_n(0, \omega_0)^{p_n-2} \right) \int_{\mathcal{C}} |\chi_n|^2 dy .$$

Contradiction: $\lim_{n \rightarrow +\infty} \Lambda_n = \Lambda$, $\limsup (p_n - 1) c_n^{2-p_n} W_n(0, \omega_0)^{p_n-2} \leq \Lambda$

Two simply connected regions separated by a continuous curve

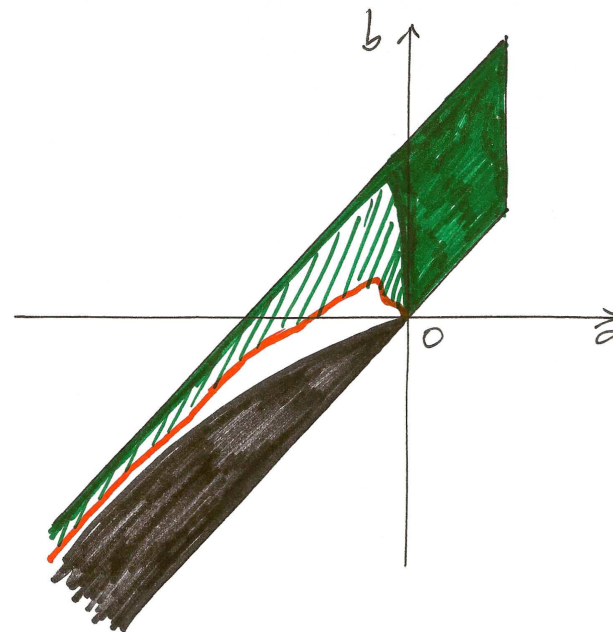
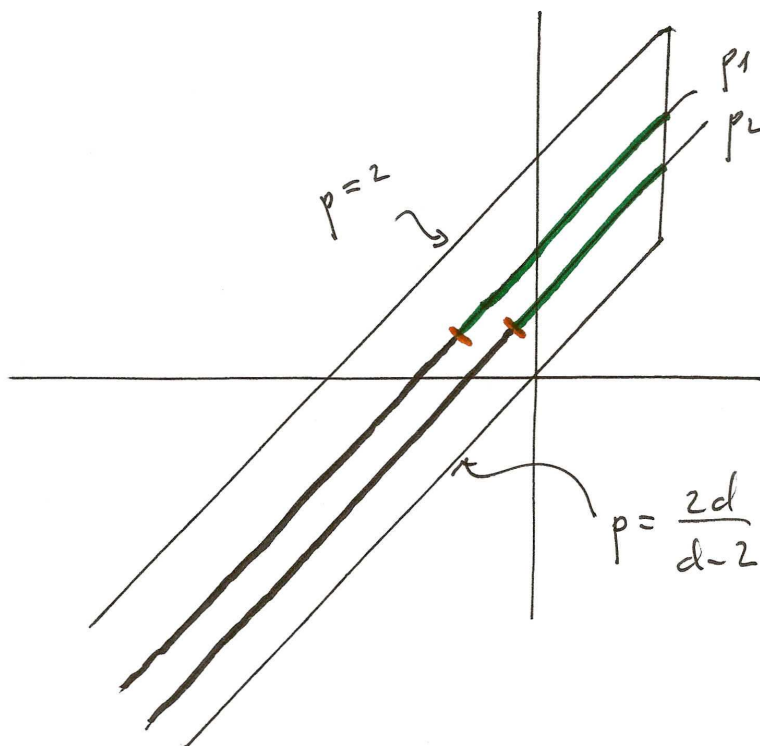
The symmetry and the symmetry breaking zones are simply connected and separated by a continuous curve.



Scaling and consequences

Let $w_\sigma(t, \omega) := w(\sigma t, \omega)$ for any $\sigma > 0$

$$\mathcal{F}_{\sigma^2 \Lambda, p}(w_\sigma) = \sigma^{1+2/p} \mathcal{F}_{\Lambda, p}(w) - \sigma^{-1+2/p} (\sigma^2 - 1) \frac{\int_{\mathcal{C}} |\nabla_\omega w|^2 dy}{(\int_{\mathcal{C}} |w|^p dy)^{2/p}}$$



Upper semicontinuity of the curve $p \rightarrow \bar{\Lambda}(p)$ is easy to prove.
For continuity, a delicate spectral analysis is needed.

CONJECTURE: $\bar{\Lambda}$ coincides with Λ_{FS} .

End of the proof of the symmetry result for b near $a + 1$

Let us now see why the identity

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implies that if for some given (Λ, p) there is a non-symmetric minimizer, then for all $\sigma > 1$), and for $(\sigma^2\Lambda, p)$ no minimizer can be symmetric.

