Matteo Bonforte

Departamento de Matemáticas, Universidad Autónoma de Madrid, Campus de Cantoblanco 28049 Madrid, Spain

email: matteo.bonforte@uam.es http://www.uam.es/matteo.bonforte

 $igg[ext{ Joint work with } extit{Alessio Figalli (UT Austin, USA)}igg]$

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Introduction

Sign-fast diffusion VS total variation flow

Sign-Fast Diffusion Equation (SFDE) as "limit" of Fast Diffusion Eq. (FDE)

$$\begin{array}{ccc} \partial_t v = \Delta \left(v^m \right) & & \partial_t v = \Delta \left(\operatorname{sign}(v) \right) \\ \text{(FDE)} & 0 < m < 1 & & \text{(SFDE)} \end{array}$$

Note that
$$v^m = |v|^{m-1}v$$
.

Total Variation Flow (TVF) as "limit" of (parabolic) p-Laplacian

$$\partial_t u = \operatorname{div} \left(|\nabla u|^{p-2} \nabla u \right)$$

$$(p-Laplacian) \quad 1
$$\partial_t u = \operatorname{div} \left(\frac{Du}{|Du|} \right)$$

$$(TVF)$$$$

Relation between TVF and SFDE in 1 spatial Dimension

If v solves the SFDE, then $u(x) := \int_0^x v(y) dy$ solves the TVF

- The above limits and relations are formal and will be justified later.
- We will consider the 1-dimensional case.



Definition of solutions

A function $u \in L^{\infty}([0,\infty),BV(I)) \cap W_{loc}^{1,2}([0,\infty),L^2(I))$ is a strong solution of the TVF, $\partial_t = \partial_x(Du/|Du|)$, if there exists $z \in L^2_{loc}([0,\infty),W^{1,2}(I))$, with $||z||_{\infty} \leq 1$, such that

$$\partial_t u = \partial_x z$$
 on $(0, \infty) \times I$,

and

$$\int_0^T \int_I z(t,x) Du(t,x) dt dx = \int_0^T \int_I |Du(t,x)| dx dt \qquad \forall T > 0.$$

- Roughly speaking, the above condition says that z = Du/|Du|.
- There is a huge literature on this topic, we refer to the book
 - F. Andreu, V. Caselles, J. M. Mazon, Parabolic quasilinear equations minimizing linear growth functionals, Progress in Mathematics, 223, Birkhäuser Verlag, Basel.

for a discussion on the different concepts of solution to the TVF. (entropy solutions, mild solutions, semigroup solution, ...)

- For the moment, we do not specify any boundary condition. The following discussion could be applied to the Cauchy problem in ℝ, as well as the Dirichlet or the Neumann problem on an interval.
- Works on TVF by: (hopeless to quote everybody, I am really sorry if I forgot someone)
 L. Ambrosio, F. Andreu, C. Ballester, G. Bellettini, V. Caselles, A. Chambolle, J. I. Diaz,
 M.-H. Giga, Y. Giga, R. Kobayashi, R. Kohn, S. Masnou, J. M. Mazon, J.-M. Morel,
 M. Novaga, P. Rybka, ...

Time Discretization

Strong solution u of the TVF is generated via Crandall-Ligget's Theorem, namely the limit of solutions of a time-discretized problem, given by the implicit Euler scheme

$$\frac{u(t_{i+1}) - u(t_i)}{t_{i+1} - t_i} = \partial_x \left(\frac{Du(t_{i+1})}{|Du(t_{i+1})|} \right). \tag{1}$$

the time-discretized solution u_h with $h = t_{i+1} - t_i > 0$ (fixed) can be characterized by

$$u_h = \operatorname{argmin} \left[\Phi_h(u) \right], \quad \text{where} \quad \Phi_h(u) = \int_I |Du| + \frac{1}{2h} \int_I |u - u_0|^2 \, \mathrm{d}x.$$

 Φ_h is strictly convex (unique minimizer), and (1) is the Euler-Lagrange eq. for Φ_h . u_0 , $u_h \in BV(I)$ implies $\partial_x z_h \in BV(I) \subset L^{\infty}(I)$, so that z_h is Lipschitz, therefore differentiable outside a countable set of points:

$$N(z_h) := \left\{ x \in \mathbb{R} \mid \lim_{\varepsilon \to 0} \frac{z_h(x + \varepsilon) - z_h(x)}{\varepsilon} \text{ does not exists} \right\}. \tag{2}$$

