# A perturbation argument for a Monge-Ampère type equation arising in optimal transportation 

Luis Caffarelli*<br>Univ. Texas Austin

María del Mar González ${ }^{\dagger}$<br>Univ. Politècnica de Catalunya

Truyen Nguyen ${ }^{\ddagger}$<br>Univ. of Akron


#### Abstract

We prove some interior regularity results for potential functions of optimal transportation problems with power costs. The main point is that our problem is equivalent to a new optimal transportation problem whose cost function is a sufficiently small perturbation of the quadratic cost, but it does not satisfy the well known condition (A.3) guaranteeing regularity. The proof consists in a perturbation argument from the standard MongeAmpère equation in order to obtain, first, interior $\mathcal{C}^{1,1}$ estimates for the potential and, second, interior Hölder estimates for second derivatives. In particular, we take a close look at the geometry of optimal transportation when the cost function is close to quadratic in order to understand how the equation degenerates near the boundary.


## 1 Introduction

Let $\Omega, \Omega^{\prime}$ be two bounded domains in $\mathbb{R}^{n}, n \geq 1$. Let $f, g$ be nonnegative mass distributions in $\Omega$ and $\Omega^{\prime}$, respectively, that satisfy the compatibility condition

$$
\int_{\Omega} f=\int_{\Omega^{\prime}} g
$$

Consider the optimal transport problem between $\Omega$ and $\Omega^{\prime}$ : given a continuous cost function $c: \mathbb{R}^{n} \rightarrow[0, \infty)$, find an admissible transport map $\Theta: \Omega \rightarrow \Omega^{\prime}$ that minimizes the total cost functional

$$
\int_{\Omega} c(x-\Theta(x)) f(x) d x
$$

We say that a Borel map $\Theta: \Omega \rightarrow \Omega^{\prime}$ is admissible if it is measure preserving, that is, $\Theta_{\#} f d x=g d y$. This is equivalent to

$$
\int_{\Omega} h(\Theta(x)) f(x) d x=\int_{\Omega^{\prime}} h(y) g(y) d y
$$

for every continuous function $h$. It is well known (c.f. [9], [24]) that the solution of the transport problem exists for any strictly convex cost, and it is given in terms of a potential function $u \in \mathcal{C}(\bar{\Omega})$, that satisfies a Monge-Ampère type equation

$$
\left\{\begin{align*}
g\left(x-\nabla c^{*}(-\nabla u)\right) \operatorname{det}\left(I+D^{2} c^{*}(-\nabla u) D^{2} u\right) & =f(x),  \tag{1.1}\\
\partial_{c} u(\Omega) & =\Omega^{\prime},
\end{align*}\right.
$$

[^0]where $I$ is the $n \times n$ identity matrix, $\partial_{c} u(\Omega):=\cup_{x \in \Omega} \partial_{c} u(x)$ and the $c$-subdifferential $\partial_{c} u(x)$ is defined by (2.3).

In this paper we concentrate on the convex cost $c(z)=\frac{1}{p}|z|^{p}(p>1)$ and investigate the regularity of its corresponding optimal map. First we study the asymptotic behavior of the equation when the two domains $\Omega$ and $\Omega^{\prime}$ are sufficiently far away. Second, we consider the perturbation problem $p>2, p \rightarrow 2$ for any given $\Omega$ and $\Omega^{\prime}$. In both cases, our aim is to prove interior $\mathcal{C}^{2, \alpha}$ estimates for the potential $u$. This is achieved by showing that both problems are equivalent to an optimal transportation problem whose cost function is close to the quadratic cost. We then perform a perturbation argument around the standard Monge-Ampère equation to derive regularity for $u$.

Before stating our results precisely, we recall that a bounded domain $U \subset \mathbb{R}^{n}$ is called strongly convex if $U$ has a $\mathcal{C}^{2}$ defining function $\rho$ satisfying $\left\langle D^{2} \rho(x) \xi, \xi\right\rangle>0$ for all $x \in \partial U$ and all vectors $\xi$ in the tangent space to $\partial U$ at $x$. It is also known that if $U$ is strongly convex then there exist a constant $C>0$ and a defining function $\tilde{\rho}$ for $U$ such that

$$
\begin{equation*}
\left\langle D^{2} \tilde{\rho}(x) w, w\right\rangle \geq C|w|^{2}, \quad \forall x \in \partial U \text { and } \forall w \in \mathbb{R}^{n} . \tag{1.2}
\end{equation*}
$$

For any $\delta>0$, denote

$$
\Omega_{\delta}:=\{x \in \Omega: \operatorname{dist}(x, \partial \Omega)>\delta\} .
$$

Theorem 1.1. Suppose $\Omega$ and $\Omega^{\prime}$ are two strongly convex bounded domains in $\mathbb{R}^{n}$. Assume that $f \in \mathcal{C}_{\text {loc }}^{\alpha}(\Omega)$ and $g \in \mathcal{C}_{\text {loc }}^{\alpha}\left(\Omega^{\prime}\right)$ are two nonnegative mass distributions satisfying

$$
\begin{equation*}
\int_{\Omega} f(x) d x=\int_{\Omega^{\prime}} g(y) d y, \quad \Lambda_{1} \leq f(x) \leq \Lambda_{2} \text { in } \Omega \quad \text { and } \quad \Lambda_{1} \leq g(y) \leq \Lambda_{2} \text { in } \Omega^{\prime}, \tag{1.3}
\end{equation*}
$$

where $\Lambda_{1}<\Lambda_{2}$ are positive constants. Let $u \in \mathcal{C}(\bar{\Omega})$ be the potential from the optimal transport problem (1.1) between $\Omega$ and $\Omega^{\prime}$ with cost $c(z)=\frac{1}{p}|z|^{p}, p>1$. Given any $\delta_{0}>0$, we have:

- There exists $\lambda_{0}>0$ depending on $\delta_{0}, n, p, \alpha, \Lambda_{1}, \Lambda_{2}, \Omega$ and $\Omega^{\prime}$ such that if dist $\left(\Omega, \Omega^{\prime}\right) \geq$ $\lambda_{0}$, then $u \in \mathcal{C}^{1,1}\left(\Omega_{\delta_{0}}\right)$.
- There exists $p_{0}>2$ depending on $\delta_{0}, n, \alpha, \Lambda_{1}, \Lambda_{2}, \Omega$ and $\Omega^{\prime}$ such that for all $2<p \leq p_{0}$, if $\operatorname{dist}\left(\Omega, \Omega^{\prime}\right)>C_{1}$ for some constant $C_{1}>0$, we also have that $u \in \mathcal{C}^{1,1}\left(\Omega_{\delta_{0}}\right)$.
The $\mathcal{C}^{1,1}$ norm depends on $\delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \alpha, \Omega, \Omega^{\prime},\|f\|_{\mathcal{C}^{\alpha}\left(\Omega_{\delta_{1}}\right)},\|g\|_{\mathcal{C}^{\alpha}\left(\Omega_{\delta_{1}}^{\prime}\right)}, C_{1}$. Here $\delta_{1}$ depends only on $\delta_{0}$ and universal constants.

Given the $\mathcal{C}^{1,1}$ estimate in Theorem 1.1, we could use Evans-Krylov theory for fully nonlinear, uniformly elliptic equations to obtain $\mathcal{C}^{2, \alpha}$ estimates for $u$. However, in our setting we are able to give an alternative and more geometric proof of this higher regularity result. A nice feature of our arguments is that the $\mathcal{C}^{2, \alpha}$ estimate is derived independently from the result in Theorem 1.1 and under a very weak assumption for $u$, namely, the local and global subdifferentials of $u$ agree (in particular $u \in \mathcal{C}^{1}$ is enough). We hope that this proof would be useful for other related problems. The precise statement is as follows:

Theorem 1.2. Under the same assumptions as in the previous theorem and if, in addition, the potential $u$ satisfies $\partial_{c}^{l o c} u(x)=\partial_{c} u(x)$ for all $x \in \Omega_{c \delta_{0}^{\theta}}$, then $u$ is $\mathcal{C}^{2, \alpha^{\prime}}\left(\Omega_{\delta_{0}}\right)$ for some $0<\alpha^{\prime}<\alpha$, and the $\mathcal{C}^{2, \alpha^{\prime}}$ norm depends only on the same quantities as in Theorem 1.1. Here $c, \theta$ are universal positive constants; the c-subdifferentials $\partial_{c} u(x)$ and $\partial_{c}^{l o c} u(x)$ are defined respectively by (2.3) and (2.4).

Determining the optimal regularity for the optimal transport map is a tricky question that has received a lot of attention. In the quadratic case, for the optimal transport with cost $c_{q}(z)=\frac{1}{2}|z|^{2}$, showing $\mathcal{C}^{1}$ regularity for $u$ is equivalent to proving strict convexity (see Caffarelli $[7]$ ). Then, second derivative Hölder estimates were obtained by Caffarelli in [5], where, more generally, he showed $W^{2, p}$ estimates; also the Park City notes [11] contain a sketch of the $\mathcal{C}^{1,1}$ estimates. The recent note by Gutiérrez, Huang and Nguyen [27] thoroughly explores the ideas in Caffarelli's work to give a $\mathcal{C}^{2, \alpha}$ estimate directly.

For a general convex cost, Ma-Trudinger-Wang [35] have introduced condition (A.3) that is sufficient for regularity. The example by Loeper [33] shows that the weak form of (A.3) is also necessary in the general case. Note that (A.3) and its weak form are conditions on the cost that involve fourth order derivatives. More precisely, ( $A .3$ ) is written as

$$
\begin{equation*}
\sum_{i, j, k, l, p, q, r, s}\left(c^{p, q} c_{i j, p} c_{q, r s}-c_{i j, r s}\right) c^{r, k} c^{s, l} \xi_{i} \xi_{j} \eta_{k} \eta_{l} \geq c_{0}|\xi|^{2}|\eta|^{2} \tag{1.4}
\end{equation*}
$$

for all $\xi, \eta \in \mathbb{R}^{n}$ satisfying $\xi \perp \eta$, where $c_{i, j}(x, y)=\frac{\partial^{2}}{\partial x_{i} \partial y_{j}} c(x, y),\left(c^{i, j}\right)$ is the inverse matrix of $\left(c_{i, j}\right)$ and $c_{0}$ is a positive constant. Weak- $(A .3)$ condition allows $c_{0}$ to be zero. Further references can be found in Caffarelli-Gutierrez-Huang [13], Figalli-Kim-McCann [17], FigalliLoeper [19], Kim-McCann [29], Liu [31], Liu-Trudinger-Wang [32] and Trudinger-Wang [37, $38,39]$.

Note that the power cost $c(z)=\frac{1}{p}|z|^{p}$ does not satisfy weak-(A.3) in general. This is still the case even if the two domains $\Omega$ and $\Omega^{\prime}$ are far away (see the computation after Proposition 2.1).

Thus the purpose of this paper is to treat (1.1) as a perturbation of the usual MongeAmpère equation

$$
\begin{equation*}
g(x+\nabla w) \operatorname{det}\left(I+D^{2} w\right)=f(x) \tag{1.5}
\end{equation*}
$$

Both equations are close when the distance $\lambda$ between $\Omega$ and $\Omega^{\prime}$ is very big, or when $p \rightarrow 2$. These perturbation techniques already appeared in Caffarelli's work [4] in the setting of fully nonlinear uniformly elliptic equations, (see also the book [12]). The underlying idea is that the quadratic case has such a strong geometry that is preserved under perturbation. In particular, we explore the notion of "almost" convexity in depth in order to obtain more refined information of the possible degeneracies. By comparing with the optimal transport problem for the quadratic cost, we are able to show that a priori loss of information can only happen in a very small region near the boundary.

We also need to quantify the closeness between solutions of (1.1) and solutions of the standard Monge-Ampère equation (1.5). Many times estimates are usually given in terms of good Dirichlet boundary data. To handle this, we localize the problem and work with sections. We prove a comparison principle that allows to control the sections of $u$ since they are close to sections for $w$, that have good shapes and hence their eccentricity is controlled.

Next, our estimates deteriorate at the boundary of the sections; this deterioration will be compensated with the fact that the sections become more and more round, but this geometric decay needs to be quantified explicitly.

We end the introduction by noting that the assumptions $f, g \in \mathcal{C}_{l o c}^{\alpha}$ are crucial for the conclusion in Theorem 1.1. When $f$ and $g$ are merely continuous, it is not clear how to prove that $\partial_{c}^{\text {loc }} u(x)=\partial_{c} u(x)$. As a consequence, in order to employ a perturbation method similar to our proof of Theorem 1.2 one has to deal with a difficult situation where, a priori, the restricted potential $u$ might not be a weak Alexandroff solution to the same Monge-Ampère
type equation in a smaller subregion of $\Omega$. We hope to return to this problem and to investigate interior $W^{2, p}$ estimates for the function $u$ elsewhere.

Some related works that have recently appeared are the following: regularity of optimal transport maps on manifolds (Delanoë [15], Delanöe-Ge [16], Figalli-Kim-McCann [18], Figalli-Rifford [20], Figalli-Rifford-Villani [21, 22], Kim-McCann [30], Loeper [34]), regularity for Gaussian-to-Gaussian maps with near Euclidean cost (Warren [42]). In particular, some of the ideas in Theorem 1.2 have appeared independently in De Philippis-Figalli [14].

The paper is structured as follows: Section 2 is the basic set up for our problem, together with a quick review on optimal transportation for convex costs, while Section 3 contains some geometrical properties of our special cost that will be needed at a later time. Section 4 is the starting point in the proof of the Theorem 1.1: a comparison to the optimal transport problem with quadratic cost to obtain a good initial control of the sections of $u$; Section 5 contains the main ingredient in proof of our first theorem. Next, Section 6 deals with a comparison principle for Dirichlet solutions, that allows to show that the potential $u$ is close to a quadratic polynomial locally. Finally, in the last Section 7 we put together all these results in order to complete the proof of Theorem 1.2.

Acknowledgements. The authors would like to thank the anonymous referee for valuable comments and suggestions which improved the presentation of the paper. M. d. M. González and T. Nguyen gratefully acknowledge the support provided by the Mathematical Sciences Research Institute at Berkeley and the Institute for Advanced Study at Princeton where parts of this work were carried out. T. Nguyen also would like to thank Neil Trudinger for fruitful discussions on regularity of optimal maps.

## 2 Set up

### 2.1 A review of convex costs

Here we review some standard background for equation (1.1) when the cost function $c$ : $\mathbb{R}^{n} \rightarrow[0, \infty)$ is strictly convex. The basic references for optimal transportation are the books by Ambrosio-Gigli-Savaré [1] and Villani ([40], or its most recent book [41]). In this case, it is well known that the optimal transport map exists and is unique. This was studied independently by Caffarelli [9] and Gangbo-McCann [23] following the ideas introduced by Brenier in [2] for the quadratic cost $c(z)=\frac{1}{2}|z|^{2}$. To prove the existence of such map, they considered the dual functional of Kantorovich

$$
\begin{equation*}
I(\phi, \psi)=-\int_{\Omega} f(x) \phi(x) d x-\int_{\Omega^{\prime}} g(y) \psi(y) d y \tag{2.1}
\end{equation*}
$$

for pairs $(\phi, \psi)$ of Lipschitz functions in the set

$$
K:=\left\{(\phi, \psi):-\phi(x)-\psi(y) \leq c(x-y) \quad \text { for all } \quad x \in \Omega, y \in \Omega^{\prime}\right\} .
$$

Now, it is possible to show that a maximizing pair $(u, v)$ for the Kantorovich functional exists and is unique up to a constant. Furthermore if we let

$$
\Theta_{u}(x):=x-\nabla c^{*}(-\nabla u),
$$

then this is the solution of the transport problem, where we have defined $c^{*}$, the Legendre transform of $c$, as

$$
c^{*}(y)=\sup _{x \in \mathbb{R}^{n}}\{\langle x, y\rangle-c(x)\} .
$$

The function $u: \Omega \rightarrow \mathbb{R}$ is known as the potential, and from the proof we obtain a representation of the form

$$
u(x)=\sup _{y \in \Omega^{\prime}}\{-c(x-y)-v(y)\}
$$

This potential $u$ is locally Lipschitz and satisfies a second order fully non-linear PDE of Monge-Ampère type:

$$
\begin{equation*}
M[u]=f \quad \text { in } \Omega, \tag{2.2}
\end{equation*}
$$

where

$$
M[u]:=g\left(x-\nabla c^{*}(-\nabla u)\right) \operatorname{det}\left(I+D^{2} c^{*}(-\nabla u) D^{2} u\right) .
$$

Equation (2.2) can be understood in the a.e. sense (where the gradient is an $L^{\infty}$ function), in the sense of Alexandroff, or in the viscosity setting. The book of Gutiérrez [26] gives a good introduction to these concepts in the quadratic case. Let us give here the necessary definitions for general strictly convex cost function (cf. Gutiérrez-Nguyen [28]).