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CONSEQUENCE: $\left(C_{\sigma^2 \Lambda, p}^*\right)^{-1} = \sigma^{1+2/p} \left(C_{\Lambda, p}^*\right)^{-1}.$

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CONSEQUENCE: $\left(C_{\sigma^2 \Lambda, p}^*\right)^{-1} = \sigma^{1+2/p} \left(C_{\Lambda, p}^*\right)^{-1}.$

Assume that $w_{\Lambda, p}$ is a non radially symmetric extremal function and apply the above identity with $w = w_{\Lambda, p}$, $w_\sigma(t, \omega) := w(\sigma t, \omega)$, $\lambda = \sigma^2 \Lambda$ and $\sigma > 1$:

$$\begin{aligned} (C_{\lambda, p})^{-1} &\leq \mathcal{F}_{\sigma^2 \Lambda, p}(w_\sigma) = \sigma^{1+\frac{2}{p}} (C_{\Lambda, p})^{-1} - \sigma^{-1+\frac{2}{p}} (\sigma^2 - 1) \frac{\int_{\mathcal{C}} |\nabla_\omega w_{\Lambda, p}|^2 dy}{\left(\int_{\mathcal{C}} |w_{\Lambda, p}|^p dy\right)^{\frac{2}{p}}} \\ &\leq \sigma^{1+\frac{2}{p}} (C_{\Lambda, p}^*)^{-1} - \sigma^{-1+\frac{2}{p}} (\sigma^2 - 1) \frac{\int_{\mathcal{C}} |\nabla_\omega w_{\Lambda, p}|^2 dy}{\left(\int_{\mathcal{C}} |w_{\Lambda, p}|^p dy\right)^{\frac{2}{p}}} < (C_{\lambda, p}^*)^{-1}, \end{aligned}$$

since $\nabla_\omega w_{\Lambda, p} \not\equiv 0$. This proves the second claim with $\lambda = \sigma^2 \Lambda$.

Generalized Caffarelli-Kohn-Nirenberg inequalities (CKN)

Let $d \geq 3$. For any $p \in [2, p(\theta, d) := \frac{2d}{d-2\theta}]$, there exists a positive constant $C(\theta, p, a)$ such that

$$\left(\int_{\mathbb{R}^d} \frac{|u|^p}{|x|^{bp}} dx \right)^{\frac{2}{p}} \leq C(\theta, p, a) \left(\int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{2a}} dx \right)^{\theta} \left(\int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} dx \right)^{1-\theta}$$

In the radial case, with $\Lambda = (a - a_c)^2$, the best constant when the inequality is restricted to radial functions is $C_{\text{CKN}}^*(\theta, p, a)$ and (see [Del Pino, Dolbeault, Filippas, Tertikas]):

$$C_{\text{CKN}}(\theta, p, a) \geq C_{\text{CKN}}^*(\theta, p, a) = C_{\text{CKN}}^*(\theta, p) \Lambda^{\frac{p-2}{2p} - \theta}$$

$$C_{\text{CKN}}^*(\theta, p) = \left[\frac{2\pi^{d/2}}{\Gamma(d/2)} \right]^{2\frac{p-1}{p}} \left[\frac{(p-2)^2}{2+(2\theta-1)p} \right]^{\frac{p-2}{2p}} \left[\frac{2+(2\theta-1)p}{2p\theta} \right]^{\theta} \left[\frac{4}{p+2} \right]^{\frac{6-p}{2p}} \left[\frac{\Gamma\left(\frac{2}{p-2} + \frac{1}{2}\right)}{\sqrt{\pi} \Gamma\left(\frac{2}{p-2}\right)} \right]^{\frac{p-2}{p}}$$

For $\theta \in (0, 1)$ fixed, $p \in [2, p(\theta, d) := \frac{2d}{d-2\theta}]$.

For $p \in [2, \frac{2d}{d-2}]$ fixed, $\theta \in [2, \theta(p, d) := \frac{d(p-2)}{2p}]$.

Existence results (Dolbeault, E.)

Let $d \geq 2$ and assume that $a \in (-\infty, \frac{d-2}{2})$.