Finally, equation (1) is equivalent to

$$\begin{cases} h \, \partial_x z_h(x) = u_h(x) - u_0(x) & \text{for all } x \in \mathbb{R} \setminus N(z_h) \\ |z_h(x)| \le 1, & \text{for all } x \in \mathbb{R} \\ z_h(x) = \pm 1, & \text{for } |Du_h| - \text{a.e.} \end{cases}$$



Behaviour near continuity points

If u_h is different from u_0 at some common continuity point x, then it is constant in an open neighborhood of x.

Behaviour at discontinuity points (jumps decrease size in time)

Let $u_0 \in BV(I)$. Then, the following inequalities hold for any $x \in I$:

if
$$u_h(x^-) \le u_h(x^+)$$
 then $u_0(x^-) \le u_h(x^-) < u_h(x^+) \le u_0(x^+)$

$$\text{if} \qquad u_h(x^+) \leq u_h(x^-) \qquad \text{then} \qquad u_0(x^+) \leq u_h(x^+) < u_h(x^-) \leq u_0(x^-) \,.$$

Moreover,

$$u_h(x^-) < u_h(x^+)$$
 implies $z_h(x) = 1$

$$u_h(x^-) > u_h(x^+)$$
 implies $z_h(x) = -1$.

Local continuity

Let $x \in I$. If u_0 is continuous at x, then u_h is continuous at x.



The dynamics of local step functions I. The time discretized case

The dynamics of local step functions I. The time discretized case. Maximum steps.

- Let us fix an interval $I = I_1 \cup I_2 \cup I_3$
- Assume that $u_0 = \alpha_1 \chi_1 + \alpha_2 \chi_2 + \alpha_3 \chi_3$ on I with $\alpha_2 > \max\{\alpha_1, \alpha_3\}$ and $\chi_k = \chi_{I_k}$ is the char. funct. of $I_k = (x_{k-1}, x_k)$.
- We make no assumptions on u_0 outside I. Fix h > 0 small.

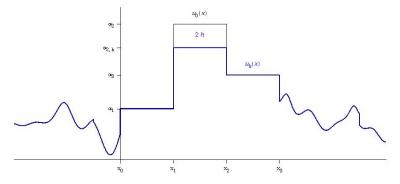


Figure: Dynamics of a maximum step. This figure shows the dynamic only inside the interval $[x_0, x_3]$.

Evolution of a general step function.

- Let $\alpha_k \in \mathbb{R}$ for $k = 0, \dots, N+1$, and $\chi_k = \chi_{I_k}$ is the char. funct. of $I_k = (x_{k-1}, x_k)$ (also the values $x_0 = -\infty$ and $x_{N+1} = +\infty$ are allowed)
- If $0 < \ell h < \min_{j=0,...,N} \left\{ \left| \alpha_j \alpha_{j+1} \right| \min \left\{ |I_j|, |I_{j+1}| \right\} \right\}$, the discrete solution after ℓ steps is given by

$$u_0 = \sum_{k=0}^{N+1} \alpha_k \chi_k$$
 gives $u_{\ell h} = \sum_{k=0}^{N+1} \alpha_{k,\ell h} \chi_k$ on I ,

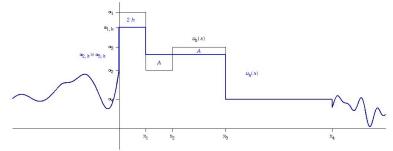
where we are able to explicitly get the values of $\alpha_{k,\ell h}$ for $k=1,\ldots N$, and some information on $\alpha_{0,\ell k}$ and $\alpha_{N+1,\ell k}$: for $k=1,\ldots N$

$$\alpha_{k,\ell h} = \left\{ \begin{array}{l} \alpha_k \,, & \text{if } \alpha_{k-1} < \alpha_k < \alpha_{k+1} \text{ or if } \alpha_{k+1} < \alpha_k < \alpha_{k-1} \\ \\ \alpha_k - \frac{2\ell h}{|I_k|} \,, & \text{if } \alpha_k > \max \left\{ \alpha_{k-1} \,,\, \alpha_{k+1} \right\} \\ \\ \alpha_k + \frac{2\ell h}{|I_k|} \,, & \text{if } \alpha_k < \min \left\{ \alpha_{k-1} \,,\, \alpha_{k+1} \right\} \end{array} \right.$$

The dynamics of local step functions I. The time discretized case

A concluding remark on the smallness of the time step h.