Let $u$ be a function on $\Omega$. We define the (global) $c$-subdifferential of $u$ at $x_{0}$ by

$$
\begin{equation*}
\partial_{c} u\left(x_{0}\right):=\left\{y \in \mathbb{R}^{n}: u(x) \geq u\left(x_{0}\right)-c(x-y)+c\left(x_{0}-y\right) \quad \text { for all } \quad x \in \Omega\right\} \tag{2.3}
\end{equation*}
$$

and the (local) $c$-subdifferential of $u$ at $x_{0}$ as

$$
\begin{equation*}
\partial_{c}^{\mathrm{loc}} u\left(x_{0}\right):=\left\{p \in \mathbb{R}^{n}: u(x) \geq u\left(x_{0}\right)-c(x-p)+c\left(x_{0}-p\right) \text { in some neighborhood of } x_{0}\right\} . \tag{2.4}
\end{equation*}
$$

For each $u \in \mathcal{C}(\bar{\Omega})$ satisfying $\partial_{c} u(\Omega) \subset \overline{\Omega^{\prime}}$, the generalized Monge-Ampère measure of $u$ associated to the cost function $c$ and the weight $g$ is the Borel measure defined by

$$
w_{c}(g, u)(E)=\int_{\partial_{c} u(E)} g(y) d y \quad \text { for all Borel sets } \quad E \subset \Omega .
$$

If, in addition $u \in \mathcal{C}^{2}(\Omega)$ and $c^{*} \in \mathcal{C}^{2}\left(\mathbb{R}^{n}\right)$, then we can give an alternate representation

$$
\begin{equation*}
w_{c}(g, u)(E)=\int_{E \cap \Gamma_{u}} g\left(x-\nabla c^{*}(-\nabla u)\right) \operatorname{det}\left(I+D^{2} c^{*}(-\nabla u) D^{2} u\right) d x \tag{2.5}
\end{equation*}
$$

for $\Gamma_{u}=\left\{x \in \Omega: \partial_{c} u(x) \neq \emptyset\right\}$.
On the other hand, a function $u: \Omega \rightarrow \mathbb{R} \cup\{+\infty\}$ is called $c$-convex in $\Omega$ if there exists $A \subset \mathbb{R}^{n} \times \mathbb{R}$ such that

$$
u(x)=\sup _{(y, a) \in A}\{-c(x-y)+a\}
$$

for all $x \in \Omega$.
With these two ingredients in mind, we say that a $c$-convex function $u \in \mathcal{C}(\bar{\Omega})$ is a solution of equation (2.2) in the sense of Alexandroff if $\partial_{c} u(\Omega) \subset \overline{\Omega^{\prime}}$ and $w_{c}(g, u)(E)=\int_{E} f$ for every Borel set $E \subset \Omega$. It can be shown that a $c$-convex function $u$ with $\partial_{c} u(\Omega) \subset \overline{\Omega^{\prime}}$ is a solution
in the sense of Alexandroff if and only if it satisfies (2.2) a.e. in $\Omega$ and $w_{c}(g, u)$ is absolutely continuous with respect to the Lebesgue measure.

However, one of the main difficulties in this problem is that the restriction to a subdomain of an Alexandroff solution may not be an Alexandroff solution any longer which makes it extremely difficult to study regularity of such solutions. This is so because the global $c$ subdifferential (2.3) may not agree with the local one (2.4).

We need one more concept (see Ma-Trudinger-Wang [35]): a set $E^{\prime} \subset \mathbb{R}^{n}$ is $c^{*}$-convex relative to a set $E \subset \mathbb{R}^{n}$ if for each $x \in E$, the image $\left\{\nabla_{x} c(x-y): y \in E^{\prime}\right\}$ is convex. They proved that if $u$ is a generalized solution of (2.2) in the $L^{\infty}$ sense, $f>0$ in $\Omega$, and $\Omega^{\prime}$ is $c^{*}$-convex relative to $\Omega$, then $u$ is also a solution in the sense of Alexandroff.

## 2.2 "Almost" quadratic costs

In the following, we set up the proof for the first statement of Theorem 1.1. Let $\Omega$ be a bounded domain. We will assume that $\Omega^{\prime}=\lambda \mathbf{a}+\Omega^{*}$ with $\Omega^{*}$ is a bounded domain and $\mathbf{a}$ is a unit vector. Thus, $\Omega^{\prime}$ is nothing but $\Omega^{*}$ translated by a distance $\lambda>0$ in the direction of the vector $\mathbf{a}$. We will rewrite the transport equation when $\lambda \geq \lambda_{0}$ for $\lambda_{0}$ big enough in the following proposition, which plays a key role in our approach. Before stating and proving the result, we need to introduce some notations. As the matrix $I+(p-2) \mathbf{a} \otimes \mathbf{a}$ is symmetric and positive definite, there is an invertible matrix $A$ such that $I+(p-2) \mathbf{a} \otimes \mathbf{a}=A^{T} A$. Also if $E$ is a subset in $\mathbb{R}^{n}$, we shall use $A(E)$ to denote the set $\{A x: x \in E\}$.

Proposition 2.1. There exists $\lambda_{0}>0$ depending only on $\operatorname{diam}\left(\Omega \cup \Omega^{*}\right)$ such that for all $\lambda \geq \lambda_{0}$, we have: $T(x)$ is the optimal map for the transport problem between $(\Omega, f)$ and $\left(\lambda \mathbf{a}+\Omega^{*}, g\right)$ with cost $c(z)=\frac{1}{p}|z|^{p}$ if and only if $\tilde{T}(x):=A\left[T\left(A^{-1} x\right)-\lambda \mathbf{a}\right]$ is the optimal map for the transport problem between $(A(\Omega), \tilde{f})$ and $\left(A\left(\Omega^{*}\right), \tilde{g}\right)$ for the convex cost

$$
\begin{equation*}
c_{\lambda}(z)=\frac{1}{2}|z|^{2}+\frac{1}{\lambda} E(z), \tag{2.6}
\end{equation*}
$$

where the term $E$ is a smooth function satisfying the following

$$
\begin{align*}
& E(z)=O\left(|z|^{3}\right), \quad \nabla E(z)=O\left(|z|^{2}\right), \quad D^{2} E(z)=O(|z|), \\
& E(A x-A y)=\lambda^{3}\left\{\frac{1}{p}\left|\mathbf{a}-\frac{x-y}{\lambda}\right|^{p}-\frac{1}{p}+\left\langle\mathbf{a}, \frac{x-y}{\lambda}\right\rangle-\frac{1}{2}\left|\frac{A x-A y}{\lambda}\right|^{2}\right\} \quad \forall x \in \Omega, \quad \forall y \in \Omega^{*} \tag{2.7}
\end{align*}
$$

Here we have set $\tilde{f}(x)=f\left(A^{-1} x\right)$ and $\tilde{g}(y)=g\left(\lambda \mathbf{a}+A^{-1} y\right)$.
Proof. We first observe that a Borel map $y: \Omega \longrightarrow \lambda \mathbf{a}+\Omega^{*}$ satisfies $y_{\#} f d x=g d x$ if and only if $\tilde{y}: A(\Omega) \longrightarrow A\left(\Omega^{*}\right)$ satisfies $\tilde{y} \# \tilde{f} d x=\tilde{g} d x$, where $\tilde{y}(x):=A\left[y\left(A^{-1} x\right)-\lambda \mathbf{a}\right]$. Moreover, by the definition of $\tilde{f}$, we have

$$
\begin{align*}
\int_{\Omega} \frac{1}{p}|x-y(x)|^{p} f(x) d x & =\int_{\Omega} \frac{1}{p}\left|x-\lambda \mathbf{a}-A^{-1} \tilde{y}(A x)\right|^{p} \tilde{f}(A x) d x \\
& =\frac{\lambda^{p-2}}{\operatorname{det} A} \int_{A(\Omega)} \frac{\lambda^{2}}{p}\left|\mathbf{a}-\frac{A^{-1}(x-\tilde{y}(x))}{\lambda}\right|^{p} \tilde{f}(x) d x . \tag{2.8}
\end{align*}
$$

We now want to expand the integrand in the above integral. Note that $\left|A^{-1}(x-\tilde{y}(x))\right| \leq$ $c:=\operatorname{diam}\left(\Omega \cup \Omega^{*}\right)$ and hence $\left|\frac{A^{-1}(x-\tilde{y}(x))}{\lambda}\right| \leq \frac{c}{\lambda} \ll 1$ for $\lambda$ big enough. Let $w=\frac{A^{-1}(x-\tilde{y}(x))}{\lambda}$
and consider the Taylor expansion of

$$
h(w)=\frac{1}{p}|\mathbf{a}-w|^{p}
$$

around the origin:

$$
h(w)=\frac{1}{p}-\langle\mathbf{a}, w\rangle+\frac{1}{2}\langle[I+(p-2) \mathbf{a} \otimes \mathbf{a}] w, w\rangle+F(w)=\frac{1}{p}-\langle\mathbf{a}, w\rangle+\frac{1}{2}\left\langle A^{T} A w, w\right\rangle+F(w),
$$

where $F(w)=O\left(|w|^{3}\right)$. Define

$$
c_{\lambda}(z):=\frac{1}{2}|z|^{2}+\frac{1}{\lambda} E(z)
$$

with $E(z)=\lambda^{3} F\left(\frac{A^{-1} z}{\lambda}\right)$. Then it follows from the above expansion that

$$
\begin{aligned}
\frac{\lambda^{2}}{p}\left|\mathbf{a}-\frac{A^{-1}(x-\tilde{y}(x))}{\lambda}\right|^{p}= & \lambda^{2} h\left(\frac{A^{-1}(x-\tilde{y}(x))}{\lambda}\right)=\frac{\lambda^{2}}{p}-\lambda\left\langle\mathbf{a}, A^{-1} x\right\rangle+\lambda\left\langle\mathbf{a}, A^{-1} \tilde{y}(x)\right\rangle \\
& +\frac{1}{2}|x-\tilde{y}(x)|^{2}+\lambda^{2} F\left(\frac{A^{-1}(x-\tilde{y}(x))}{\lambda}\right) \\
= & \frac{\lambda^{2}}{p}-\lambda\left\langle\mathbf{a}, A^{-1} x\right\rangle+\lambda\left\langle\mathbf{a}, A^{-1} \tilde{y}(x)\right\rangle+c_{\lambda}(x-\tilde{y}(x)) .
\end{aligned}
$$

Plug this into (2.8) and by using the fact $\tilde{y}_{\#} \tilde{f}^{f} d x=\tilde{g} d x$, we obtain

$$
\begin{aligned}
\int_{\Omega} \frac{1}{p}|x-y(x)|^{p} f(x) d x= & \frac{\lambda^{p-2}}{\operatorname{det} A}\left\{\int_{A(\Omega)}\left(\frac{\lambda^{2}}{p}-\lambda\left\langle\mathbf{a}, A^{-1} x\right\rangle\right) \tilde{f}(x) d x\right. \\
& \left.+\lambda \int_{A(\Omega)}\left\langle\mathbf{a}, A^{-1} \tilde{y}(x)\right\rangle \tilde{f}(x) d x+\int_{A(\Omega)} c_{\lambda}(x-\tilde{y}(x)) \tilde{f}(x) d x\right\} \\
= & \lambda^{p-1}\left\{\int_{\Omega}\left(\frac{\lambda}{p}-\langle\mathbf{a}, x\rangle\right) f(x) d x-\int_{\lambda \mathbf{a}+\Omega^{*}}(\lambda-\langle\mathbf{a}, y\rangle) g(y) d y\right\} \\
& +\frac{\lambda^{p-2}}{\operatorname{det} A}\left\{\int_{A(\Omega)} c_{\lambda}(x-\tilde{y}(x)) \tilde{f}(x) d x\right\} .
\end{aligned}
$$

Since the first term in the last expression is independent of $y(x)$ and $\tilde{y}(x)$, the proposition is proved. Notice that the last identity in (2.7) follows from the above definition of $E$ and the fact that $F(w)=h(w)-\frac{1}{p}+\langle\mathbf{a}, w\rangle-\frac{1}{2}|A w|^{2}$.

By Proposition 2.1, it is enough to prove regularity for the potential $u$ in the optimal transportation problem between $A(\Omega)$ and $A\left(\Omega^{*}\right)$ with respect to the convex cost function $c(z):=c_{\lambda}(z)=\frac{1}{2}|z|^{2}+\frac{1}{\lambda} E(z)$, where $E$ is given by (2.7). Notice that $c$ can be rewritten as

$$
c(z)=\frac{\lambda^{2}}{p}\left|\frac{A^{-1} z}{\lambda}-\mathbf{a}\right|^{p}-\frac{\lambda^{2}}{p}+\left\langle\mathbf{a}, \frac{A^{-1} z}{\lambda}\right\rangle .
$$

Let us check the weak-(A.3) condition for this cost in dimension $n=2$. For simplicity, consider $\mathbf{a}=e_{1}=(1,0)$. In that case, the matrix $A$ is given by $A=\left(\begin{array}{cc}\sqrt{p-1} & 0 \\ 0 & 1\end{array}\right)$. Then by a straightforward but tedious calculation, we have

$$
\begin{aligned}
& \sum_{k, l, s, t} c^{k l}(z) c_{11 k}(z) c_{s t l}(z) c^{s 2}(z) c^{t 2}(z)-\sum_{s, t} c_{11 s t}(z) c^{s 2}(z) c^{t 2}(z) \\
& =\frac{(p-2) \lambda^{p-2}}{(p-1)^{3}\left(\Theta^{2}+z_{2}^{2}\right)^{\frac{p+6}{2}}}\left\{(p-1)^{2} \Theta^{6}+z_{2}^{6}-(p-1)(p+3) z_{2}^{2} \Theta^{4}-\left(2 p^{2}-3\right) z_{2}^{4} \Theta^{2}\right\},
\end{aligned}
$$

where $\Theta=\Theta(z):=\frac{z_{1}}{\sqrt{p-1}}-\lambda$. Thus for $1<p<2$,

$$
\sum_{k, l, s, t} c^{k l} c_{11 k} c_{s t l} c^{s 2} c^{t 2}-\sum_{s, t} c_{11 s t} c^{s 2} c^{t 2} \nsupseteq 0 .
$$

That is, the weak-(A.3) condition is not satisfied when $1<p<2$ no matter how big $\lambda$ is. The above computation also shows that in dimension $n=2$, if $\lambda=\operatorname{dist}\left(\Omega, \Omega^{\prime}\right)$ is sufficiently large, then the cost does satisfy weak-(A.3) when $p>2$, but not (A.3). We stress that our results are not restricted to the power costs $\frac{1}{p}|z|^{p}$ with $p>1$. Indeed:

Remark 2.2. Theorems 1.1 and 1.2 are true for any convex cost that can be written as

$$
c(z)=\frac{1}{2}|z|^{2}\left(1+\frac{F(z)}{\lambda}\right)
$$

for some $\mathcal{C}^{2}$ function $F$ and some constant $\lambda$ big enough. Looking at expression (1.4), since the quadratic cost is sharp for condition weak-(A.3), unless the perturbation term has a particular structure (involving fourth derivatives of $F$ ), this cost will not be weak-(A.3) in general.