– For any $p \in (2, 2^*)$ and any $\theta \in (\vartheta(p, d), 1)$, the Caffarelli-Kohn-Nirenberg inequality (CKN)

$$\left(\int_{\mathbb{R}^d} \frac{|u|^p}{|x|^{bp}} dx \right)^{\frac{2}{p}} \leq C(\theta, p, a) \left(\int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{2a}} dx \right)^{\theta} \left(\int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} dx \right)^{1-\theta}$$

admits an extremal function in $\mathcal{D}_a^{1,2}(\mathbb{R}^d)$

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Critical case: there exists a continuous function $a^* : (2, 2^*) \rightarrow (-\infty, a_c)$ such that the inequality also admits an extremal function in $\mathcal{D}_a^{1,2}(\mathbb{R}^d)$ if $\theta = \vartheta(p, d)$ and $a \in (a^*(p), a_c)$

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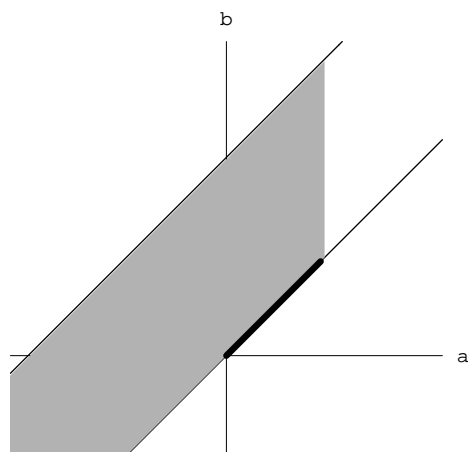
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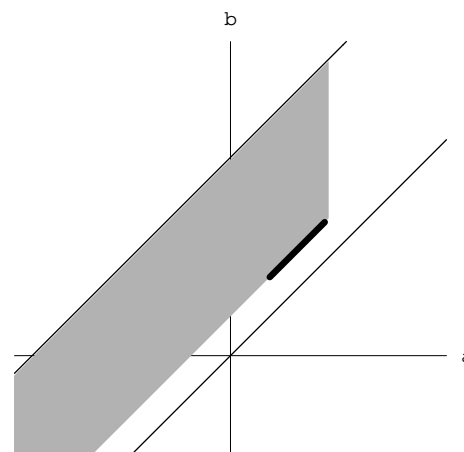
$$\left(\int_{\mathbb{R}^d} \frac{|u|^p}{|x|^{b p}} dx \right)^{\frac{2}{p}} \leq C(\theta, p, a) \left(\int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{2a}} dx \right)^{\theta} \left(\int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} dx \right)^{1-\theta}$$

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$d = 3, \theta = 1$



$d = 3, \theta = 0.8$

Strategy of the proofs

Emden-Fowler transformation and minimization on the cylinder of the functionals

$$\mathcal{E}_\theta[v] := \left(\|\nabla v\|_{L^2(\mathcal{C})}^2 + \Lambda \|v\|_{L^2(\mathcal{C})}^2 \right)^\theta \|v\|_{L^2(\mathcal{C})}^{2(1-\theta)}$$

under the constraints $\|v\|_{L^p(\mathcal{C})} = 1$.

Strategy of the proofs

Emden-Fowler transformation and minimization on the cylinder of the functionals

$$\mathcal{E}_\theta[v] := \left(\|\nabla v\|_{L^2(\mathcal{C})}^2 + \Lambda \|v\|_{L^2(\mathcal{C})}^2 \right)^\theta \|v\|_{L^2(\mathcal{C})}^{2(1-\theta)}$$

under the constraints $\|v\|_{L^p(\mathcal{C})} = 1$.

Existence: Convergence of minimizing sequences **if they are bounded in $H^1(\mathcal{C})$** : by a more or less classical use of concentration-compactness techniques.

For $p \in (2, p(\theta, d))$, existence is easy to prove, for all $\Lambda > 0$.

For $p = p(\theta, d)$, a priori estimates can be obtained by relatively elaborate direct interpolation inequalities or estimating the energy strictly from above.

A possible loss of compactness

- Gagliardo-Nirenberg interpolation inequalities: if $p \in (2, 2^*)$,

$$\|u\|_{L^p(\mathbb{R}^d)}^2 \leq C_{GN}(p) \|\nabla u\|_{L^2(\mathbb{R}^d)}^{2\vartheta(p,d)} \|u\|_{L^2(\mathbb{R}^d)}^{2(1-\vartheta(p,d))} \quad \forall u \in H^1(\mathbb{R}^d)$$