- Since we are mainly interested in the limit $h \to 0$, condition on smallness of h is always fulfilled.
- Anyway it is interesting to observe that the dynamic becomes more complicated to understand for general values of h, since the "locality" property is lost.
- Figure below shows a situation when a maximum and a minimum disappear in one step (for this to happen, the area A has to be less than 2h). Of course one can construct much more complicated examples.
- We can observe that the value of u_h inside $[x_1, x_2]$ depends on the values of u_0 on both $[x_1, x_2]$ and $[x_2, x_3]$.





The dynamics of local step functions II. The continuous time case.

Letting $h \to 0^+$ in the time discretized solution we obtain

$$u_0(x) = \sum_{k=0}^{N+1} \alpha_k \chi_{I_k}(x) \quad \rightsquigarrow \quad u(t,x) = u_0(x) + t \sum_{k=0}^{N+1} \beta_{k,\ell h} \chi_k(x) \quad \text{on } [0,t_1] \times I,$$

with
$$t_1 < \min_{j=0,...,N} \left\{ \left| \alpha_j - \alpha_{j+1} \right| \min \left\{ |I_j|, |I_{j+1}| \right\} \right\}$$
, and

$$\beta_{k,\ell h} := \left\{ \begin{array}{ll} 0\,, & \text{if } \alpha_{k-1} < \alpha_k < \alpha_{k+1} \text{ or if } \alpha_{k+1} < \alpha_k < \alpha_{k-1} \\ \\ -\frac{2}{|I_k|}\,, & \text{if } \alpha_k > \max\left\{\alpha_{k-1}\,,\,\alpha_{k+1}\right\} \\ \\ \frac{2}{|I_k|}\,, & \text{if } \alpha_k < \min\left\{\alpha_{k-1}\,,\,\alpha_{k+1}\right\} \end{array} \right.$$

for $k = 1, \ldots, N$, and

$$\beta_{0,\ell h} \left\{ \begin{array}{l} \geq 0 \,, \qquad \quad \text{if } \alpha_0 < \alpha_1 \\ \leq 0 \,, \qquad \quad \text{if } \alpha_0 > \alpha_1 \end{array} \right. \qquad \beta_{N+1,\ell h} \left\{ \begin{array}{l} \geq 0 \,, \qquad \quad \text{if } \alpha_N > \alpha_{N+1} \\ \leq 0 \,, \qquad \quad \text{if } \alpha_N < \alpha_{N+1}. \end{array} \right.$$

On I_0 and I_{N+1} it is monotonically increasing/decreasing, depending on the value on I_1 and I_N .

- This formula will then continue to hold until a maximum/minimum disappear.
- After repeating this at most N times, all the maxima and minima inside I disappear, and u(t) is monotonically decreasing/increasing on I.
- For instance, if $I = \mathbb{R}$ and the initial data is a compactly supported step function, then $u \equiv 0$ after some finite time T (which we call *extinction time*).
- If u_0 is an increasing (resp. decreasing) step function, then it will remain constant in time.

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with
$$t_1 < \min_{j=0,...,N} \left\{ \left| \alpha_j - \alpha_{j+1} \right| \min \left\{ |I_j|, |I_{j+1}| \right\} \right\}$$
, and

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for $k = 1, \dots, N$, and

$$\beta_{0,\ell h} \left\{ \begin{array}{l} \geq 0 \,, \qquad \quad \text{if } \alpha_0 < \alpha_1 \\ \leq 0 \,, \qquad \quad \text{if } \alpha_0 > \alpha_1 \end{array} \right. \qquad \beta_{N+1,\ell h} \left\{ \begin{array}{l} \geq 0 \,, \qquad \quad \text{if } \alpha_N > \alpha_{N+1} \\ \leq 0 \,, \qquad \quad \text{if } \alpha_N < \alpha_{N+1}. \end{array} \right.$$

On I_0 and I_{N+1} it is monotonically increasing/decreasing, depending on the value on I_1 and I_N .