Lemma 2.3. Let $\lambda_{0}$ be as in Proposition 2.1. Then for all $\lambda \geq \lambda_{0}$ we have: $\Omega^{\prime}$ is $c^{*}$-convex with respect to $\Omega$ if and only if $A\left(\Omega^{*}\right)$ is $c_{\lambda}^{*}$-convex with respect to $A(\Omega)$.

Proof. Let $x \in \Omega$, and consider the set

$$
E_{x}:=\left\{\nabla_{x} c(x-y): y \in \Omega^{\prime}\right\}=\left\{\nabla_{x} c(x-\lambda \mathbf{a}-y): y \in \Omega^{*}\right\} .
$$

From the proof of Proposition 2.1, we know that

$$
c(x-\lambda \mathbf{a}-y)=\frac{\lambda^{p}}{p}-\lambda^{p-1}\langle\mathbf{a}, x-y\rangle+\frac{\lambda^{p-2}}{2}\left\langle A^{T} A(x-y), x-y\right\rangle+\lambda^{p} F\left(\frac{x-y}{\lambda}\right) .
$$

Therefore we obtain

$$
E_{x}=-\lambda^{p-1} \mathbf{a}+\lambda^{p-2} A^{T}\left(\left\{A(x-y)+\lambda\left(A^{T}\right)^{-1} \nabla F\left(\frac{x-y}{\lambda}\right): y \in \Omega^{*}\right\}\right) .
$$

On the other hand, if we let $z=A x$ and consider the set

$$
K_{A x}=K_{z}:=\left\{\nabla_{z} c_{\lambda}(z-w): w \in A\left(\Omega^{*}\right)\right\},
$$

then since $c_{\lambda}(z)=\frac{1}{2}|z|^{2}+\lambda^{2} F\left(\frac{A^{-1} z}{\lambda}\right)$, we have

$$
\begin{aligned}
K_{A x} & =\left\{z-w+\lambda\left(A^{-1}\right)^{T} \nabla F\left(\frac{A^{-1}(z-w)}{\lambda}\right): w \in A\left(\Omega^{*}\right)\right\} \\
& =\left\{A(x-y)+\lambda\left(A^{T}\right)^{-1} \nabla F\left(\frac{x-y}{\lambda}\right): y \in \Omega^{*}\right\} .
\end{aligned}
$$

It follows that

$$
E_{x}=-\lambda^{p-1} \mathbf{a}+\lambda^{p-2} A^{T}\left(K_{A x}\right),
$$

and hence the set $E_{x}$ is convex if and only if the set $K_{A x}$ is convex.
We will also need:

Lemma 2.4. The potential $u$ in the optimal transportation problem of Theorem 1.1 is a weak solution in the Alexandroff sense, for $\lambda$ big enough depending on the convexity of $\Omega^{\prime}$ and $p$. Moreover, $\partial_{c}^{l o c} u(\Omega) \subset \overline{\Omega^{\prime}}$.

Proof. Since $\Omega^{\prime}$ is strongly convex by the assumption, so is $\Omega^{*}$. Let $\tilde{\rho}$ be a $\mathcal{C}^{2}$ defining function for $\Omega^{*}$ such that (1.2) holds. Then $\rho(x):=\tilde{\rho}\left(A^{-1} x\right)$ is a defining function for $A\left(\Omega^{*}\right)$ satisfying

$$
\begin{aligned}
\left\langle D^{2} \rho(x) \xi, \xi\right\rangle & =\left\langle D^{2} \tilde{\rho}\left(A^{-1} x\right) A^{-1} \xi, A^{-1} \xi\right\rangle \\
& \geq C\left\langle A^{-1} \xi, A^{-1} \xi\right\rangle=C\left\langle\left(A^{-1}\right)^{t} A^{-1} \xi, \xi\right\rangle \geq c_{p}|\xi|^{2}
\end{aligned}
$$

for all $x \in \partial A\left(\Omega^{*}\right)$ and all $\xi \in \mathbb{R}^{n}$. Thus $A\left(\Omega^{*}\right)$ is also a strongly convex domain. Next we claim that $A\left(\Omega^{*}\right)$ is $c_{\lambda}^{*}$-convex with respect to $A(\Omega)$ for $\lambda$ big enough (depending on the convexity of $\Omega^{\prime}$ and $p$ ). Indeed, $\nabla c_{\lambda}(x-y)=x-y+\frac{1}{\lambda} O\left(|x-y|^{2}\right)$ so the image $\left\{\nabla_{x} c_{\lambda}(x-y): y \in A\left(\Omega^{*}\right)\right\}$ is the same as $x-U_{x}$ with $U_{x}:=\left\{y-\frac{1}{\lambda} O\left(|x-y|^{2}\right): y \in A\left(\Omega^{*}\right)\right\}$. Let $F(z)$ be the inverse of the map $y \longmapsto y-\frac{1}{\lambda} O\left(|x-y|^{2}\right)$. Then $\bar{\rho}(z):=\rho(F(z))$ is a defining function for $U_{x}$. Since $F(z)=z+\frac{1}{\lambda} O\left(|x-z|^{2}\right), D F(z)=I+\frac{1}{\lambda} O(|x-z|), D^{2} F_{k}(z)=\frac{1}{\lambda} O(1)$ and

$$
\left\langle D^{2} \bar{\rho}(z) \xi, \xi\right\rangle=\left\langle D^{2} \rho(F(z)) D F(z) \xi, D F(z) \xi\right\rangle+\sum_{k=1}^{n} \frac{\partial \rho}{\partial y_{k}}(F(z))\left\langle D^{2} F_{k}(z) \xi, \xi\right\rangle
$$

we infer that for all $\lambda>0$ sufficiently large

$$
\left\langle D^{2} \bar{\rho}(z) \xi, \xi\right\rangle \geq c_{p}\left\langle D F(z)^{t} D F(z) \xi, \xi\right\rangle-\frac{C}{\lambda}|\xi|^{2} \geq\left(\frac{c_{p}}{2}-\frac{C}{\lambda}\right)|\xi|^{2} \geq 0 \quad \forall z \in \partial U_{x}, \forall \xi \in \mathbb{R}^{n} .
$$

Therefore the connected set $U_{x}$ is convex and the claim is proved. Consequently, $\Omega^{\prime}$ is $c^{*}$ convex with respect to $\Omega$ by Lemma 2.3. It then follows from [35] that $\partial_{c}^{\operatorname{loc}} u(\Omega) \subset \overline{\Omega^{\prime}}$ and $u$ is an Alexandroff solution.

Let us summarize our problem: $u \in \mathcal{C}(\bar{\Omega})$ is a weak solution in the sense of Alexandroff of the equation

$$
\left\{\begin{align*}
g\left(x-\nabla c^{*}(-\nabla u)\right) \operatorname{det}\left(I+D^{2} c^{*}(-\nabla u) D^{2} u\right) & =f \text { in } \Omega,  \tag{2.9}\\
\partial_{c} u(\Omega) & =\Omega^{\prime},
\end{align*}\right.
$$

for the cost $c=c_{\lambda}$, when $\lambda$ is big enough. This PDE has a very precise structure. Indeed, consider the operator

$$
M_{\lambda}[u]=g\left(x-\nabla c^{*}(-\nabla u)\right) \operatorname{det}\left(I+D^{2} c^{*}(-\nabla u) D^{2} u\right) .
$$

Because of the expression for the cost (2.6), it can be rewritten as

$$
\begin{equation*}
M_{\lambda}[u]=g(x+\nabla u+G(\nabla u)) \operatorname{det}\left(I+D^{2} u+H(\nabla u)\left(D^{2} u\right)\right), \tag{2.10}
\end{equation*}
$$

where $G$ and $H$ in particular satisfy

$$
\begin{equation*}
\|H(\nabla u)\|_{L^{\infty}} \leq \frac{C}{\lambda}\|\nabla u\|_{L^{\infty}(\Omega)}, \quad\|G(\nabla u)\|_{L^{\infty}} \leq \frac{C}{\lambda}\|\nabla u\|_{L^{\infty}(\Omega)}^{2} . \tag{2.11}
\end{equation*}
$$

We will need one more important concept for a $c$-convex function $u \in \mathcal{C}(\bar{\Omega})$ : the notion of sections. Let $y_{0} \in \partial_{c} u(\Omega)$. Define $\bar{u}(x):=u(x)+c\left(x-y_{0}\right)$ and let $x_{0} \in \Omega$ be a minimum
point of $\bar{u}$ in $\Omega$. We note that this choice of $x_{0}$ is equivalent to assuming that $y_{0} \in \partial_{c} u\left(x_{0}\right)$ since

$$
u(x)+c\left(x-y_{0}\right) \geq u\left(x_{0}\right)+c\left(x_{0}-y_{0}\right) \text { for all } x \in \Omega .
$$

In particular, if $u$ is differentiable at $x_{0}$ then $y_{0}:=x_{0}-\nabla c^{*}\left(-\nabla u\left(x_{0}\right)\right)$. For any $\mu>0$ define the $\mu$-section of $u$ at $x_{0}$ as

$$
\begin{equation*}
S_{\mu}\left(u, x_{0}, y_{0}\right):=\left\{x \in \Omega: u(x) \leq u\left(x_{0}\right)-c\left(x-y_{0}\right)+c\left(x_{0}-y_{0}\right)+\mu\right\} . \tag{2.12}
\end{equation*}
$$

In order to simplify the notation, we simply write $S_{\mu}\left(u, x_{0}\right)$ for the section, while the dependence on $y_{0}$ is made implicitly. With some abuse of notation, this section can be rewritten as

$$
\begin{equation*}
S_{\mu}\left(\bar{u}, x_{0}\right)=\left\{x \in \Omega: \bar{u}(x) \leq \bar{u}\left(x_{0}\right)+\mu\right\} . \tag{2.13}
\end{equation*}
$$

We would like to study its geometrical properties when $\mu \rightarrow 0$.
In our last remark here, we provide the main idea for the proof of the second statement of Theorem 1.1. When $p>2, p \rightarrow 2$, we have the following estimates for the cost $c(z)=\frac{1}{p}|z|^{p}$. For simplicity, assume that $1<\operatorname{dist}\left(\Omega, \Omega^{\prime}\right)<C$. Then

$$
\begin{aligned}
& \frac{1}{p}|z|^{2} \leq c(z) \leq \frac{1}{p}|z|^{2}\left[1+(p-2)(\log |z|)^{p-1}\right], \\
& \frac{2}{p} z_{i} \leq \nabla_{i} c(z) \leq \frac{2}{p} z_{i}\left[1+(p-2)(\log |z|)^{p-1}\right]
\end{aligned}
$$

and

$$
\frac{2}{p} \delta_{i j}+(p-2)\left[|z|^{p-4} z_{i} z_{j}\right] \leq D_{i j}^{2} c(z) \leq \frac{2}{p} \delta_{i j}+(p-2)\left[\delta_{i j}(\log |z|)^{p-1}+|z|^{p-4} z_{i} z_{j}\right] .
$$

In particular, there exists $p_{0}>2$, close to 2 , and depending only on the distance between $\Omega$ and $\Omega^{\prime}$, such that we are in the same situation as in the first statement of Theorem 1.1.

## 3 On the geometry of the equation

In this section we summarize some miscellaneous geometric properties that will be needed for the proof of the main theorems, in particular, a precise understanding of the almost convexity of the cost.

### 3.1 Geometric interpretation of the cost $c_{\lambda}$ - almost convexity

Let $u \in \mathcal{C}(\bar{\Omega})$ be a $c$-convex function in $\Omega$. Fix $x_{0} \in \Omega$ and $y_{0} \in \partial_{c} u\left(x_{0}\right)$. Then

$$
\begin{equation*}
u(x) \geq-c\left(x-y_{0}\right)+c\left(x_{0}-y_{0}\right)+u\left(x_{0}\right) \quad \text { for all } \quad x \in \Omega . \tag{3.1}
\end{equation*}
$$

If $c=c_{q}$ is precisely the quadratic cost, then (3.1) reduces to

$$
\begin{equation*}
u(x) \geq-\frac{1}{2}|x|^{2}+\left\langle x-x_{0}, y_{0}\right\rangle+\frac{1}{2}\left|x_{0}\right|^{2}+u\left(x_{0}\right) \quad \text { for all } \quad x \in \Omega, \tag{3.2}
\end{equation*}
$$

i.e, $\bar{u}(x):=u(x)+\frac{1}{2}|x|^{2}$ is a convex function. If $c=c_{\lambda}$ as defined in (2.6), then (3.1) reads

$$
\begin{equation*}
u(x) \geq-\frac{1}{2}|x|^{2}+\left\langle x-x_{0}, p_{0}\right\rangle-\frac{1}{2 \lambda}\left\langle D^{2} E\left(\zeta-y_{0}\right)\left(x-x_{0}\right), x-x_{0}\right\rangle+\frac{1}{2}\left|x_{0}\right|^{2}+u\left(x_{0}\right), \tag{3.3}
\end{equation*}
$$

where $\zeta$ belongs to the segment $\left[x, x_{0}\right]$ and $p_{0}:=y_{0}-\frac{1}{\lambda} \nabla E\left(x_{0}-y_{0}\right)$. Essentially, it is almost quadratic as in (3.2) plus a small perturbation that goes to zero as $\lambda \rightarrow \infty$.