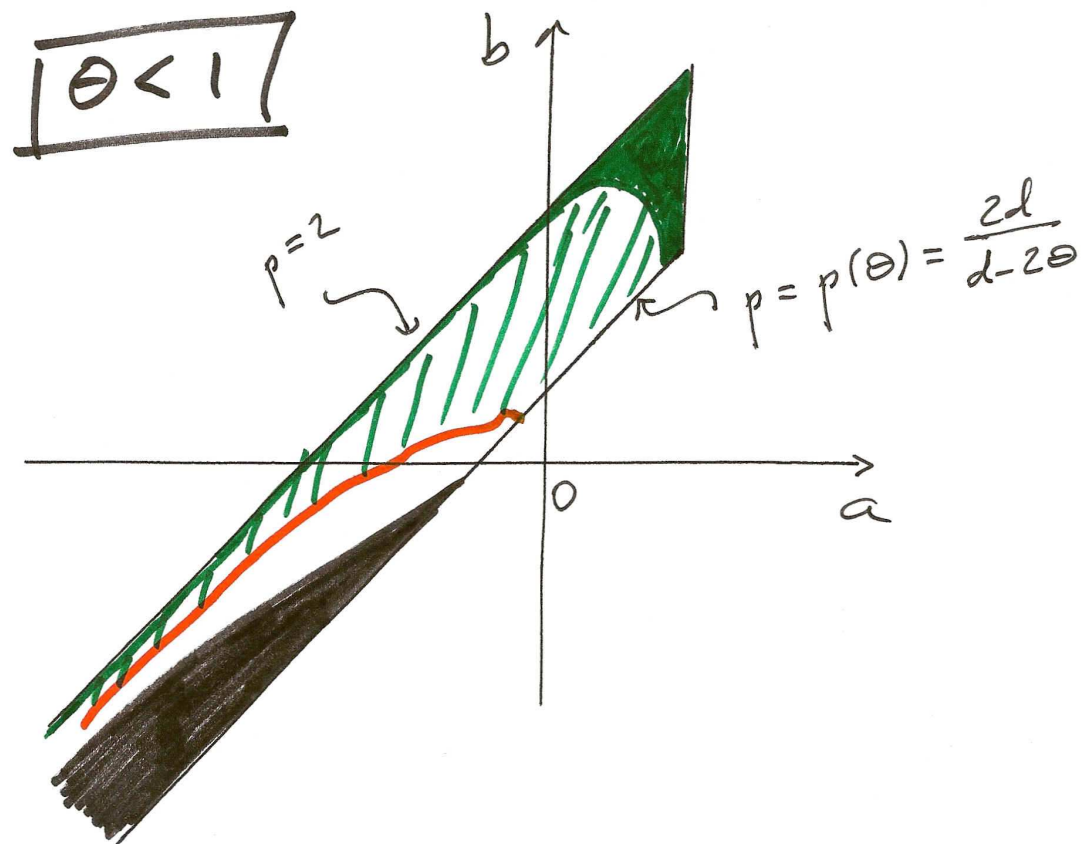
If u is a radial minimizer for $1/C_{GN}(p)$ and $u_n(x) := u(x + n\mathbf{e})$, $\mathbf{e} \in \mathbb{S}^{d-1}$

$$\begin{aligned} \frac{1}{C_{CKN}(\vartheta(p,d), p, a)} &\leq \frac{\| |x|^{-a} \nabla u_n \|_{L^2(\mathbb{R}^d)}^{2\vartheta(p,d)} \| |x|^{-(a+1)} u_n \|_{L^2(\mathbb{R}^d)}^{2(1-\vartheta(p,d))}}{\| |x|^{-b} u_n \|_{L^p(\mathbb{R}^d)}^2} \\ &= \frac{1}{C_{GN}(p)} (1 + \mathcal{R} n^{-2} + O(n^{-4})) \end{aligned}$$

So, $C_{GN} \leq C_{CKN}$.

A priori estimates (and relative compactness of minimizing sequences) can also be obtained if $C_{GN} < C_{CKN}$.

Symmetry and symmetry breaking for $\theta \in (0, 1)$



New symmetry breaking results

We already know that

$$C_{GN} \leq C_{CKN} .$$

If we are able to find a positive Λ and a function g such that

$$\frac{1}{C_{GN}} \leq \mathcal{E}_{GN}[g] < \frac{1}{C_{CKN}^*(\vartheta(p, d), p, \Lambda)} ,$$

then,

$$\frac{1}{C_{CKN}(\vartheta(p, d), p, \Lambda)} \leq \frac{1}{C_{GN}} < \frac{1}{C_{CKN}^*(\vartheta(p, d), p, \Lambda)}$$