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- This analysis can be extended to the case of suitable initial value problems on intervals with boundary condition:
- The dynamic of the **Dirichlet problem** is analogous to the one described above for the Cauchy problem with compactly supported initial data.
- The **Neumann problem** on some closed interval $[a, b] = I_0 \cup ... \cup I_{N+1}$. The dynamics on $I_1 \cup \ldots \cup I_N$ is known by our analysis (which, as we observed before, is "local"). Neumann condition at the level of discretized problem allows to uniquely characterize the value of u in I_0 and I_{N+1} .
 - \star For example, if $u_0 = \sum_{k=1}^{N+1} \alpha_k \chi_{I_k}$ with $\alpha_1 \leq \ldots \leq \alpha_{N+1}$ (i.e. u_0 is monotonically increasing), then

$$u(t) = u_0 + t \left(\frac{1}{|I_0|} \chi_{I_0} - \frac{1}{|I_{N+1}|} \chi_{N+1} \right)$$

(i.e. the value on I_0 increases, while the one on I_{N+1} decreases). This holds true until a jump disappears, and then one simply repeat the construction.



Theorem. (Local continuity)

Assume that u_0 is continuous on some open interval I. Then also the corresponding solution u(t) is continuous on the same interval I and the oscillation is contractive, namely

$$\sup_{I} u(t) - \inf_{I} u(t) =: \operatorname{osc}_{I} (u(t)) \leq \operatorname{osc}_{I} (u_{0}).$$

The above theorem still holds if u_0 is not continuous on I: if $u^+(t)$ and $u^-(t)$ are the solution starting respectively from

$$u^{+}(x) := \begin{cases} u_{0}(x) & \text{if } x \notin I; \\ \operatorname{esssup}_{I} u_{0} & \text{if } x \in I; \end{cases} \qquad u^{-}(x) := \begin{cases} u_{0}(x) & \text{if } x \notin I; \\ \operatorname{essinf}_{I} u_{0} & \text{if } x \in I; \end{cases}$$

then $u^-(t) \le u(t) \le u^+(t)$, $u^+(t)$ and $u^-(t)$ are both constant on I, and

$$||u^{-}(t,x) - u^{-}(t,x)||_{L^{\infty}(I)}$$
 is decreasing in time.



Further Local Properties of solutions of the TVF.

Arguing by approximation, using the stability in L^p , $1 \le p \le \infty$ we deduce *local properties of the TVF, valid on any subinterval I* where the solution u(t) is considered.

- **1** The set of discontinuity points of u(t) is contained in the set of discontinuity points of u_0 , i.e. "the TVF does not create new discontinuities".
- The number of maxima and minima decreases in time.
- If u_0 is monotone on I, then u(t) has the same monotonicity as u_0 on I. If u_0 is monotone on \mathbb{R} , then it is a stationary sol. to the Cauchy problem.
- **O** $C^{0,\alpha}$ -regularity is preserved along the flow for any $\alpha \in (0,1]$. Similar results for the denoising problem and for the Neumann problem for the TVF in V. Caselles, A. Chambolle, M. Novaga, Rev. Mat. Iberoamericana **27**, (2011). Moreover, if $u_0 \in W^{1,1}(\mathbb{R})$, then $u(t) \in W^{1,1}(\mathbb{R})$ (this is a consequence of the fact that the oscillation does not increase on any subinterval).
- **③** If $u_0 ∈ BV_{loc}(\mathbb{R})$, a priori we do not have a well-defined semigroup. However, in this case u_0 is locally bounded and the set of its discontinuity points is countable, and so in particular has Lebesgue measure zero. Then, by approximation we can still define a dynamics, which will still be contractive in any L^p space.