Next, we fix $x_{0} \in \Omega, y_{0} \in \partial_{c} u\left(x_{0}\right)$. Take any other pair $\tilde{x} \in \Omega, \tilde{y} \in \partial_{c} u(\tilde{x})$. Then there is a supporting cost that touches $u$ from below at $\tilde{x}$, and it is written as

$$
z=u(\tilde{x})-c(x-\tilde{y})+c(\tilde{x}-\tilde{y}) .
$$

Now we let

$$
\begin{equation*}
\bar{u}(x):=u(x)+c\left(x-y_{0}\right) . \tag{3.4}
\end{equation*}
$$

The supporting cost at $\tilde{x}$ for $\bar{u}$ is given by

$$
\begin{align*}
z & =\bar{u}(\tilde{x})-c(x-\tilde{y})+c(\tilde{x}-\tilde{y})-c\left(\tilde{x}-y_{0}\right)+c\left(x-y_{0}\right) \\
& =\bar{u}(\tilde{x})+\langle\tilde{p}, x-\tilde{x}\rangle+\langle\tilde{Q}(x-\tilde{x}), x-\tilde{x}\rangle \tag{3.5}
\end{align*}
$$

where $\tilde{p}:=\nabla c\left(\tilde{x}-y_{0}\right)-\nabla c(\tilde{x}-\tilde{y})$ and $\tilde{Q}=\frac{1}{2}\left[D^{2} c\left(\xi-y_{0}\right)-D^{2} c(\xi-\tilde{y})\right]$ for some $\xi \in[x, \tilde{x}]$.
Since our cost is of the form $c(z)=\frac{1}{2}|z|^{2}+\frac{1}{\lambda} E(z)$, an estimate for $\tilde{Q}$ is

$$
\begin{equation*}
\|\tilde{Q}\|=\frac{1}{2 \lambda}\left\|D^{2} E\left(\xi-y_{0}\right)-D^{2} E(\xi-\tilde{y})\right\| \leq \frac{C}{\lambda}\left|\tilde{y}-y_{0}\right| . \tag{3.6}
\end{equation*}
$$

On the other hand, by its construction

$$
\tilde{p}=\nabla c\left(\tilde{x}-y_{0}\right)-\nabla c(\tilde{x}-\tilde{y})=\tilde{y}-y_{0}+\frac{1}{\lambda}\left[\nabla E\left(\tilde{x}-y_{0}\right)-\nabla E(\tilde{x}-\tilde{y})\right],
$$

i.e.,

$$
\begin{equation*}
\tilde{p}=\left(\tilde{y}-y_{0}\right)\left(1+O\left(\frac{C}{\lambda}\right)\right), \tag{3.7}
\end{equation*}
$$

where the constants in both (3.6) and (3.7) depend on $C=C\left(y_{0}, \tilde{y}\right)$.
If the cost is quadratic, then $Q \equiv 0$, and the supporting cost for $\bar{u}$ is just a hyperplane; this is the standard construction for convex functions. We have proved that:

Lemma 3.1. Let $u \in \mathcal{C}(\bar{\Omega})$ be c-convex and $y_{0} \in \partial_{c} u\left(x_{0}\right)$ for some $x_{0} \in \Omega$. Then the function $\bar{u}$ defined by (3.4) is almost a convex function, precisely, at any point $\tilde{x} \in \Omega$, there exists (at least) one supporting almost-hyperplane of the form

$$
z=\bar{u}(\tilde{x})+\langle\tilde{p}, x-\tilde{x}\rangle+\langle\tilde{Q}(x-\tilde{x}), x-\tilde{x}\rangle,
$$

where $\tilde{p}$ and $\tilde{Q}$ are given by (3.7) and (3.6), respectively.
In the light of the previous arguments, it is expected that a section $S$ of $u$, as defined in (2.12), is not going to be too far from its convex hull. For the rest of the paper, we will be using the notation $[A]$ to denote the convex hull of the set $A$.

Proposition 3.2. Fix the cost $c=c_{\lambda}$. Suppose that $u \in \mathcal{C}(\bar{\Omega})$ is c-convex satisfying $\partial_{c} u(\Omega) \subset$ $\Omega^{\prime}$. Let $x_{1} \in \Omega, y_{1} \in \partial_{c} u\left(x_{1}\right), \mu>0$, and consider the section $S:=S_{\mu}\left(u, x_{1}, y_{1}\right)$ as defined in (2.12). Assume that $S \Subset \Omega$. Then there exists a constant $C>0$ such that

$$
\begin{equation*}
\operatorname{dist}(x, \partial[S]) \leq \frac{C}{\lambda} \text { for all } x \in \partial S \tag{3.8}
\end{equation*}
$$

The constant $C$ depends on $\Omega$ and $\partial_{c} u(\Omega)$, but not on $\lambda$.

Proof. Let $x_{0} \in \partial S \subset \Omega$, and $y_{0} \in \partial_{c} u\left(x_{0}\right)$. Then we have for all $x$ in $\Omega$,

$$
\begin{aligned}
u(x) & \geq u\left(x_{0}\right)-c\left(x-y_{0}\right)+c\left(x_{0}-y_{0}\right) \\
& =u\left(x_{1}\right)-c\left(x_{0}-y_{1}\right)+c\left(x_{1}-y_{1}\right)+\mu-c\left(x-y_{0}\right)+c\left(x_{0}-y_{0}\right) .
\end{aligned}
$$

Thus it follows from the definition of $S$ that for all $x \in S$,

$$
\begin{aligned}
0 & \geq c\left(x-y_{1}\right)-c\left(x_{0}-y_{1}\right)+c\left(x_{0}-y_{0}\right)-c\left(x-y_{0}\right) \\
& =\left\langle\nabla c\left(x_{0}-y_{1}\right)-\nabla c\left(x_{0}-y_{0}\right), x-x_{0}\right\rangle+\frac{1}{2}\left\langle\left[D^{2} c\left(\zeta-y_{1}\right)-D^{2} c\left(\zeta-y_{0}\right)\right]\left(x-x_{0}\right), x-x_{0}\right\rangle
\end{aligned}
$$

for some $\zeta$ in the segment joining $x$ to $x_{0}$. Or

$$
0 \geq\left\langle p, x-x_{0}\right\rangle+\frac{1}{2}\left\langle Q\left(x-x_{0}\right), x-x_{0}\right\rangle \text { for all } x \in S,
$$

where $p=\nabla c\left(x_{0}-y_{1}\right)-\nabla c\left(x_{0}-y_{0}\right)$ and $Q=D^{2} c\left(\zeta-y_{1}\right)-D^{2} c\left(\zeta-y_{0}\right)$. Since our cost is of the form $c(z)=\frac{1}{2}|z|^{2}+\frac{1}{\lambda} E(z)$, an estimate for $Q$ is

$$
\|Q\|=\frac{1}{\lambda}\left\|D^{2} E\left(\zeta-y_{1}\right)-D^{2} E\left(\zeta-y_{0}\right)\right\| \leq \frac{C}{\lambda}\left|y_{0}-y_{1}\right| .
$$

Therefore, $S$ is the intersection of almost-halfplanes,

$$
0 \geq\left\langle p, x-x_{0}\right\rangle+\frac{1}{2}\left\langle Q\left(x-x_{0}\right), x-x_{0}\right\rangle
$$

The section $S$ would be convex if $Q \equiv 0$ for all $x_{0} \in \partial S$. In our case, these supporting functions are planes plus a small quadratic perturbation; we would like to measure the distance of how far is $S$ from $[S]$. If $S$ is not convex, for $x_{0} \in \partial S$, we can find $x \in S$ such that $x$ is outside the half plane

$$
H_{x_{0}}: 0 \geq\left\langle\frac{p}{|p|}, x-x_{0}\right\rangle
$$

Then if we let $d=\operatorname{dist}\left(x, H_{x_{0}}\right)$ then $d=\left\langle\frac{p}{|p|}, x-x_{0}\right\rangle>0$. We also have

$$
\left\langle p, x-x_{0}\right\rangle \leq-\frac{1}{2}\left\langle Q\left(x-x_{0}\right), x-x_{0}\right\rangle .
$$

Hence by using the estimate of $Q$,

$$
d \leq \frac{C}{\lambda} \frac{\left|y_{0}-y_{1}\right|}{|p|}\left|x-x_{0}\right|^{2} \leq \frac{C}{\lambda} \frac{\left|y_{0}-y_{1}\right|}{|p|} .
$$

On the other hand, by its construction

$$
p=\nabla c\left(x_{0}-y_{1}\right)-\nabla c\left(x_{0}-y_{0}\right)=y_{0}-y_{1}+\frac{1}{\lambda}\left[\nabla E\left(x_{0}-y_{1}\right)-\nabla E\left(x_{0}-y_{0}\right)\right] .
$$

Thus, $\left(1-\frac{C}{\lambda}\right)\left|y_{0}-y_{1}\right| \leq|p| \leq\left(1+\frac{C}{\lambda}\right)\left|y_{0}-y_{1}\right|$ and by combining with the above estimate we obtain $d \leq \frac{C}{\lambda}$. We have in fact shown for each $x_{0} \in \partial S$,

$$
\operatorname{dist}\left(x, H_{x_{0}}\right) \leq \frac{C}{\lambda} \text { for all } x \in S
$$

It follows immediately from this that

$$
\operatorname{dist}\left(x_{0}, \partial[S]\right) \leq \frac{C}{\lambda}
$$

### 3.2 On the local $c$-subdifferential

In Lemma 3.1 we have given a characterization of the global $c$-subdifferentials of $u$. It is well known that this concept may be different from the local $c$-subdifferential, and this is one of the main difficulties in the study of the regularity for optimal transportation. More precisely, the restriction of an Alexandroff solution of the transport problem to a smaller domain may not be an Alexandroff solution any longer because extra mass may appear (globalflocal). We give here the first attempt to understand the behavior of the subdifferential under restriction to a smaller domain. In particular, we show that the characterization of Lemma 3.1 on "almost"-convexity is still true for (local) $c$-subgradients.

We have defined in Section 2 the (global) c-subdifferential of $u$ at $\tilde{x}$ by

$$
\partial_{c} u(\tilde{x}):=\left\{\tilde{y} \in \mathbb{R}^{n}: u(x) \geq u(\tilde{x})-c(x-\tilde{y})+c(\tilde{x}-\tilde{y}) \quad \text { for all } x \in \Omega\right\}
$$

and the (local) c-subdifferential of $u$ at $\tilde{x}$ as

$$
\partial_{c}^{\mathrm{loc}} u(\tilde{x}):=\left\{p \in \mathbb{R}^{n}: u(x) \geq u(\tilde{x})-c(x-p)+c(\tilde{x}-p) \quad \text { in some neighborhood of } \tilde{x}\right\} .
$$

More generally for an open set $V \subset \Omega$ and $\tilde{x} \in V$, we write

$$
\partial_{c}(u, V)(\tilde{x}):=\left\{q \in \mathbb{R}^{n}: u(x) \geq u(\tilde{x})-c(x-q)+c(\tilde{x}-q) \quad \text { for all } x \in V\right\} .
$$

It is clear that $\partial_{c}^{\text {loc }} u(\tilde{x}) \supset \partial_{c} u(\tilde{x})$, but in general they are not equal.
We will also need to work with the following set of subgradients and $c$-subgradients. For $\tilde{x} \in \Omega$, define

$$
\begin{aligned}
& \partial^{-} u(\tilde{x}):=\left\{p \in \mathbb{R}^{n}: u(x) \geq u(\tilde{x})+\langle p, x-\tilde{x}\rangle+o(|x-\tilde{x}|) \text { for } x \text { near } \tilde{x}\right\}, \\
& \partial_{c}^{-} u(\tilde{x}):=\left\{\tilde{y} \in \mathbb{R}^{n}: u(x) \geq u(\tilde{x})-c(x-\tilde{y})+c(\tilde{x}-\tilde{y})+o(|x-\tilde{x}|) \text { for } x \text { near } \tilde{x}\right\} .
\end{aligned}
$$

It follows that

$$
\begin{equation*}
\partial^{-} u(\tilde{x})=-\nabla c\left(\tilde{x}-\partial_{c}^{-} u(\tilde{x})\right) . \tag{3.9}
\end{equation*}
$$

The above mappings can be extended to boundary points as follows. Let $\tilde{x} \in \partial \Omega$, we denote $\partial^{-} u(\tilde{x}):=\left\{p \in \mathbb{R}^{n}: p=\lim _{k \rightarrow \infty} p_{k}\right\}$, where $p_{k} \in \partial^{-} u\left(x_{k}\right)$ and $\left\{x_{k}\right\}$ is a sequence of interior points of $\Omega$ converging to $\tilde{x}$, and let $\partial_{c}^{-} u(\tilde{x})$ be given by (3.9).

First we present two well known lemmas that relate the $c$-subdifferentials of two functions:
Lemma 3.3. Let $u, v \in \mathcal{C}(\bar{\Omega})$ and $E \subset \Omega$ be an open set such that

$$
\begin{cases}u=v & \text { on } \partial E, \\ u \geq v & \text { in } E .\end{cases}
$$

Then $\partial_{c}(u, E)(E) \subset \partial_{c}(v, E)(E)$. Consequently, $\partial_{c} u(E) \subset \partial_{c}^{l o c} v(E)$.
Proof. This result can be found in [28, Lemma 5.1] but we include the proof here for convenience. Let $p \in \partial_{c}(u, E)(E)$. Then there exists $x_{0} \in E$ such that $u(x) \geq u\left(x_{0}\right)-c(x-p)+$ $c\left(x_{0}-p\right)$ for all $x \in E$. Let

$$
a:=\sup _{x \in \bar{E}}\left\{u\left(x_{0}\right)-c(x-p)+c\left(x_{0}-p\right)-v(x)\right\} .
$$

Since $u\left(x_{0}\right) \geq v\left(x_{0}\right)$, we get $a \geq 0$. Also there is $x_{1} \in \bar{E}$ such that $a=u\left(x_{0}\right)-c\left(x_{1}-p\right)+$ $c\left(x_{0}-p\right)-v\left(x_{1}\right)$. Thus $v(x) \geq v\left(x_{1}\right)-c(x-p)+c\left(x_{1}-p\right)$ for all $x \in E$. Then two things can happen:

If $a>0$, then since $u\left(x_{1}\right) \geq a+v\left(x_{1}\right)$ we must have $x_{1} \notin \partial E$ and hence, $p \in \partial_{c}(v, E)\left(x_{1}\right) \subset$ $\partial_{c}(v, E)(E)$. On the other hand if $a=0$, then we already have that $v(x) \geq v\left(x_{0}\right)-c(x-$ $p)+c\left(x_{0}-p\right)$ in $E$ which in turn gives $p \in \partial_{c}(v, E)\left(x_{0}\right) \subset \partial_{c}(v, E)(E)$.

For the rest of the paper, $\mathcal{N}_{\nu}(E)$ will denote the $\nu$-neighborhood of the set $E$, i.e.,

$$
\mathcal{N}_{\nu}(E):=\left\{x \in \mathbb{R}^{n}: \operatorname{dist}(x, E) \leq \nu\right\} .
$$

Lemma 3.4. Let $u, v \in \mathcal{C}\left(\overline{B_{R}(\bar{x})}\right)$ and suppose that $u$ is $c_{\lambda}$-convex and $v$ is $c_{q}$-convex. There exists a universal constant $C>0$ such that if $\|u-v\|_{L^{\infty}\left(B_{R}(\bar{x})\right)} \leq \epsilon$, then

$$
\begin{equation*}
\partial_{c} u\left(B_{R-\frac{\epsilon}{\delta}}(\bar{x})\right) \subset \mathcal{N}_{\left(4 \delta+\frac{C}{\lambda}\right)}\left\{\partial_{c_{q}} v\left(B_{R}(\bar{x})\right)\right\} \quad \text { for all } \quad \delta>\frac{\epsilon}{R} . \tag{3.10}
\end{equation*}
$$

Consequently,

$$
\begin{equation*}
\partial_{c}^{l o c} u\left(B_{R-\frac{\epsilon}{\delta}}(\bar{x})\right) \subset \mathcal{N}_{\left(4 \delta+\frac{C}{\lambda}\right)}\left\{\left[\partial_{c_{q}} v\left(B_{R}(\bar{x})\right)\right]\right\}, \tag{3.11}
\end{equation*}
$$

where $\left[\partial_{c_{q}} v\left(B_{R}(\bar{x})\right)\right]$ is the convex hull of $\partial_{c_{q}} v\left(B_{R}(\bar{x})\right)$.
Proof. We note that a version of this result for quadratic cost appears in [11] (Lemma 2 of part 4). Let $p \in \partial_{c} u\left(B_{R-\frac{\epsilon}{\delta}}(\bar{x})\right)$, then there exists $x_{0} \in B_{R-\frac{\epsilon}{\delta}}(\bar{x})$ such that

$$
u(x) \geq u\left(x_{0}\right)-c(x-p)+c\left(x_{0}-p\right) \quad \text { for all } \quad x \in B_{R}:=B_{R}(\bar{x}) .
$$

Consider $v^{*}(x):=v(x)+\epsilon+\frac{2 \delta}{R}\left(|x-\bar{x}|^{2}-R^{2}\right)$, then $\left.v^{*}\right|_{\partial B_{R}} \geq u$ and $\left.v^{*}\right|_{B_{R-\frac{\epsilon}{\delta}}(\bar{x})} \leq u$. Define

$$
a:=\sup _{x \in \bar{B}_{R}}\left\{u\left(x_{0}\right)-c(x-p)+c\left(x_{0}-p\right)-v^{*}(x)\right\} .
$$

Since $u\left(x_{0}\right) \geq v^{*}\left(x_{0}\right)$, we get $a \geq 0$. On the other hand, there exists $x_{1} \in \bar{B}_{R}$ such that $a=u\left(x_{0}\right)-c\left(x_{1}-p\right)+c\left(x_{0}-p\right)-v^{*}\left(x_{1}\right)$, and thus $v^{*}(x) \geq v^{*}\left(x_{1}\right)-c(x-p)+c\left(x_{1}-p\right)$ for all $x \in B_{R}$. Moreover, $u\left(x_{1}\right) \geq u\left(x_{0}\right)-c\left(x_{1}-p\right)+c\left(x_{0}-p\right)=v^{*}\left(x_{1}\right)+a$. Therefore, if $a>0$, then $x_{1}$ does not belong to the boundary of $B_{R}$, and we conclude $p \in \partial_{c} v^{*}\left(x_{1}\right) \subset \partial_{c} v^{*}\left(B_{R}\right)$. If, on the contrary, $a=0$, then $v^{*}(x) \geq v^{*}\left(x_{0}\right)-c(x-p)+c\left(x_{0}-p\right)$ for all $x \in B_{R}$, so automatically $p \in \partial_{c} v^{*}\left(x_{0}\right) \subset \partial_{c} v^{*}\left(B_{R}\right)$. Thus we have shown that $\partial_{c} u\left(B_{R-\frac{\epsilon}{\delta}}(\bar{x})\right) \subset \partial_{c} v^{*}\left(B_{R}\right)$.