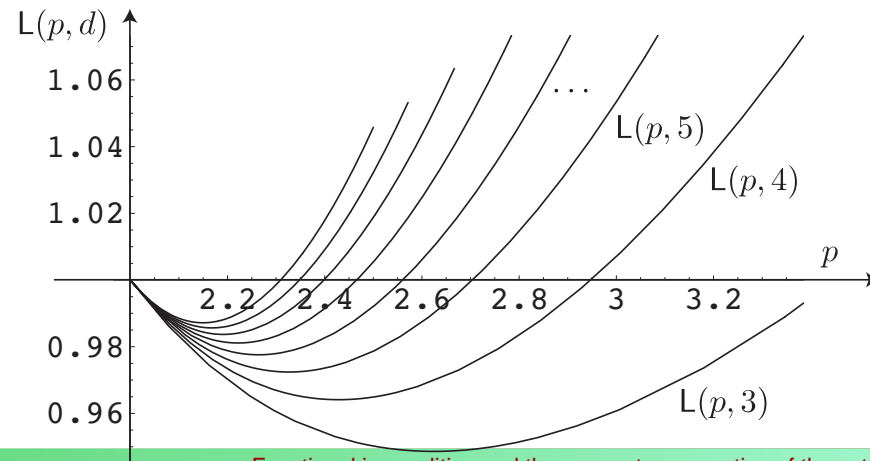
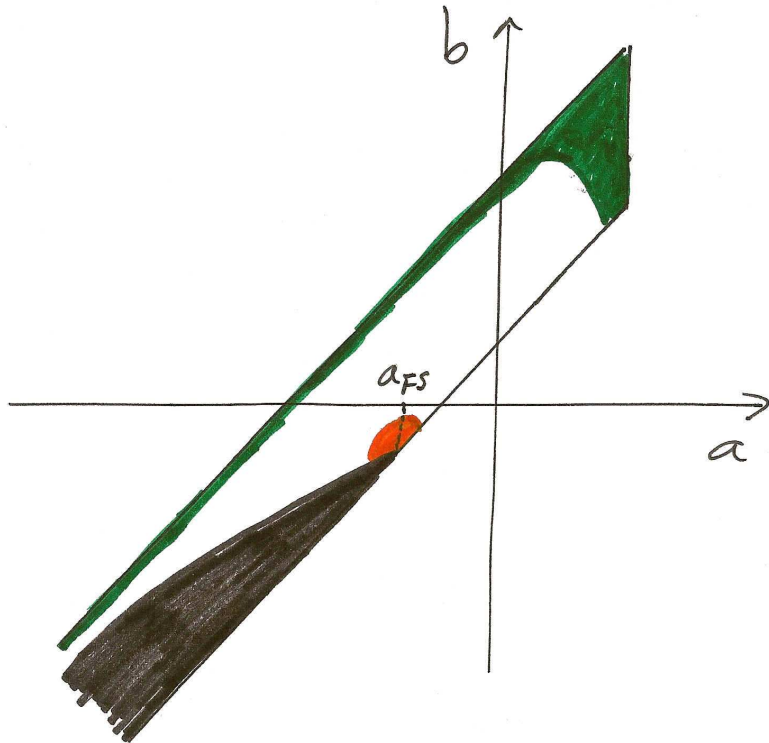
A new symmetry breaking result (2010, Dolbeault, E., Tarantello, Tertikas)

Let $g(x) := (2\pi)^{-d/4} \exp(-|x|^2/4)$. Choose $\Lambda = \Lambda_{FS}(p(\theta, d), d)$