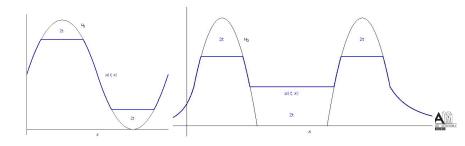
Behaviour near maxima and minima. Assume that u_0 has a local maximum at x_0 . Then, at least for short time, the solution is explicitly given near x_0 by

$$u(t,x) = \min\{u_0(x), h(t)\},\,$$

where the constant value h(t) is implicitly defined by

$$\int_{I_0} \left[u_0(x) - h(t) \right]_+ \, \mathrm{d}x = 2t \,,$$

 I_0 being the connected component of $\{u_0 > h(t)\}$ containing x_0 . The dynamics goes on in this way until a local minimum "merges" with a local maximum, and then one can simply start again the above description starting from the new configuration. For a minimum point the argument is analogous.



Loss of mass and extinction time

Let u(t) be the solution to the Cauchy problem in \mathbb{R} for the TVF, starting from a non-negative compactly supported initial datum $u_0 \in L^1(\mathbb{R})$. Then the following estimates hold:

$$\int_{\mathbb{R}} u(t,x) \, \mathrm{d}x = \int_{\mathbb{R}} u_0(x) \, \mathrm{d}x - 2t = 2(T-t) \qquad \text{for all } t \ge 0 \,,$$

and the extinction time for u is given by $T = T(u_0) = \frac{1}{2} \int_{\mathbb{R}} u_0(x) dx$.

Remark. There is no general explicit formula for the extinction time when u_0 changes sign. The rescaled flow.

We now are interested in describing the behavior of the solution near the extinction time.

We perform a logarithmic time rescaling, mapping the interval [0, T) into $[0, +\infty)$, where T is the extinction time corresponding to the initial datum u_0 . We define

$$w(s,x) = \frac{T}{T-t} u(t,x)$$
, $Z(s,x) = z(t,x)$, $s = T \log \left(\frac{T}{T-t}\right)$

where u(t) is a solution to the TVF. Then

$$\partial_s w(s,x) = \partial_x Z + \frac{w}{T}, \qquad Z \cdot D_x w = |D_x w|, \qquad w(0,x) = u_0(x).$$



Stationary solutions S(x) for the rescaled equation for w correspond to separation of variable solutions in the original variable, namely

$$-\partial_x Z = rac{S}{T}$$
 provides the separate variable solution $U_T(t,x) := rac{T-t}{T} S(x)$.

The "extended support" of a function f is the smallest interval that includes the support of f:

$$\operatorname{supp}^*(f) = \inf \{ [a, b] \mid \operatorname{supp}(f) \subseteq [a, b] \}.$$

Theorem. Stationary solutions

All compactly supported solutions of the equation $-\partial_x Z = \frac{S}{T}$, $Z \cdot D_x S = |D_x S|$, are of the form

$$S(x) = \frac{2T}{b-a} \chi_{[a,b]}(x)$$
, with $[a,b] \subseteq \mathbb{R}$.

Proposition. Mass conservation for rescaled solutions

Let w(s) be the rescaled solution, corresponding to $0 \le u_0 \in BV(\mathbb{R}) \cap L^1(\mathbb{R})$. Then

$$\int_{\mathbb{R}} w(s,x) \, dx = \int_{\mathbb{R}} u_0(x) \, dx.$$



Corollary. Separate variable solutions

All compactly supported solutions of the TVF obtained by separation of variables are of the form

$$U_T(t,x) = 2\frac{T-t}{b-a}\chi_{(a,b)}(x)$$
, where $T > 0$ and $[a,b] \subseteq \mathbb{R}$.

Proposition. Stationary solutions are asymptotic profiles

Let w(s,x) be a solution to the rescaled TVF corresponding to a non-negative initial datum $u_0 \in BV(\mathbb{R}) \cap L^1(\mathbb{R})$. Then there exists a subsequence $s_n \to \infty$ such that $w(s_n,\cdot) \to S$ in $L^1(I)$ as $n \to \infty$ where S is a stationary solution as in (1). Equivalently we have that there exists a sequence of times $t_n \to T$ as $n \to \infty$ such that

$$\left\|\frac{u(t_n,\cdot)}{T-t_n}-\frac{S}{T}\right\|_{L^1}\xrightarrow[n\to\infty]{}0.$$

where *S* is a stationary solution.

The above result has been proved by F. Andreu, V. Caselles, M. Mazon in a series of paper and in their book, for the Cauchy, Dirichlet or Neumann problem.