Next suppose $y^{*} \in \partial_{c} v^{*}\left(x_{0}\right), x_{0} \in B_{R}$. Then by (3.9), $y^{*}=x_{0}-\nabla c^{*}\left(x_{0}-p^{*}\right)$ for some $p^{*} \in$ $\partial^{-} v^{*}\left(x_{0}\right)$. Define $p:=p^{*}-\frac{4 \delta}{R}\left(x_{0}-\bar{x}\right)$ and we claim that $x_{0}+p \in \partial \bar{v}\left(x_{0}\right)$, where $\bar{v}:=v+\frac{1}{2}|x|^{2}$ is a convex function and $\partial \bar{v}\left(x_{0}\right):=\left\{q: \bar{v}(x) \geq \bar{v}\left(x_{0}\right)+\left\langle q, x-x_{0}\right\rangle \forall x \in B_{R}\right\}$. To see this, let us write $p^{*}=\sum_{i=1}^{N} t_{i} p_{i}^{*}$ where $p_{i}^{*}$ are extremal points of $\partial^{-} v^{*}\left(x_{0}\right)$ and $0 \leq t_{i} \leq 1$ satisfying $\sum_{i=1}^{N} t_{i}=1$. For each $i$, by a result in [7] we can find a sequence $\left\{x_{k}\right\}$ of differentiable points of $v$ such that $x_{k} \rightarrow x_{0}$ and $p_{i}^{*}=\lim _{k \rightarrow \infty} \nabla v^{*}\left(x_{k}\right)=\lim _{k \rightarrow \infty} \nabla v\left(x_{k}\right)+\frac{4 \delta}{R}\left(x_{0}-\bar{x}\right)$. This yields $x_{0}+p_{i}^{*}-\frac{4 \delta}{R}\left(x_{0}-\bar{x}\right)=\lim _{k \rightarrow \infty} \nabla \bar{v}\left(x_{k}\right) \in \partial \bar{v}\left(x_{0}\right)$ for each $i$. We then obtain the claim since $x_{0}+p=\sum_{i=1}^{N} t_{i}\left[x_{0}+p_{i}^{*}-\frac{4 \delta}{R}\left(x_{0}-\bar{x}\right)\right]$ and $\partial \bar{v}\left(x_{0}\right)$ is a convex set. Observe that $y^{*}=x_{0}+p^{*}+\frac{1}{\lambda} \nabla E\left(x_{0}-y^{*}\right)$ by the definitions of $p^{*}$ and $c_{\lambda}$. The relation (3.10) now follows from the claim, the fact $\partial \bar{v}\left(x_{0}\right)=\partial_{c_{q}} v\left(x_{0}\right)$ and

$$
\left|y^{*}-\left[x_{0}+p\right]\right|=\left|x_{0}+p^{*}+\frac{1}{\lambda} \nabla E\left(x_{0}-y^{*}\right)-\left[x_{0}+p^{*}-\frac{4 \delta}{R}\left(x_{0}-\bar{x}\right)\right]\right| \leq 4 \delta+\frac{C}{\lambda} .
$$

Too see (3.11), let $y \in \partial_{c}^{\operatorname{loc}} u(x)$ for $x \in B_{R-\frac{\epsilon}{\delta}}(\bar{x})$. Then there exists $p \in \partial^{-} u(x)$ such that $y=x-\nabla c^{*}(-p)$. We write $p=\sum_{i=1}^{N} t_{i} p_{i}$ as a finite convex combination of extremal points $p_{i}$ of $\partial^{-} u\left(x_{0}\right)$. Since $u$ is $c$-convex, $y_{i}:=x-\nabla c^{*}\left(-p_{i}\right) \in \partial_{c} u(x)$ for all $i$ which implies that

$$
\sum_{i=1}^{N} t_{i} y_{i} \in \mathcal{N}_{\left(4 \delta+\frac{C}{\lambda}\right)}\left\{\left[\partial_{c_{q}} v\left(B_{R}(\bar{x})\right)\right]\right\}
$$

by (3.10). This yields (3.11) because

$$
\left|y-\sum_{i=1}^{N} t_{i} y_{i}\right|=\left|x+p+\frac{1}{\lambda} \nabla E(x-y)-\sum_{i=1}^{N} t_{i}\left[x+p_{i}+\frac{1}{\lambda} \nabla E\left(x-y_{i}\right)\right]\right| \leq \frac{C}{\lambda} .
$$

Next, we present a characterization of subgradients. Although the so called "double mountain above sliding mountain" lemma is not true anymore without assuming the ( $A .3 w$ ) condition, we still have a useful characterization using the asymptotic behavior of the cost $c:=c_{\lambda}$, that generalizes Lemma 3.1.

Proposition 3.5. Suppose that $u \in \mathcal{C}(\bar{\Omega})$ is a c-convex function in $\Omega$. Then for $\tilde{x} \in \bar{\Omega}$, $\tilde{p} \in \partial^{-} u(\tilde{x})$, we have

$$
u(x)+\frac{1}{2}|x|^{2} \geq u(\tilde{x})+\frac{1}{2}|\tilde{x}|^{2}+\langle\tilde{p}+\tilde{x},-\tilde{x}\rangle+\frac{1}{\lambda} O\left(|x-\tilde{x}|^{2}\right) \quad \text { for all } x \in \Omega .
$$

Consequently if $\tilde{y} \in \partial_{c}^{-} u(\tilde{x})$ with $\tilde{x} \in \bar{\Omega}$, then

$$
u(x) \geq u(\tilde{x})-c(x-\tilde{y})+c(\tilde{x}-\tilde{y})+\frac{1}{\lambda} O\left(|x-\tilde{x}|^{2}\right) \quad \text { for all } x \in \Omega .
$$

Proof. First assume that $\tilde{x} \in \Omega$. By construction, $\partial^{-} u(\tilde{x})$ is a closed convex set and hence we can write $\tilde{p}$ as a convex combination of finite extremal points, that is, $\tilde{p}=\sum_{i=1}^{N} t_{i} q_{i}$, where $t_{i} \geq 0, \sum_{i=1}^{N} t_{i}=1$ and $q_{i}$ are extremal points in $\partial^{-} u(\tilde{x})$.

For each $i$, let $\bar{y}_{i} \in \mathbb{R}^{n}$ be such that $q_{i}=-\nabla c\left(\tilde{x}-\bar{y}_{i}\right)$. Since $q_{i}=\lim _{m \rightarrow \infty} \nabla u\left(z_{m}\right)$ for a sequence $z_{m} \rightarrow \tilde{x}$ which are points of differentiability of $u$, it follows that $\bar{y}_{i} \in \partial_{c} u(\tilde{x})$. Then, for each $\bar{y}_{i}$, using the fact $c(z)=\frac{1}{2}|z|^{2}+\frac{1}{\lambda} E(z)$ given by (2.6), we have that

$$
u(x)+\frac{1}{2}|x|^{2} \geq u(\tilde{x})+\frac{1}{2}|\tilde{x}|^{2}+\left\langle\bar{y}_{i}, x-\tilde{x}\right\rangle+\frac{1}{\lambda}\left[E\left(\tilde{x}-\bar{y}_{i}\right)-E\left(x-\bar{y}_{i}\right)\right] \quad \text { for all } x \in \Omega,
$$

and consequently, as $\bar{y}_{i}=q_{i}+\tilde{x}+\frac{1}{\lambda} \nabla E\left(\tilde{x}-\bar{y}_{i}\right)$ by our choice of $\bar{y}_{i}$,

$$
\begin{aligned}
& u(x)+\frac{1}{2}|x|^{2} \\
& \geq u(\tilde{x})+\frac{1}{2}|\tilde{x}|^{2}+\left\langle q_{i}+\tilde{x}, x-\tilde{x}\right\rangle-\frac{1}{\lambda}\left[E\left(x-\bar{y}_{i}\right)-E\left(\tilde{x}-\bar{y}_{i}\right)-\left\langle\nabla E\left(\tilde{x}-\bar{y}_{i}\right), x-\tilde{x}\right\rangle\right] \\
& =u(\tilde{x})+\frac{1}{2}|\tilde{x}|^{2}+\left\langle q_{i}+\tilde{x}, x-\tilde{x}\right\rangle+\frac{1}{\lambda} O\left(|x-\tilde{x}|^{2}\right) \quad \text { for all } x \in \Omega
\end{aligned}
$$

Therefore if we take the convex combination with respect to $i$, we obtain

$$
\begin{equation*}
u(x)+\frac{1}{2}|x|^{2} \geq u(\tilde{x})+\frac{1}{2}|\tilde{x}|^{2}+\langle\tilde{p}+\tilde{x}, x-\tilde{x}\rangle+\frac{1}{\lambda} O\left(|x-\tilde{x}|^{2}\right) \quad \text { for all } x \in \Omega . \tag{3.12}
\end{equation*}
$$

Now assume that $\tilde{x} \in \partial \Omega$. Then by the definition of subdifferential at boundary points, we have $\tilde{p}=\lim _{k \rightarrow \infty} p_{k}$ where $p_{k} \in \partial^{-} u\left(x_{k}\right)$ and $\left\{x_{k}\right\}$ is a sequence of interior points of $\Omega$ converging to $\tilde{x}$. The above proof shows that (3.12) holds if $\tilde{x}$ and $\tilde{p}$ are replaced by $x_{k}$ and $p_{k}$, and therefore by letting $k$ tend to infinity we see that (3.12) is still true when $\tilde{x} \in \partial \Omega$. Thus the first statement of the lemma follows.

To prove the second statement, let $\tilde{p}:=-\nabla c(\tilde{x}-\tilde{y})$. Then we obtain (3.12) since $\tilde{p} \in$ $\partial^{-} u(\tilde{x})$ by the assumption $\tilde{y} \in \partial_{c}^{-} u(\tilde{x})$. It follows from (3.12) and by a similar argument to the one used above that

$$
u(x) \geq u(\tilde{x})-c(x-\tilde{y})+c(\tilde{x}-\tilde{y})+\frac{1}{\lambda} O\left(|x-\tilde{x}|^{2}\right) \quad \text { for all } x \in \Omega .
$$

## 4 Initial step

This is the first ingredient in the proof of Theorem 1.1. Let $u \in \mathcal{C}(\bar{\Omega})$ be a solution of the optimal transport problem between $\Omega$ and $\Omega^{\prime}$, i.e.,

$$
\left\{\begin{align*}
M_{\lambda}[u] & =f \text { in } \Omega,  \tag{4.1}\\
\partial_{c_{\lambda}} u(\Omega) & =\Omega^{\prime},
\end{align*}\right.
$$

where $M_{\lambda}$ is defined in (2.10) and satisfies (2.11), with the hypothesis stated in Theorem 1.1.
The main idea is to compare the solution $u$ of the optimal transport equation (2.9) to $w$, where $w$ is the solution of the optimal transport problem between $\Omega$ and $\Omega^{\prime}$ with quadratic $\operatorname{cost} c_{q}(z)=\frac{1}{2}|z|^{2}$, and $\Omega^{\prime}$ convex. Note that $w$ is unique up to an additive constant, strictly convex and smooth (see the article by Caffarelli [7], further regularity is contained in [8], [10]). The equation satisfied by $w$ is the following

$$
\left\{\begin{align*}
g(x+\nabla w) \operatorname{det}\left(I+D^{2} w\right) & =f \text { in } \Omega,  \tag{4.2}\\
\partial_{c_{q}} w(\Omega) & =\Omega^{\prime} .
\end{align*}\right.
$$

### 4.1 An approximation result

We prove a comparison result between the solution $u$ of (4.1) to the solution $w$ of (4.2):
Proposition 4.1. For any $\epsilon_{0}>0$, there exists $\lambda_{0}>0$ depending only on $\epsilon_{0}, n, \Lambda_{1}, \Lambda_{2}$, $\Omega$ and $\Omega^{\prime}$ such that if $\lambda \geq \lambda_{0}$, $f$ and $g$ satisfy (1.3), $w \in \mathcal{C}(\bar{\Omega})$ is a solution of (4.2) and $u \in \mathcal{C}(\bar{\Omega})$ is a solution of the transport equation (4.1) with $u\left(x_{0}\right)=w\left(x_{0}\right)$ for some $x_{0} \in \Omega$, then we have

$$
\|u-w\|_{L^{\infty}(\Omega)} \leq \epsilon_{0} .
$$

The proof of this proposition is a simple consequence of the following strengthened result in which the constant $\lambda_{0}$ is independent of the domains $\Omega$ and $\Omega^{\prime}$, as long as they lie between two universal balls.

Proposition 4.2. Let $\Omega$ and $\Omega^{\prime}$ be two closed subsets of $\mathbb{R}^{n}$ satisfying $B_{1}(0) \subset \Omega \subset \bar{B}_{n}(0)$ and $B_{1}(0) \subset \Omega^{\prime} \subset \bar{B}_{n}(0)$. Let $f$ and $g$ satisfy (1.3).