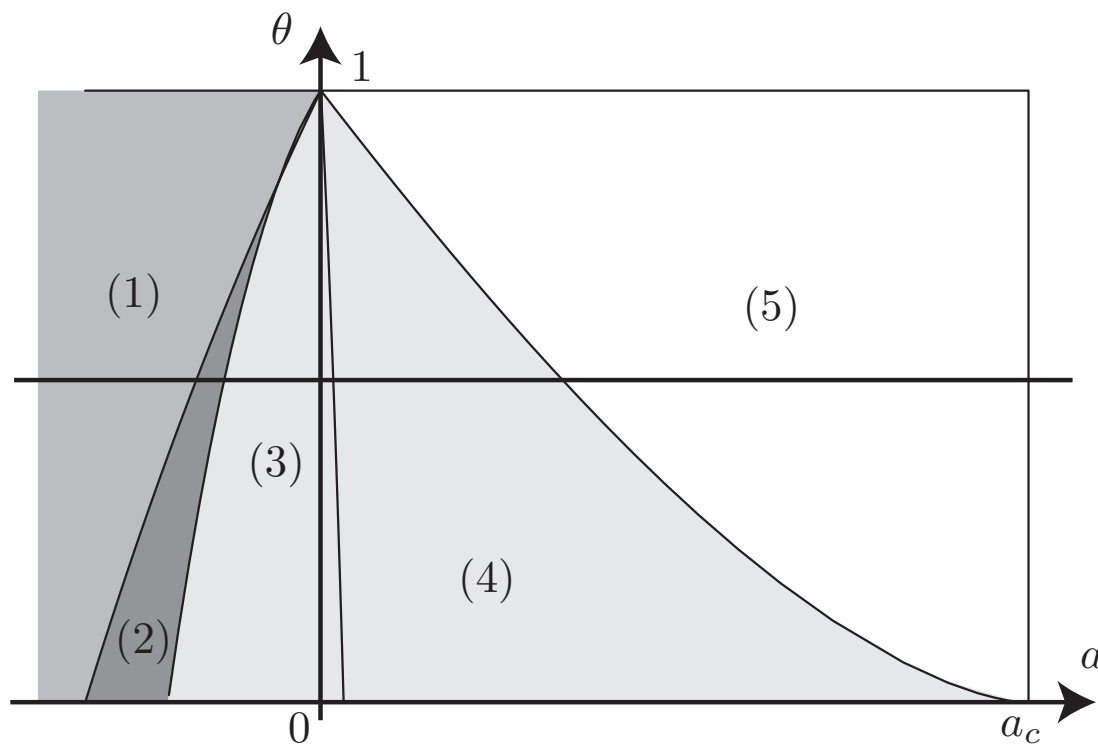
Symmetry breaking occurs if

$$L(p, d) := \frac{\mathcal{E}_{GN}[g]}{\frac{1}{C_{CKN}^*(\vartheta(p, d), p, \Lambda)}} < 1$$

We have the following result:



Analytical and numerical results



(1) Symmetry breaking by instability of the radial extremals.

(2) Symmetry breaking by comparison with GN best constant.

(5) existence obtained because of symmetry by kind of Schwarz' symmetrization.

(4) existence by a priori estimates obtained analytically.

(3) existence obtained by (strict) comparison of best constant with best constant of GN.

A new symmetry result (June 2011, Dolbeault, E., Loss)

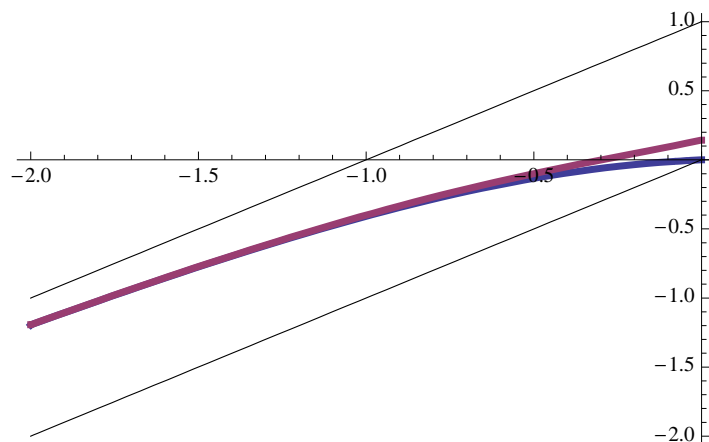
For $\theta = 1$ and $d \geq 2$,

using an elaborate argument involving Lieb-Thirring inequalities, Beckner's inequalities on the sphere and a rigidity argument due to Bidaut-Veron and Veron for positive solutions of $-\Delta_g v + \mu v = v^q$ on manifolds (M, g) ,

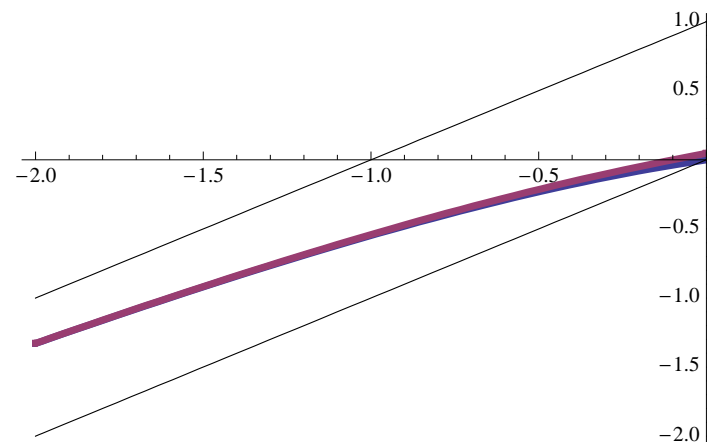
there exists a unique minimizer for the (CKN) problem, and it is symmetric, for all $\Lambda \leq \tilde{\Lambda}(p)$, for all $p \in (2, \frac{2d}{d-2})$.

$$\tilde{\Lambda}(p) := \frac{(d-1)(6-p)}{4(p-2)} < \Lambda_{FS}(p).$$

$d = 3$



$d = 5$



Strategy of the proofs

Let L^2 be the Laplace-Beltrami operator on S^{d-1} . So that $-\Delta$ on the cylinder becomes $-\partial_s^2 - L^2$.