Theorem. Extinction profile for solutions to the TVF

Let u(t,x) be a solution to the TVF corresponding to a non negative initial datum $u_0 \in BV(\mathbb{R})$ with supp* $(u_0) = [a,b]$, and set

$$T = \frac{1}{2} \int_a^b u_0(x) \, \mathrm{d}x \, .$$

Then supp(u(t)) = [a, b] for all $t \in (0, T)$ and

$$\left\|\frac{u(t,\cdot)}{T-t}-2\frac{\chi_{[a,b]}}{b-a}\right\|_{L^1([a,b])}\xrightarrow[t\to T]{}0.$$

Remarks. The above theorem shows to important facts:

- The support of the solution becomes instantaneously the "extended support" of the initial datum, which is the support of the extinction profile.
- (ii) On $[a, b] = \text{supp}^*(u_0)$ we consider the quotient $u(t, x)/U_T(t, x)$, where U_T is the separate variable solution. Then the above convergence result can be rewritten as

$$\left\| \frac{u(t,\cdot)}{U_T(t,\cdot)} - 1 \right\|_{L^1([a,b])} \xrightarrow[t \to T]{} 0.$$
 convergence in relative error

Equivalently, L^1 -norm of the difference decays at least with the rate

$$||u(t,\cdot)-U_T(t,\cdot)||_{L^1(\mathbb{R})} \leq o(T-t).$$

We will show that the o(1) appearing in the above rate cannot be quantified/improved, so that the above convergence result is sharp, as we will see in the next slide.

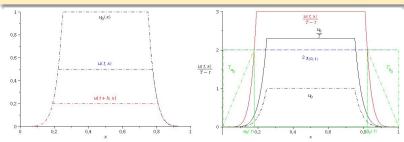
Definition. Let $\xi:[0,\infty)\to[0,\infty)$ be a continuous increasing function, with $\xi(0)=0$. We say that ξ is a *rate function* if, for any solution u(t) of the TVF,

$$\left\| \frac{u(t)}{T-t} - \frac{S}{T} \right\|_{L^1(I)} \le \xi(T-t)$$
 for any t close to the extinction time T .

Theorem. Absence of universal convergence rates

For any rate function $\xi:[0,\infty)\to[0,\infty)$, there exists an initial datum $u_0\in BV(\mathbb{R})$, with $supp^*(u_0) = [0, 1]$, such that

$$2\xi(T-t) \le \left\| \frac{u(t)}{T-t} - 2\chi_{[0,1]} \right\|_{L^1(I)}, \quad \text{for any } 0 \le T-t \le 1.$$



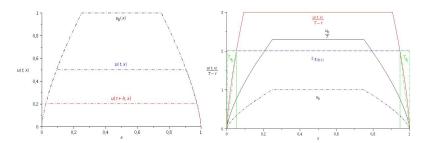
Left: Dynamic of u(t): black: $u_0(x)$, blue: u(t, x), red: u(t + h). **Right:** Rescaled dynamic: black: $u_0(x)$ (dashdot) and u_0/T (cont.), blue: u(t, x), red: u(t, x), red: u(t, x), red: u(t, x)

Remark. The above Theorem shows that there cannot be universal rates of convergence. A similar construction will provide (nontrivial) initial data for which the convergence is as fast as desired.

Fast decaying initial data

For any rate function $\xi:[0,\infty)\to[0,\infty)$, there exists an initial datum $u_0\in L^1(I)$ such that the corresponding solution u(t) satisfies

$$\left\| \frac{u(t)}{T-t} - 2\chi_{[0,1]} \right\|_{L^1(I)} \le \xi(8(T-t)), \quad \text{for any } 0 \le T - t \le 1.$$
 (4)