For any $\epsilon_{0}>0$, there exists $\lambda_{0}>0$ depending only on $\epsilon_{0}, n, \Lambda_{1}$ and $\Lambda_{2}$ such that if $\lambda \geq \lambda_{0}, w \in \mathcal{C}(\bar{\Omega})$ is a solution of (4.2) and $u \in \mathcal{C}(\bar{\Omega})$ is a solution of the transport equation (4.1) with $u\left(x_{0}\right)=w\left(x_{0}\right)$ for some $x_{0} \in \Omega$, then we have

$$
\begin{equation*}
\|u-w\|_{L^{\infty}\left(B_{1}(0)\right)} \leq \epsilon_{0} . \tag{4.3}
\end{equation*}
$$

Proof. By contradiction, assume that there exist some $\epsilon_{0},\left\{\lambda_{k}\right\} \subset \mathbb{R},\left\{\Omega_{k}\right\},\left\{\Omega_{k}^{\prime}\right\},\left\{f_{k}\right\}$, $\left\{g_{k}\right\},\left\{w_{k}\right\},\left\{u_{k}\right\}$ and $\left\{x_{k}\right\}$ with $x_{k} \in \Omega_{k}$ such that $\lambda_{k} \rightarrow \infty, B_{1}(0) \subset \Omega_{k} \subset \bar{B}_{n}(0)$ closed, $B_{1}(0) \subset \Omega_{k}^{\prime} \subset \bar{B}_{n}(0)$ closed, $f_{k}$ and $g_{k}$ satisfy (1.3) on its domains $\Omega_{k}$ and $\Omega_{k}^{\prime}, w_{k} \in \mathcal{C}\left(\bar{\Omega}_{k}\right)$ is a solution of

$$
\left\{\begin{aligned}
g_{k}\left(x+\nabla w_{k}\right) \operatorname{det}\left(I+D^{2} w_{k}\right) & =f_{k} \quad \text { in } \Omega_{k}, \\
\partial_{c_{q}} w_{k}\left(\Omega_{k}\right) & =\Omega_{k}^{\prime},
\end{aligned}\right.
$$

and $u_{k} \in \mathcal{C}\left(\bar{\Omega}_{k}\right)$ is a solution of

$$
\left\{\begin{aligned}
g_{k}\left(x-\nabla c^{*}\left(-\nabla u_{k}\right)\right) \operatorname{det}\left(I+D^{2} c_{k}^{*}\left(-\nabla u_{k}\right) D^{2} u_{k}\right) & =f_{k} \text { in } \Omega_{k}, \\
\partial_{c_{k}} u_{k}\left(\Omega_{k}\right) & =\Omega_{k}^{\prime},
\end{aligned}\right.
$$

where $c_{k}=c_{\lambda_{k}}, u_{k}\left(x_{k}\right)=w_{k}\left(x_{k}\right)$ and

$$
\begin{equation*}
\left\|u_{k}-w_{k}\right\|_{L^{\infty}\left(B_{1}(0)\right)} \geq \epsilon_{0} \quad \text { for all } \quad k \geq 1 . \tag{4.4}
\end{equation*}
$$

By selecting a subsequence, we know from [36, Theorem 1.8.4] that $\left\{\Omega_{k}\right\}$ converges to $\Omega$ and $\left\{\Omega_{k}^{\prime}\right\}$ converges to $\Omega^{\prime}$ in the Hausdorff distance topology. Moreover, $\Omega$ and $\Omega^{\prime}$ are closed sets satisfying $B_{1}(0) \subset \Omega \subset \bar{B}_{n}(0)$ and $B_{1}(0) \subset \Omega^{\prime} \subset \bar{B}_{n}(0)$. It also follows that $x_{k} \rightarrow \bar{x}$ for some $\bar{x} \in \Omega$.

Note that $\left\|u_{k}-w_{k}\right\|_{L^{\infty}\left(B_{1}(0)\right)}=\left\|\left[u_{k}-u_{k}\left(x_{k}\right)\right]-\left[w_{k}-w_{k}\left(x_{k}\right)\right]\right\|_{L^{\infty}\left(B_{1}(0)\right)}$. By replacing $u_{k}$ and $w_{k}$ by $u_{k}-u_{k}\left(x_{k}\right)$ and $w_{k}-w_{k}\left(x_{k}\right)$ respectively, we can assume in addition that $u_{k}\left(x_{k}\right)=w_{k}\left(x_{k}\right)=0$. Next for convenience, we extend $u_{k}$ as a $c_{k}$-convex function on $\bar{B}_{n}(0)$. Let us describe precisely this natural extension. Define the Kantorovich functional as in (2.1),

$$
I_{k}(\phi, \psi):=-\int_{\Omega_{k}} f_{k}(x) \phi(x) d x-\int_{\Omega_{k}^{\prime}} g_{k}(y) \psi(y) d y
$$

on the set
$K_{c_{k}}:=\left\{(\phi, \psi) \in \operatorname{Lip}\left(\Omega_{k}\right) \times \operatorname{Lip}\left(\Omega_{k}^{\prime}\right):-\phi(x)-\psi(y) \leq c_{k}(x-y) \quad\right.$ for all $\left.\quad x \in \Omega_{k}, y \in \Omega_{k}^{\prime}\right\}$. Let $v_{k}$ be the $c_{k}$-transform of $u_{k}$, that is,

$$
\begin{equation*}
v_{k}(y):=\sup _{x \in \Omega_{k}}\left\{-c_{k}(x-y)-u_{k}(x)\right\} \quad \text { for } y \in \mathbb{R}^{n} . \tag{4.5}
\end{equation*}
$$

We know that $\left(u_{k}, v_{k}\right)$ is a maximizer of the functional $I_{k}(\phi, \psi)$ over $K_{c_{k}}$, and furthermore

$$
u_{k}(x)=\sup _{y \in \Omega_{k}^{\prime}}\left\{-c_{k}(x-y)-v_{k}(y)\right\} \quad \text { for all } \quad x \in \Omega_{k} .
$$

Now let us extend $u_{k}$ by defining

$$
\begin{equation*}
u_{k}^{*}(x):=\sup _{y \in \bar{B}_{n}(0)}\left\{-c_{k}(x-y)-v_{k}(y)\right\} \quad \text { for } x \in \bar{B}_{n}(0) \supset \Omega_{k} . \tag{4.6}
\end{equation*}
$$

Since $\Omega_{k}^{\prime} \subset \bar{B}_{n}(0)$, it is clear that $u_{k}^{*}(x) \geq u_{k}(x)$ for all $x \in \Omega_{k}$. On the other hand, $\sup _{y \in \bar{B}_{n}(0)}\left\{-c_{k}(x-y)-v_{k}(y)\right\} \leq u_{k}(x)$ for $x \in \Omega_{k}$ by (4.5). Thus, $\left.u_{k}^{*}\right|_{\Omega_{k}} \equiv u_{k}$. In particular, we have

$$
\begin{aligned}
\sup _{x \in \bar{B}_{n}(0)}\left\{-c_{k}(x-y)-u_{k}^{*}(x)\right\} & \geq \sup _{x \in \Omega_{k}}\left\{-c_{k}(x-y)-u_{k}^{*}(x)\right\} \\
& =\sup _{x \in \Omega_{k}}\left\{-c_{k}(x-y)-u_{k}(x)\right\}=v_{k}(y) .
\end{aligned}
$$

Also, for $y \in \bar{B}_{n}(0)$, it follows from the definition of $u_{k}^{*}$ in (4.6) that $u_{k}^{*}(x) \geq-c_{k}(x-y)-v_{k}(y)$ for all $x \in \bar{B}_{n}(0)$ giving $v_{k}(y) \geq \sup _{x \in \bar{B}_{n}(0)}\left\{-c_{k}(x-y)-u_{k}^{*}(x)\right\}$. Hence we obtain

$$
v_{k}(y)=\sup _{x \in \bar{B}_{n}(0)}\left\{-c_{k}(x-y)-u_{k}^{*}(x)\right\} \quad \forall y \in \bar{B}_{n}(0) .
$$

From now on, we identify $u_{k}$ with its extension $u_{k}^{*}$ given by (4.6). Likewise, we also identify the function $w_{k}$ with its extension on $\bar{B}_{n}(0)$. Let us check that $\left\{u_{k}\right\}$ is equicontinuous on $\bar{B}_{n}(0)$ : fix $x_{1}, x_{2} \in \bar{B}_{n}(0)$. By the extension of $u_{k}$, there exists $y_{k} \in \partial_{c_{k}} u_{k}\left(x_{2}\right)$ such that $y_{k} \in \bar{B}_{n}(0)$. Then by definition,

$$
u_{k}(x)-u_{k}(y) \geq-c_{k}\left(x-y_{k}\right)+c_{k}\left(y-y_{k}\right) .
$$

Because $\left\|y_{k}\right\| \leq n$, we get

$$
\left|c_{k}\left(x_{1}-y_{k}\right)-c_{k}\left(x_{2}-y_{k}\right)\right| \leq C\left|x_{1}-x_{2}\right|
$$

for some constant $C$ independent of $k$. Thus $u_{k}\left(x_{1}\right)-u_{k}\left(x_{2}\right) \geq-C\left|x_{1}-x_{2}\right|$. It then follows by interchanging $x_{1}$ and $x_{2}$ that

$$
\left|u_{k}\left(x_{1}\right)-u_{k}\left(x_{2}\right)\right| \leq C\left|x_{1}-x_{2}\right|,
$$

and we have proved that $\left\{u_{k}\right\}$ is equicontinuous on $\bar{B}_{n}(0)$. Moreover $\left\{u_{k}\right\}$ is uniformly bounded, because we have just seen that $\left|u_{k}(x)\right|=\left|u_{k}(x)-u_{k}\left(x_{k}\right)\right| \leq C\left|x-x_{k}\right| \leq C$. Therefore by Arzela-Ascoli theorem, we have that a subsequence, still denoted by $u_{k}$, converges uniformly on $\bar{B}_{n}(0)$ to some $u_{\infty} \in \mathcal{C}\left(\bar{B}_{n}(0)\right)$. Similarly, there exists a subsequence $\left\{w_{k}\right\}$ converging uniformly on $\bar{B}_{n}(0)$ to some $w_{\infty} \in \mathcal{C}\left(\bar{B}_{n}(0)\right)$. These yield $u_{\infty}(\bar{x})=w_{\infty}(\bar{x})$ for some $\bar{x}$ as a consequence.

Notice that $f_{k}$ and $g_{k}$ can be viewed as functions defined on $B_{n}(0)$ by extending them to be zero outside $\Omega_{k}$ and $\Omega_{k}^{\prime}$ respectively. Then by extracting subsequences, we can assume that $f_{k} \stackrel{*}{\rightharpoonup} f$ weakly* in $L^{\infty}\left(B_{n}(0)\right)$ and $g_{k} \xrightarrow{*} g$ weakly* in $L^{\infty}\left(B_{n}(0)\right)$. Let $\mu_{k}:=f_{k}(x) d x$, $\mu:=f(x) d x, \nu_{k}:=g_{k}(y) d y$ and $\nu:=g(y) d y$.

Claim 1: $\bar{B}_{1}(0) \subset \operatorname{supp} \mu \subset \Omega$ and $\bar{B}_{1}(0) \subset \operatorname{supp} \nu \subset \Omega^{\prime}$. Indeed since $f_{k} \stackrel{*}{\rightharpoonup} f$ weakly* in $L^{\infty}\left(B_{n}(0)\right)$, we get $\left\{\mu_{k}\right\}$ converges weakly* to $\mu$ as measures, i.e.,

$$
\lim _{k \rightarrow \infty} \int_{B_{n}(0)} \varphi(x) d \mu_{k}=\int_{B_{n}(0)} \varphi(x) d \mu \quad \text { for all } \quad \varphi \in \mathcal{C}_{c}\left(B_{n}(0)\right) .
$$

Now let $x \in \operatorname{supp} \mu$. Then by following the proof of [1, Proposition 5.1.8] where probability measures are considered, we see that there exists $x_{k} \in \operatorname{supp} \mu_{k}=\Omega_{k}$ such that $x_{k} \rightarrow x$. But then we must have $x \in \Omega$ since $\left\{\Omega_{k}\right\}$ converges to $\Omega$ in the Hausdorff topology. Thus $\operatorname{supp} \mu \subset \Omega$. Next if $x_{0} \in B_{1}(0)$, then $\Lambda_{1} \leq \frac{1}{\left|B\left(x_{0}, r\right)\right|} \int_{B\left(x_{0}, r\right)} f d x \leq \Lambda_{2}$ for all $r>0$ sufficiently small because $\Lambda_{1} \leq \frac{1}{\left|B\left(x_{0}, r\right)\right|} \int_{B\left(x_{0}, r\right)} f_{k} d x \leq \Lambda_{2}$ and

$$
\frac{1}{\left|B\left(x_{0}, r\right)\right|} \int_{B\left(x_{0}, r\right)} f_{k} d x \rightarrow \frac{1}{\left|B\left(x_{0}, r\right)\right|} \int_{B\left(x_{0}, r\right)} f d x \quad \text { as } \quad k \rightarrow \infty .
$$

Hence by letting $r \rightarrow 0^{+}$, we obtain $\Lambda_{1} \leq f\left(x_{0}\right) \leq \Lambda_{2}$ for a.e. $x_{0} \in B_{1}(0)$. This implies $\bar{B}_{1}(0) \subset \operatorname{supp} \mu$ and we have shown that $\bar{B}_{1}(0) \subset \operatorname{supp} \mu \subset \Omega$. Thus Claim 1 is proved since the case for $\nu$ is completely similar.

We would like to show next that $u_{\infty}$ and $w_{\infty}$ are both solutions of (4.2). Let us prove this only for $u_{\infty}$ as the case of $w_{\infty}$ is much simpler and follows in a similar way. Recall that $v_{k}(y)=\sup _{x \in \bar{B}_{n}(0)}\left\{-c_{k}(x-y)-u_{k}(x)\right\}$ for all $y \in \bar{B}_{n}(0)$, and let

$$
v_{\infty}(y):=\sup _{x \in \bar{B}_{n}(0)}\left\{-c_{q}(x-y)-u_{\infty}(x)\right\} \quad \text { for } y \in \bar{B}_{n}(0),
$$

where $c_{q}$ denotes the quadratic cost. Then since $c_{k}(z)=c_{q}(z)+\frac{E_{k}(z)}{\lambda_{k}}$, it is easy to see that

$$
\left\|v_{k}-v_{\infty}\right\|_{L^{\infty}\left(\bar{B}_{n}(0)\right)} \leq\left\|u_{k}-u_{\infty}\right\|_{L^{\infty}\left(\bar{B}_{n}(0)\right)}+\frac{C}{\lambda_{k}}
$$

yielding $v_{k} \rightarrow v_{\infty}$ uniformly in $\bar{B}_{n}(0)$.
Claim 2: We have

$$
\lim _{k \rightarrow \infty} \int_{\Omega_{k}} f_{k} u_{k}=\int_{\Omega} f u_{\infty}, \quad \lim _{k \rightarrow \infty} \int_{\Omega_{k}^{\prime}} g_{k} v_{k}=\int_{\Omega^{\prime}} g v_{\infty} \quad \text { and } \quad \int_{\Omega} f=\int_{\Omega^{\prime}} g .
$$

To see this, observe that $\int_{\Omega} u_{\infty} f=\int_{\operatorname{supp} \mu} u_{\infty} f=\int_{B_{n}(0)} u_{\infty} f$ since supp $\mu \subset \Omega$ by Claim 1 . Therefore,

$$
\begin{aligned}
\int_{\Omega_{k}} f_{k} u_{k}-\int_{\Omega} f u_{\infty} & =\int_{\Omega_{k}} f_{k}\left(u_{k}-u_{\infty}\right)+\left[\int_{\Omega_{k}} f_{k} u_{\infty}-\int_{\Omega} f u_{\infty}\right] \\
& =\int_{\Omega_{k}} f_{k}\left(u_{k}-u_{\infty}\right)+\left[\int_{B_{n}(0)} f_{k} u_{\infty}-\int_{B_{n}(0)} f u_{\infty}\right] \longrightarrow 0
\end{aligned}
$$

giving $\lim _{k \rightarrow \infty} \int_{\Omega_{k}} f_{k} u_{k}=\int_{\Omega} f u_{\infty}$ as desired. Similarly, $\lim _{k \rightarrow \infty} \int_{\Omega_{k}^{\prime}} g_{k} v_{k}=\int_{\Omega^{\prime}} g v_{\infty}$. Also

$$
\int_{\Omega} f=\int_{B_{n}(0)} f=\lim _{k \rightarrow \infty} \int_{B_{n}(0)} f_{k}=\lim _{k \rightarrow \infty} \int_{B_{n}(0)} g_{k}=\int_{B_{n}(0)} g=\int_{\Omega^{\prime}} g
$$

and the claim follows.
We are ready to show that $u_{\infty}$ is a solution of (4.2). Let $I(\phi, \psi)$ be the Kantorovich functional given by (2.1), and define

$$
K_{c_{q}}:=\left\{(\phi, \psi) \in \operatorname{Lip}(\Omega) \times \operatorname{Lip}\left(\Omega^{\prime}\right):-\phi(x)-\psi(y) \leq c_{q}(x-y) \quad \text { for all } \quad x \in \Omega, y \in \Omega^{\prime}\right\} .
$$