THEOREM. Let $d \geq 2$ and let u be a non-negative function on $\mathcal{C} = \mathbb{R} \times S^{d-1}$ that satisfies

$$-\partial_s^2 v - L^2 v + \Lambda v = v^{p-1}$$

and consider the symmetric solution v_* . Assume that

$$\int_{\mathcal{C}} |v(s, \omega)|^p ds d\omega \leq \int_{\mathbb{R}} |v_*(s)|^p ds$$

for some $2 < p < 6$ satisfying $p \leq \frac{2d}{d-2}$. If $\Lambda \leq \tilde{\Lambda}(p)$, then for a.e. $\omega \in S^{d-1}$ and $s \in \mathbb{R}$, we have $v(s, \omega) = v_*(s - C)$ for some constant C .

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REMARK 1. With the above normalization, we have

$$\frac{1}{C_{\Lambda, p}} = \inf \frac{\int_{\mathcal{C}} |\nabla v|^2 + \Lambda v^2 dx}{\left(\int_{\mathcal{C}} |v|^p dx\right)^{2/p}} = \left(\int_{\mathcal{C}} |v(s, \omega)|^p ds d\omega\right)^{\frac{p-2}{p}}.$$

REMARK 2. We choose $d\omega$ to be a probability measure on S^{d-1} .

Lieb-Thirring in 1-d

LEMMA. Let $V = V(s)$ be a non-negative real valued potential in $L^{\gamma+1/2}(\mathbb{R})$ for some $\gamma > 1/2$ and let $-\lambda_1(V)$ be the lowest eigenvalue of the Schrödinger operator $-\frac{d^2}{ds^2} - V$. Define

$$c_{\text{LT}}(\gamma) = \frac{\pi^{-1/2}}{\gamma - 1/2} \frac{\Gamma(\gamma + 1)}{\Gamma(\gamma + 1/2)} \left(\frac{\gamma - 1/2}{\gamma + 1/2} \right)^{\gamma+1/2}.$$

Then

$$\lambda_1(V)^\gamma \leq c_{\text{LT}}(\gamma) \int_{\mathbb{R}} V^{\gamma+1/2}(s) ds$$

with equality if and only if, up to scalings and translations,

$$V(s) = \frac{\gamma^2 - 1/4}{\cosh^2(s)} =: V_0(s)$$

in which case

$$\lambda_1(V_0) = (\gamma - 1/2)^2.$$

Furthermore, the corresponding ground state eigenfunction is given by

$$\psi_\gamma(s) = \pi^{-1/4} \left(\frac{\Gamma(\gamma)}{\Gamma(\gamma - 1/2)} \right)^{1/2} [\cosh(s)]^{-\gamma+1/2}.$$

With $V = v^{p-2}$, the equation $-\Delta v + \Lambda v = v^{p-1}$ can be seen as the “linear” equation $-\Delta v - Vv = -\Lambda v$.

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Let us define $f(\omega) := \sqrt{\int_{\mathbb{R}} |v(s, \omega)|^2 ds}$. By the Lieb-Thirring Lemma, we find that a.e. in ω ,

$$-\Lambda \int_{\mathcal{C}} |v(s, \omega)|^2 ds d\omega = \int_{S^{d-1}} \int_{\mathbb{R}} (v_s^2 - v^p) ds d\omega + \int_{\mathcal{C}} |Lv|^2 ds d\omega$$

$$\int_{S^{d-1}} \int_{\mathbb{R}} (v_s^2 - V v^2) ds d\omega + \int_{\mathcal{C}} |Lv|^2 ds d\omega =: \mathcal{F}[v] .$$

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$$\mathcal{F}[v] \geq -c_{\text{LT}}(\gamma)^{1/\gamma} \left(\int_{\mathbb{R}} |v(s, \omega)|^p ds \right)^{1/\gamma} |f|^2 + \int_{S^{d-1}} |Lf|^2 d\omega .$$

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Now, setting $D := c_{\text{LT}}(\gamma)^{1/\gamma} \left(\int_{\mathcal{C}} v^p ds d\omega \right)^{\frac{1}{\gamma}}$, by using Hölders’s inequality, we obtain

$$\mathcal{F}[v] \geq \int_{S^{d-1}} (Lf)^2 d\omega - D \left(\int_{S^{d-1}} f^{\frac{2\gamma}{\gamma-1}} d\omega \right)^{\frac{\gamma-1}{\gamma}} =: \mathcal{E}[f] .$$