Solutions to the SFDE VS solutions to the TVF

- Formally: The TVF and SFDE are formally related by the fact that "u solves the TVF if and only if $D_x u$ solves the SFDE".
- In order to make this rigorous, we need first to explain what do we mean by a solution of the SFDE, and then we will prove the above relation by approximating the TVF with the p-Laplacian and the SFDE by the porous medium equation.
- The notion of solution we consider for the SFDE is the one of mild solution. We use:
 P. Benilan, M. G. Crandall, Indiana Univ. Math. J. 30 (1981), no. 2, 161–177.
- The multivalued graph of the function $r \mapsto \text{sign}(r)$ is maximal monotone (MMG).
- There exists a unique solution $u \in C([0,\infty); L^1(\mathbb{R})) \cap L^\infty([0,\infty) \times \mathbb{R})$ corresponding to the initial datum $u_0 \in L^1(\mathbb{R}) \cap L^\infty(\mathbb{R})$ that solves the problem

$$\begin{cases} u_t = \Delta \varphi(u), & \text{in } \mathcal{D}'((0, \infty) \times \mathbb{R}) \\ u(0, x) = u_0(x), & x \in \mathbb{R} \end{cases}$$

where the first equation is meant in the sense that

$$u_t = \Delta w$$
 in $\mathcal{D}'((0,\infty) \times \mathbb{R})$, with $w(t,x) \in \varphi(u(t,x))$ a.e. $t,x \in \mathbb{R}$.

• Given an approximating sequence of MMG $\varphi_n \to \varphi$ then one can prove that $u_n \to u$ in $C([0,\infty);L^1(\mathbb{R}))$.



TVF vs SFDE

Assume u_0 is a smooth compactly supported function. Let 1 , and <math>m = p - 1.

Then the following diagram is commutative:

$$T_t^p u_0 \in W^{1,p}(\mathbb{R})$$
 $\xrightarrow{p \to 1^+}$ $T_t^1 u_0 \in W^{1,1}(\mathbb{R})$

$$\begin{array}{ccc} \partial_x & & & \partial_x \\ & & & & \\ S^m_t \big(\partial_x u_0 \big) \in L^{1+m}(\mathbb{R}) & & & \longrightarrow & S^0_t \big(\partial_x u_0 \big) \in L^1(\mathbb{R}). \end{array}$$

Note that the convergence in meant in the sense of distributions

- S_t^m is the semigroup associated to the FDE equation $\partial_t v = \Delta(v^m)$. We have that $S_t^m v_0 \to S_t^0 v_0$ as $m \to 0^+$, in $C([0,\infty); L^1(\mathbb{R}))$, for any initial datum $v_0 \in L^1$.
- We can consider the p-Laplacian semigroup T_t^p for p = 1 + m. If $u_0 \in W^{1,p}(\mathbb{R})$, then $T_t^p u_0 \in W^{1,p}(\mathbb{R})$, so that as $p \to 1^+$, strong solutions to the p-Laplacian converge to strong solutions to the TVF. So that $T_t^p u_0 \to T_t^1 u_0$ in $C([0,\infty); L^1(\mathbb{R}))$ as $p \to 1^+$, where T_t^1 denotes the TVF-semigroup.
- If p=1+m, we have that $\partial_x (T_t^p u_0)$ solves (in the distributional and semigroup sense) the FDE with initial datum $\partial_x u_0$, i.e. $\partial_x (T_t^p u_0) = S_t^m (\partial_x u_0)$. Hence, by letting $m \to 0^+$, we recover such a relation in the limit p=1 and m=0.



• If $v_0 \in L^1(\mathbb{R})$, the unique mild solution of the SFDE is given by

$$S_t^0 v_0 = \partial_x \left(T_t^1 \left(\int_{-\infty}^x v_0(dy) \right) \right). \tag{5}$$

• Using (5) we can uniquely extend the generator S_t^0 to measure initial data (actually, since the semigroup T_t^1 is well-defined on $L^2(\mathbb{R})$, one could even extend the SFDE to distributional initial data in $W^{-1,2}(\mathbb{R})$).



Dynamic of the 1-dimensional SFDE

General properties of the SFDE flow. Arguing by approximation (or again using the direct relation with the TVF), as a consequence we have the following properties of the SFDE flow:

- Nonnegative initial data. Let $v_0 \ge 0$ be a locally finite measure, and define $u_0(x) := \int_{-\infty}^{x} v_0(dy) \ge 0$. Since u_0 is monotone non-decreasing, it does not evolve under the TVF. Hence v_0 is a stationary solution to the SFDE.
 - (Actually, since monotone profiles are the only stationary state for the TVF, the only stationary solutions for the SFDE are nonnegative/nonpositive initial data.)
- Only zero mean valued initial data extinguish in finite time. If v_0 is a finite measure, v(t) converges in finite time to a stationary solution \bar{v} such that $\int_{\mathbb{D}} \bar{v}(dy) = \int_{\mathbb{D}} v_0(dy)$.
 - Moreover, $\bar{v} \equiv 0$ (i.e. v_0 extinguish in finite time) if and only if $\int_{\mathbb{D}} v_0(dy) = 0$.