Let $\left(\phi_{0}, \psi_{0}\right) \in K_{c_{q}}$ be a maximizer of the functional $I(\phi, \psi)$ over $K_{c_{q}}$ and recall that ( $u_{k}, v_{k}$ ) is a maximizer of the functional $I_{\bar{k}}(\phi, \psi)$ over $K_{c_{k}}$. As done for $\left(u_{k}, v_{k}\right)$, we can extend $\phi_{0}$ and $\psi_{0}$ as Lipschitz functions on $\bar{B}_{n}(0)$ satisfying $\psi_{0}(y)=\sup _{x \in \bar{B}_{n}(0)}\left\{-c_{q}(x-y)-\phi_{0}(x)\right\}$ for all $y \in \bar{B}_{n}(0)$. Using the facts $\left(u_{k}+\frac{C}{\lambda_{k}}, v_{k}\right) \in K_{C_{q}}$ and $\left(\phi_{0}+\frac{C}{\lambda_{k}}, \psi_{0}\right) \in K_{c_{k}}$, we obtain $I_{k}\left(\phi_{0}+\frac{C}{\lambda_{k}}, \psi_{0}\right)-I\left(\phi_{0}, \psi_{0}\right) \leq I_{k}\left(u_{k}, v_{k}\right)-I\left(\phi_{0}, \psi_{0}\right) \leq I_{k}\left(u_{k}, v_{k}\right)-I\left(u_{k}+\frac{C}{\lambda_{k}}, v_{k}\right)$ giving

$$
\begin{aligned}
& -\frac{C}{\lambda_{k}} \int_{\Omega_{k}} f_{k}+\int_{\Omega} f \phi_{0}-\int_{\Omega_{k}} f_{k} \phi_{0}+\int_{\Omega} g \psi_{0}-\int_{\Omega_{k}^{\prime}} g_{k} \psi_{0} \\
& \leq I_{k}\left(u_{k}, v_{k}\right)-I\left(\phi_{0}, \psi_{0}\right) \leq \frac{C}{\lambda_{k}} \int_{\Omega} f+\int_{\Omega} f u_{k}-\int_{\Omega_{k}} f_{k} u_{k}+\int_{\Omega^{\prime}} g v_{k}-\int_{\Omega_{k}^{\prime}} g_{k} v_{k} \text { for all } k .
\end{aligned}
$$

Hence it follows from Claim 2 that $\lim _{k \rightarrow \infty} I_{k}\left(u_{k}, v_{k}\right)=I\left(\phi_{0}, \psi_{0}\right)$. Consequently,

$$
\begin{aligned}
I\left(u_{\infty}, v_{\infty}\right) & =-\int_{\Omega} f(x) u_{\infty}(x) d x-\int_{\Omega^{\prime}} g(y) v_{\infty}(y) d y \\
& =\lim _{k \rightarrow \infty}\left\{-\int_{\Omega_{k}} f_{k}(x) u_{k}(x) d x-\int_{\Omega_{k}^{\prime}} g_{k}(y) v_{k}(y) d y\right\} \\
& =\lim _{k \rightarrow \infty} I_{k}\left(u_{k}, v_{k}\right)=I\left(\phi_{0}, \psi_{0}\right)=\sup _{(\phi, \psi) \in K_{c_{q}}} I(\phi, \psi),
\end{aligned}
$$

where we have again used Claim 2 to obtain the second equality. Since $\int_{\Omega} f=\int_{\Omega^{\prime}} g$, we infer that $u_{\infty}$ is a solution of (4.2) as desired. Similarly, $w_{\infty}$ is also a solution of (4.2). Thus $u_{\infty} \equiv w_{\infty}$ on $\operatorname{supp} \mu$ as $u_{\infty}(\bar{x})=w_{\infty}(\bar{x})$ and (4.2) has a unique solution modulo additive constants on the support of $\mu=f(x) d x$. This together with the fact $B_{1}(0) \subset \operatorname{supp} \mu$ yields $\lim _{k \rightarrow \infty}\left\|u_{k}-w_{k}\right\|_{L^{\infty}\left(B_{1}(0)\right)}=\left\|u_{\infty}-w_{\infty}\right\|_{L^{\infty}\left(B_{1}(0)\right)}=0$, a contradiction to (4.4).
Remark 4.3. If we assume in addition that $\Omega$ is convex in Proposition 4.2, then the conclusion (4.3) can be strengthened by $\|u-w\|_{L^{\infty}(\Omega)} \leq \epsilon_{0}$. The reason is that in this case we actually have supp $\mu=\Omega$ instead of supp $\mu \subset \Omega$. Indeed, let $x \in \operatorname{Int}(\Omega)$ and $B(x, r) \subset \Omega$. Then $B(x, r) \subset B_{\frac{r}{4}}\left(\Omega_{k}\right):=\cup_{y \in \Omega_{k}} B\left(y, \frac{r}{4}\right)$ for all $k$ sufficiently large because $\Omega \subset B_{\frac{r}{4}}\left(\Omega_{k}\right)$ from the definition of the Hausdorff distance. Since $\Omega_{k}$ are convex, it is clear that $\operatorname{dist}\left(x, \partial \Omega_{k}\right)=$ $\operatorname{dist}\left(x, \partial B_{\frac{r}{4}}\left(\Omega_{k}\right)\right)-\frac{r}{4} \geq \operatorname{dist}(x, \partial \Omega)-\frac{r}{4} \geq r-\frac{r}{4}=\frac{3 r}{4}$ yielding in particular $B\left(x, \frac{r}{2}\right) \subset \Omega_{k}$ for all $k$ sufficiently large. Consequently,

$$
\mu(B(x, r))=\lim _{k \rightarrow \infty} \mu_{k}(B(x, r)) \geq \Lambda_{1} \lim _{k \rightarrow \infty}\left|B(x, r) \cap \Omega_{k}\right|=\Lambda_{1}|B(x, r)|
$$

implying $x \in \operatorname{supp} \mu$, i.e., $\operatorname{Int}(\Omega) \subset \operatorname{supp} \mu$. This and the Claim 1 above give supp $\mu=\Omega$.

### 4.2 Some initial estimates

With the information given by Proposition 4.1, we can start localizing the problem. Indeed, we will show that points far from the boundary of $\Omega$ are mapped (uniformly) to points far from the boundary of $\Omega^{\prime}$. We remind the reader that we are using the notation

$$
\Omega_{\delta_{0}}=\left\{x \in \Omega: \operatorname{dist}(x, \partial \Omega)>\delta_{0}\right\} .
$$

Proposition 4.4. Given any $\delta_{0}>0$, there exists $\lambda_{0}>0$ depending only on $\delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \Omega$ and $\Omega^{\prime}$ such that for all $\lambda \geq \lambda_{0}$ and $u \in \mathcal{C}(\bar{\Omega})$ is a solution of (4.1), we have

$$
\begin{equation*}
\operatorname{dist}\left(\partial_{c} u\left(\Omega_{\delta_{0}}\right), \partial \Omega^{\prime}\right) \geq C \delta_{0}^{\gamma}-\frac{C}{\lambda} \tag{4.7}
\end{equation*}
$$

for some universal constants $C, \gamma>0$. Consequently,

$$
\begin{equation*}
\operatorname{dist}\left(\partial_{c} u\left(\Omega_{\delta_{0}}\right), \partial_{c} u(\partial \Omega)\right) \geq C \delta_{0}^{\gamma}-\frac{C}{\lambda} . \tag{4.8}
\end{equation*}
$$

Proof. For any $\epsilon>0$, Proposition 4.1 gives that $\left\|u-\left[w+u\left(x_{0}\right)-w\left(x_{0}\right)\right]\right\|_{L^{\infty}(\Omega)}<\epsilon$ for all $\lambda \geq \lambda_{0}=\lambda_{0}\left(\epsilon, n, \Lambda_{1}, \Lambda_{2}, \Omega, \Omega^{\prime}\right)$, where $x_{0} \in \Omega$ and $w \in \mathcal{C}(\bar{\Omega})$ is a solution of (4.2). By working with $\tilde{w}:=w+u\left(x_{0}\right)-w\left(x_{0}\right)$ if necessary, we can assume that $\|u-w\|_{L^{\infty}(\Omega)}<\epsilon$. Let $\phi$ be the smooth convex function satisfying

$$
\left\{\begin{array}{rlll}
\operatorname{det} D^{2} \phi & =1 & \text { in } \quad \Omega_{\frac{\delta_{0}}{2}}, \\
\phi= & 0 & \text { on } & \partial \Omega_{\frac{\delta_{0}}{2}} .
\end{array}\right.
$$

It is known that $-C_{1}\left(n, \Omega_{\frac{\delta_{0}}{2}}\right) \leq \min _{\Omega_{\delta_{0}}} \phi \leq-C_{2}\left(n, \Omega_{\frac{\delta_{0}}{2}}\right.$ (see [26, Proposition 3.2.3]). Hence by using [26, Lemma 5.1.6] we obtain

$$
\begin{equation*}
\Omega_{\delta_{0}} \subset\left\{x \in \Omega_{\frac{\delta_{0}}{2}}: \phi(x) \leq-c \delta_{0}\right\}, \tag{4.9}
\end{equation*}
$$

where $c>0$ depending on $n$ and $\Omega_{\frac{\delta_{0}}{2}}$. In fact, we can choose $c$ to be a constant which depends only on $n$ and $\Omega$ for all $\delta_{0}>0$ sufficiently small. Next, we construct the function

$$
w^{*}(x)=w(x)+\epsilon+2 \epsilon\left(c \delta_{0}\right)^{-1} \phi(x) .
$$

Then we have that $w^{*}(x)>u(x)$ for $x \in \partial \Omega_{\frac{\delta_{0}}{2}}$. We take $V$ to be the set

$$
V:=\left\{x \in \Omega_{\frac{\delta_{0}}{2}}: w^{*}(x)<u(x)\right\} .
$$

It is clear from (4.9) that $\Omega_{\delta_{0}} \subset V \Subset \Omega_{\frac{\delta_{0}}{2}}$. Therefore, Lemma 3.3 assures that

$$
\begin{equation*}
\partial_{c} u\left(\Omega_{\delta_{0}}\right) \subset \partial_{c} u(V) \subset \partial_{c}\left(w^{*}, V\right)(V)=\partial_{c} w^{*}(V) . \tag{4.10}
\end{equation*}
$$

On the other hand, by straightforward computation and using that our cost $c_{\lambda}$ is just a $C / \lambda$-perturbation of the quadratic $\operatorname{cost} c_{q}$, then

$$
\begin{equation*}
\partial_{c} w^{*}(V) \subset \mathcal{N}_{\frac{2 \epsilon}{c \delta_{0}}\|\nabla \phi\|_{L \infty\left(\Omega_{\left.\delta_{0} / 2\right)}\right)}+\frac{C}{\lambda}}\left\{\partial_{c_{q}} w(V)\right\}=\mathcal{N}_{\frac{2 \epsilon}{c \delta_{0}}\|\nabla \phi\|_{L \infty\left(\Omega_{\delta_{0} / 2}\right)}+\frac{C}{\lambda}}\{\nabla \hat{w}(V)\}, \tag{4.11}
\end{equation*}
$$

where $\hat{w}=w+\frac{|x|^{2}}{2}$ and we note that $\|\nabla \phi\|_{L^{\infty}\left(\Omega_{\frac{\delta_{0}}{2}}\right)} \leq C$ (see [25, Theorem 17.21]).
Next, for the solution of the optimal transportation problem $\hat{w}$ with quadratic cost, Caffarelli [8] has proved that

1. $\nabla \hat{w}(\Omega)=\Omega^{\prime}$.
2. $\nabla \hat{w}(\partial \Omega) \subset \partial \Omega^{\prime}$.
3. $|\nabla \hat{w}(x)-\nabla \hat{w}(y)| \geq C|x-y|^{\gamma}$ for all $x, y \in \Omega$, where $\gamma>0$ is a universal constant.

As a consequence, $\nabla \hat{w}(\partial \Omega)=\partial \Omega^{\prime}$ and

$$
\begin{equation*}
\operatorname{dist}\left(\nabla \hat{w}(V), \partial \Omega^{\prime}\right) \geq \operatorname{dist}\left(\nabla \hat{w}\left(\Omega_{\frac{\delta_{0}}{2}}\right), \partial \Omega^{\prime}\right) \geq C \delta_{0}^{\gamma} . \tag{4.12}
\end{equation*}
$$

Combining (4.10), (4.11) and (4.12) we get

$$
\operatorname{dist}\left(\partial_{c} u\left(\Omega_{\delta_{0}}\right), \partial \Omega^{\prime}\right) \geq C \delta_{0}^{\gamma}-\left[\frac{2 \epsilon}{c \delta_{0}}\|\nabla \phi\|_{L^{\infty}\left(\Omega_{\delta_{0} / 2}\right)}+\frac{C}{\lambda}\right] .
$$

Choosing suitable $\epsilon$, we obtain the estimate (4.7) of the lemma. On the other hand, (4.8) follows from (4.7) and the fact that

$$
\partial_{c} u(\partial \Omega) \subset \partial \Omega^{\prime}
$$

which was proven by Figalli-Kim-McCann [17].

Using the above two propositions, one has a very good control of the sections of $u$ : they are very near the sections of $w$, and as a consequence, they stay away from the boundary $\partial \Omega$. Let $x_{0} \in \Omega_{\delta_{0}}$ and $y_{0} \in \partial_{c} u\left(x_{0}\right)$. The $\mu$-section of $u$ at $x_{0}$ is defined in (2.12), or for $\bar{u}(x):=u(x)+c\left(x-y_{0}\right)$ as in (2.13). One can also define the (quadratic) $t$-sections for $w$. Let $\bar{x}$ be the minimum point of the function $\bar{w}(x):=w(x)+\frac{1}{2}\left|x-y_{0}\right|^{2}$. We define

$$
\begin{equation*}
T_{t}(w, \bar{x})=\left\{x \in[\Omega]: w(x) \leq w(\bar{x})-c_{q}(x-\bar{y})+c_{q}(\bar{x}-\bar{y})+t\right\}, \tag{4.13}
\end{equation*}
$$

with

$$
\partial_{c_{q}} w(\bar{x})=\{\bar{y}:=\bar{x}+\nabla w(\bar{x})\} .
$$

By using the function $\bar{w}$ we can write

$$
\begin{equation*}
T_{t}(\bar{w}, \bar{x})=\{x \in[\Omega]: w(x) \leq w(\bar{x})+\langle\nabla \bar{w}(\bar{x}), x-\bar{x}\rangle+t\}, \tag{4.14}
\end{equation*}
$$

that is a convex set.
Proposition 4.5. Given $\delta_{0}>0$, there exists $\mu_{1}=\mu_{1}\left(\delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \Omega, \Omega^{\prime}\right)>0$ satisfying that, for all $h, \epsilon>0$ with $h+\epsilon \leq \mu_{1}$, we can find $\lambda_{0}=\lambda_{0}\left(\epsilon, \delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \Omega, \Omega^{\prime}\right)>0$ such that if $\lambda \geq \lambda_{0}$ and $u$ is a solution of (4.1), then

$$
\begin{equation*}
T_{h-\epsilon}(w, \bar{x}) \subset S_{h}\left(u, x_{0}, y_{0}\right) \subset T_{h+\epsilon}(w, \bar{x}) \subset \Omega_{c \delta_{0}^{\theta}} \tag{4.15}
\end{equation*}
$$

for all $x_{0} \in \Omega_{\delta_{0}}$ and $y_{0} \in \partial_{c} u\left(x_{0}\right)$, where $w$ is a solution of the quadratic transport problem (4.2) and $\bar{x}$ is the minimum point of $\bar{w}:=w+\frac{1}{2}\left|x-y_{0}\right|^{2}$. Here $c$ and $\theta$ are positive constants depending only on $n, \Lambda_{1}, \Lambda_{2}, \Omega$ and $\Omega^{\prime}$.