The generalized Poincaré inequality on the sphere states that for all $q \in (1, \frac{d+1}{d-3}]$,

$$\frac{q-1}{d-1} \int_{S^{d-1}} (Lf)^2 d\omega \geq \left(\int_{S^{d-1}} f^{q+1} d\omega \right)^{\frac{2}{q+1}} - \int_{S^{d-1}} f^2 d\omega$$

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Choosing $q + 1 = \frac{2\gamma}{\gamma-1} = 2 \frac{p+2}{6-p}$,

$$\mathcal{E}[f] \geq \left(\frac{d-1}{q-1} - D \right) \left(\int_{S^{d-1}} f^{q+1} d\omega \right)^{\frac{2}{q+1}} - \frac{d-1}{q-1} \int_{S^{d-1}} f^2 d\omega .$$

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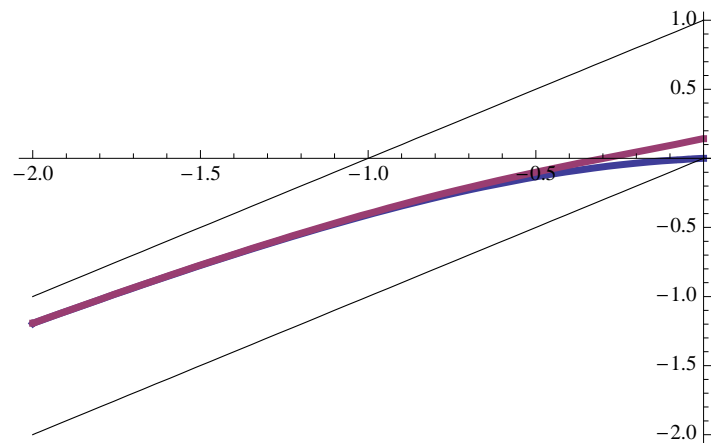
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Thus, if $D \leq \frac{d-1}{q-1}$, and if $\Lambda \leq \tilde{\Lambda}(p)$, we get

$$-\Lambda \int_{S^{d-1}} f^2 d\omega \geq \mathcal{E}[f] \geq -D \int_{S^{d-1}} f^2 d\omega \geq -\Lambda \int_{S^{d-1}} f^2 d\omega .$$

So, here we are: for $\theta = 1$,



BUT when $\theta \in (0, \theta_0)$, $\theta_0 < 1$,

we have symmetry breaking outside the Felli-Schneider zone!

Weighted logarithmic Hardy inequalities (WLH)

Let $d \geq 1$, $a < (d-2)/2$, $\gamma \geq d/4$ and $\gamma > 1/2$ if $d = 2$. Then there exists a positive constant C_{WLH} such that, for any $u \in \mathcal{D}_a^{1,2}(\mathbb{R}^d)$ normalized by $\int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} dx = 1$,

$$\int_{\mathbb{R}^d} \frac{|u|^2}{|x|^{2(a+1)}} \log \left(|x|^{d-2-2a} |u|^2 \right) dx \leq 2\gamma \log \left[C_{\text{WLH}} \int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{2a}} dx \right]$$

(Del Pino, Dolbeault, Filippas, Tertikas)

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In the radial case, the best constant when the inequality is restricted to radial functions is $C_{\text{WLH}}^*(\gamma, a)$ and $C_{\text{WLH}}(\gamma, a) \geq C_{\text{WLH}}^*(\gamma, a)$.

$$C_{\text{WLH}}^* = \frac{1}{\gamma} \frac{\left[\Gamma\left(\frac{d}{2}\right) \right]^{\frac{1}{2\gamma}}}{(8\pi^{d+1}e)^{\frac{1}{4\gamma}}} \left(\frac{4\gamma-1}{(d-2-2a)^2} \right)^{\frac{4\gamma-1}{4\gamma}} \quad \left(\gamma > \frac{1}{4} \right), \quad C_{\text{WLH}}^* = 4 \frac{\left[\Gamma\left(\frac{d}{2}\right) \right]^2}{8\pi^{d+1}e} \quad \left(\gamma = \frac{1}{4} \right)$$

If $\gamma > \frac{1}{4}$, equality is achieved by the function

$$u = \frac{\tilde{u}}{\int_{\mathbb{R}^d} \frac{|\tilde{u}|^2}{|x|^2} dx} \quad \text{where} \quad \tilde{u}(x) = |x|^{-\frac{d-2-2a}{2}} \exp \left(-\frac{(d-2-2a)^2}{4(4\gamma-1)} [\log |x|]^2 \right)$$