Dynamic of the 1-dimensional SFDE

Example 1. Delta masses as initial data.

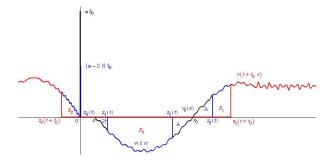
Let us assume that $v_0 = \sum_{i=1}^{N} a_i \delta_{x_i}$, with $x_1 < \cdots < x_N$. Then, for t > 0 small (dep. on $|a_i|$)

$$v(t) = \sum_{i=1}^{N} a_i(t)\delta_{x_i}, \quad \text{with } a_i(0) = a_i \text{ and}$$

$$a_{i}(t) = \begin{cases} a_{i}, & \text{if } \operatorname{sign}(a_{i-1}) = \operatorname{sign}(a_{i+1}) = \operatorname{sign}(a_{i}) \\ \operatorname{sign}(a_{i})(|a_{i}| - 4t)_{+}, & \text{if } \operatorname{sign}(a_{i-1}) = \operatorname{sign}(a_{i+1}) = -\operatorname{sign}(a_{i}) \\ \operatorname{sign}(a_{i})(|a_{i}| - 2t)_{+}, & \text{if } \operatorname{sign}(a_{i-1}) \operatorname{sign}(a_{i+1}) = -1, \end{cases}$$
(6)

where we use the convention $sign(a_0) := sign(a_1)$ and $sign(a_{N+1}) := sign(a_N)$. This formula holds true until one mass disappear at some time $t'_1 > 0$, and then it suffices to $v(t'_1)$ as initial data and repeat the construction

Example 2. Interaction between a delta and a continuous part.





Sign VS Logarithmic Fast Diffusion.

- Consider the fast diffusion equation $u_t = \Delta u^m$, with 0 < m < 1, and assume that $v_0 \ge 0$ (so $v(t) \ge 0$ for all $t \ge 0$).
- Changing the time scale $t \mapsto mt$, the above equation can be written in two different ways which lead to two different limiting equations, indeed setting $\rho(t,x) = v(t/m,x)$,

$$\frac{\partial_t v = \Delta(v^m)}{\partial_t \rho = \operatorname{div}(\rho^{m-1} \nabla \rho)} \qquad \frac{\partial_t v = \Delta(\operatorname{sign}(v))}{\partial_t \rho = \operatorname{div}(\rho^{-1} \nabla \rho)} \qquad \frac{\partial_t \rho = \operatorname{div}(\rho^{-1} \nabla \rho) = \Delta(\operatorname{log}(\rho))}{\partial_t \rho = \operatorname{div}(\rho^{-1} \nabla \rho)} = \frac{\partial_t v}{\partial_t \rho} = \frac{\partial_t v}{\partial$$

- The evolution of $\rho(t,x)$ on [0,T] corresponds to the evolution of v(t,x) on [0,T/m]. The diffusion of v is slower than the diffusion of ρ by a factor 1/m.
- So, when analyzing the limit as $m \to 0^+$, one gets two different limits: the evolution of $\rho(t, x)$ on the time interval $0 \le t \le T$ corresponds to the evolution of v(t, x) on the time interval $0 \le t < \infty$ for every T > 0.
- The solution to the SFDE corresponds to an evolution "infinitely slower" than the LFDE.

$u_0 \ge 0$ LFDE: NON-Trivial LFDE	let Problem
10 = 1	NON-trivial
$sign(u_0) = \pm 1$ SFDE: NON-trivial SFDE:	: trivial ^(**)
	NON-trivial
$sign(u_0) = \pm 1$ LFDE: NOT-possbile LFDE: 1	NOT-possbile

$$(*) u(t, \cdot) = u_0$$
, $(**)$ Immediate extinction.

Logarithmic Diffusion has been studied by A. Rodriguez, J.L. Vázquez,... and many other authors (hopeless to quote everybody).



The End

Thank you!!!