Proof. Let $x_{0} \in \Omega_{\delta_{0}}$ and $y_{0} \in \partial_{c} u\left(x_{0}\right)$. It follows from Proposition 4.4 that $y_{0} \in \Omega_{\frac{C}{2} \delta_{0}^{\gamma}}^{\prime}$ for all $\lambda \geq \lambda_{0}=\lambda_{0}\left(\delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \Omega, \Omega^{\prime}\right)$. Define $\bar{u}:=u+c\left(x-y_{0}\right)$ and $\hat{w}:=w+\frac{1}{2}|x|^{2}$. Notice that $x_{0}$ is a minimum point for $\bar{u}$ in $\Omega$. By definition,

$$
\begin{aligned}
T_{t}(w, \bar{x}) & =\{x \in \Omega: \bar{w}(x) \leq \bar{w}(\bar{x})+\langle\nabla \bar{w}(\bar{x}), x-\bar{x}\rangle+t\} \\
& =\{x \in \Omega: \hat{w}(x) \leq \hat{w}(\bar{x})+\langle\nabla \hat{w}(\bar{x}), x-\bar{x}\rangle+t\},
\end{aligned}
$$

a convex set. As $\nabla \bar{w}(\bar{x})=0$, we obtain $\nabla \hat{w}(\bar{x})=\nabla w(\bar{x})+\bar{x}=y_{0} \in \Omega_{\frac{C}{\delta} \delta_{0}^{\prime}}^{\prime}$. On the other hand, we know from [8] (see the discussion before (4.12)) that $\nabla \hat{w}(\partial \Omega)=\partial \Omega^{\prime}$ and $|\nabla \hat{w}(x)-\nabla \hat{w}(y)| \leq C|x-y|^{\gamma}$ for all $x, y \in \bar{\Omega}$. Therefore we conclude $\operatorname{dist}(\bar{x}, \partial \Omega) \geq c \delta_{0}^{\gamma^{2}}>0$. But since $\hat{w}$ is a strictly convex function (c.f. [7]), there exists $\mu_{1}>0$ depending on $\delta_{0}, n$, $\Lambda_{1}, \Lambda_{2}, \Omega$ and $\Omega^{\prime}$ such that

$$
\begin{equation*}
T_{\mu_{1}}(w, \bar{x}) \subset \Omega_{\frac{c}{2} \delta_{0}^{\gamma^{2}}}, \tag{4.16}
\end{equation*}
$$

which gives the last inclusion in (4.15) whenever $h+\epsilon \leq \mu_{1}$.
Next for any $\epsilon>0$, by Proposition 4.1 there exists $\lambda_{0}$ depending only on $\epsilon, n, \Lambda_{1}, \Lambda_{2}, \Omega$ and $\Omega^{\prime}$ such that for all $\lambda \geq \lambda_{0}$ we have $\left\|u-\left[w+u\left(x_{0}\right)-w\left(x_{0}\right)\right]\right\|_{L^{\infty}(\Omega)} \leq \epsilon / 4$ (of course, we are abusing the notation, just take $\lambda_{0}$ to be the smallest of all the choices). Since the sections of $w$ are the same as the corresponding sections of $w+u\left(x_{0}\right)-w\left(x_{0}\right)$, we can assume that $\|u-w\|_{L^{\infty}(\Omega)} \leq \epsilon / 4$. This implies that

$$
\begin{equation*}
\|\bar{u}-\bar{w}\|_{L^{\infty}(\Omega)} \leq \frac{\epsilon}{4}+\frac{C}{\lambda} \leq \frac{\epsilon}{2} \quad \text { for all } \quad \lambda \geq \lambda_{0} . \tag{4.17}
\end{equation*}
$$

Now it is easy to see that the first two inclusions in (4.15) hold true. Indeed if $x \in T_{h-\epsilon}(w, \bar{x})=$ $\{x \in \Omega: \bar{w}(x) \leq \bar{w}(\bar{x})+h-\epsilon\}$, then

$$
\bar{u}(x) \leq \bar{w}(x)+\frac{\epsilon}{2} \leq \bar{w}(\bar{x})+h-\frac{\epsilon}{2} \leq \bar{w}\left(x_{0}\right)+h-\frac{\epsilon}{2} \leq \bar{u}\left(x_{0}\right)+h .
$$

On the other hand, if $x \in S_{h}\left(u, x_{0}, y_{0}\right)=\left\{x \in \Omega: \bar{u}(x) \leq \bar{u}\left(x_{0}\right)+h\right\}$ then

$$
\bar{w}(x) \leq \bar{u}(x)+\frac{\epsilon}{2} \leq \bar{u}\left(x_{0}\right)+h+\frac{\epsilon}{2} \leq \bar{u}(\bar{x})+h+\frac{\epsilon}{2} \leq \bar{w}(\bar{x})+h+\epsilon .
$$

These yield $T_{h-\epsilon}(w, \bar{x}) \subset S_{h}\left(u, x_{0}, y_{0}\right) \subset T_{h+\epsilon}(w, \bar{x})$ as desired.
Corollary 4.6. Given $\delta_{0}>0$, let $\mu_{1}>0$ be the corresponding constant given by Proposition 4.5. Then for all $h \leq \frac{\mu_{1}}{2}$, there exists $\lambda_{0}=\lambda_{0}\left(h, \delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \Omega, \Omega^{\prime}\right)>0$ such that if $\lambda \geq \lambda_{0}$ and $u$ is a solution of (4.1), we have the following: for all $x_{0} \in \Omega_{\delta_{0}}$ and $y_{0} \in \partial_{c} u\left(x_{0}\right)$,
(i) $S_{h}\left(u, x_{0}, y_{0}\right) \subset \Omega_{c \delta_{0}^{\theta}}$;
(ii) There is an affine transformation $T x=A x+b$ with $C_{1} \leq|\operatorname{det} A|^{\frac{2}{n}} h \leq C_{2}$ and $\|A\| \leq$ $\mathrm{Ch}^{-1}$ such that

$$
B_{K}(0) \subset \tilde{\Omega}:=T S_{h}\left(u, x_{0}, y_{0}\right) \subset B_{n}(0) .
$$

Here all the constants are universal (depending on $n, \Lambda_{1}$ and $\Lambda_{2}$ ). As a consequence,

$$
\operatorname{dist}(x, \partial[\tilde{\Omega}]) \leq \frac{C h^{-1}}{\lambda} \quad \forall x \in \partial \tilde{\Omega} .
$$

Proof. Applying Proposition 4.5 for $\epsilon=h / 2$, we obtain

$$
T_{\frac{h}{2}}(w, \bar{x}) \subset S_{h}\left(u, x_{0}, y_{0}\right) \subset T_{\frac{3 h}{2}}(w, \bar{x}) \subset \Omega_{c \delta_{0}^{\theta}},
$$

where $w$ is the solution of the quadratic transport problem (4.2). Let $T x=A x+b$ be an affine transformation normalizing $T_{3 h / 2}(w, \bar{x})$, i.e., $B_{1} \subset T\left(T_{3 h / 2}(w, \bar{x})\right) \subset B_{n}$. Then as $\left|T_{3 h / 2}(w, \bar{x})\right| \approx h^{\frac{n}{2}}$ (see [26, Corollary 3.2.4]), it is clear that $C_{1} \leq|\operatorname{det} A|^{\frac{2}{n}} h \leq C_{2}$. Also it follows from Theorem 3.3.8 in [26] that $B_{K} \subset T\left(T_{h / 2}(w, \bar{x})\right)$ with $K>0$ a universal constant. Thus $B_{K} \subset T S_{h}\left(u, x_{0}, y_{0}\right) \subset B_{n}$. Note that as $B_{C h}(\bar{x}) \subset T_{3 h / 2}(w, \bar{x})$, we get $T\left(B_{C h}(\bar{x})\right) \subset B_{n}$ implying that $\|A\| \leq C h^{-1}$. This together with Proposition 3.2 yields the last estimate.

In the above corollary, $f$ and $g$ are only required to satisfy the condition (1.3). However if in addition $f \in \mathcal{C}_{\text {loc }}^{\alpha}(\Omega)$ and $g \in \mathcal{C}_{\text {loc }}^{\alpha}\left(\Omega^{\prime}\right)$ as assumed in Theorem 1.1, then we in fact have the following better estimates

$$
\begin{equation*}
S_{h}\left(u, x_{0}, y_{0}\right) \subset B_{C \sqrt{h}}\left(x_{0}\right), \quad\|A\| \leq C h^{\frac{-1}{2}}, \quad\left\|A^{-1}\right\| \leq C h^{\frac{1}{2}} \tag{4.18}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{dist}(x, \partial[\tilde{\Omega}]) \leq \frac{C h^{\frac{-1}{2}}}{\lambda} \quad \forall x \in \partial \tilde{\Omega} . \tag{4.19}
\end{equation*}
$$

Indeed the regularity of $w$ in [7] gives $C I \leq D^{2} \bar{w} \leq C^{\prime} I$ in $\Omega_{c \delta_{0}^{\theta}}$ which in turn implies that

$$
\begin{equation*}
B_{C_{1} \sqrt{h}}(\bar{x}) \subset T_{3 h / 2}(w, \bar{x}) \subset B_{C_{2} \sqrt{h}}(\bar{x}) . \tag{4.20}
\end{equation*}
$$

Then (4.18) and (4.19) follow from this, the fact $T^{-1}\left(B_{K}(0)\right) \subset S_{h}\left(u, x_{0}, y_{0}\right) \subset T_{3 h / 2}(w, \bar{x})$ and the arguments used in the proof of Corollary 4.6.

Another consequence of our construction is a gradient bound for points in the section $S_{h}:=S_{h}\left(u, x_{0}, y_{0}\right)$. We denote
$\tilde{u}(x):=\frac{1}{h}\left[u\left(T^{-1} x\right)-u\left(x_{0}\right)-c\left(x_{0}-y_{0}\right)-h\right], \quad \tilde{\Omega}:=T S_{h}\left(u, x_{0}, y_{0}\right) \quad$ and $\quad \tilde{c}(x):=\frac{1}{h} c\left(T^{-1} x\right)$,
where $T$ is the transformation from the previous corollary. Then:
Corollary 4.7. Under the same assumptions as in Corollary 4.6, assume further that $f \in$ $\mathcal{C}_{\text {loc }}^{\alpha}(\Omega)$ and $g \in \mathcal{C}_{\text {loc }}^{\alpha}\left(\Omega^{\prime}\right)$, then we have that

$$
\partial_{c}^{l o c} u\left(S_{h}\right) \subset B_{C\left(\sqrt{h}+\frac{1}{\lambda}\right)}\left(y_{0}\right)
$$

and

$$
\partial_{\tilde{c}}^{l o c} \tilde{u}(\tilde{\Omega}) \subset B_{C\left(1+\frac{1}{\lambda \sqrt{h}}\right)}\left(T y_{0}\right) .
$$

Proof. As in the proof of Corollary 4.6, we choose $\epsilon=h / 2$ and have $\|u-w\|_{L^{\infty}(\Omega)} \leq \epsilon / 4$. Also it follows from (4.20) that $S_{h} \subset T_{3 h / 2}(w, \bar{x}) \subset B_{C_{2} \sqrt{h}}(\bar{x})$. Now applying Lemma 3.4 with $\delta:=C \sqrt{h}$, we then obtain $\partial_{c}^{\operatorname{loc}} u\left(S_{h}\right) \subset \mathcal{N}_{\left(C \sqrt{h}+\frac{C}{\lambda}\right)}\left\{\left[\nabla \hat{w}\left(B_{C \sqrt{h}}(\bar{x})\right)\right]\right\}$, where $\hat{w}:=w+\frac{1}{2}|x|^{2}$. Since $w$ is the solution of optimal transport with quadratic cost from $\Omega$ to $\Omega^{\prime}$ and $B_{C \sqrt{h}}(\bar{x})$ is away from the boundary of $\Omega$, we get $\left\|D^{2} w\right\|_{L^{\infty}\left(B_{C \sqrt{h}}(\bar{x})\right)} \leq C$ and hence

$$
\left|\nabla \hat{w}(x)-y_{0}\right|=|\nabla \hat{w}(x)-\nabla \hat{w}(\bar{x})| \leq C|x-\bar{x}| \leq C \sqrt{h} \quad \forall x \in B_{C \sqrt{h}}(\bar{x}) .
$$

Therefore, we conclude that $\partial_{c}^{\mathrm{loc}} u\left(S_{h}\right) \subset B_{C\left(\sqrt{h}+\frac{1}{\lambda}\right)}\left(y_{0}\right)$. The last conclusion follows from the fact $\partial_{\tilde{c}}^{\mathrm{loc}} \tilde{u}\left(T S_{h}\right)=T \partial_{c}^{\mathrm{loc}} u\left(S_{h}\right)$ and the estimate (4.18) for $\|A\|$.

## 5 Interior $\mathcal{C}^{1,1}$ estimates

In this section we present the proof of Theorem 1.1, which is a direct consequence of the following result:

Theorem 5.1. Assume that $f \in \mathcal{C}_{\text {loc }}^{\alpha}(\Omega), g \in \mathcal{C}_{\text {loc }}^{\alpha}\left(\Omega^{\prime}\right)$ and (1.3) holds. Given $\delta_{0}>0$, there exist $\mu_{0}=\mu_{0}\left(\delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \Omega, \Omega^{\prime}\right)>0$ and $\lambda_{0}=\lambda_{0}\left(\delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \Omega, \Omega^{\prime}\right)>0$ such that if $\lambda \geq \lambda_{0}$ and $u$ is a solution of (4.1), then

$$
\begin{equation*}
B_{C^{-1} \sqrt{\rho}}\left(x_{0}\right) \subset S_{\rho}\left(u, x_{0}, y_{0}\right) \subset B_{C \sqrt{\rho}}\left(x_{0}\right) \tag{5.1}
\end{equation*}
$$

for all $x_{0} \in \Omega_{\delta_{0}}, y_{0} \in \partial_{c} u\left(x_{0}\right)$ and $0<\rho \leq \mu_{0}$, where $C>0$ depends only on $n, \Lambda_{1}, \Lambda_{2}$, $\|f\|_{\mathcal{C}^{\alpha}\left(\Omega_{\delta_{1}}\right)}$ and $\|g\|_{\mathcal{C}^{\alpha}\left(\Omega_{\delta_{1}}^{\prime}\right)}$ with $\delta_{1}=\delta_{1}\left(\delta_{0}, n, \Lambda_{1}, \Lambda_{2}\right)$.

Remark 5.2. This immediately implies that $u$ is $\mathcal{C}^{1,1}$.

Proof of Theorem 5.1. Let $\mu_{0}:=\mu^{2}$, with $\mu:=\mu_{1} / 2$ and $\mu_{1}>0$ is the constant given by Proposition 4.5. By applying Proposition 4.1 with $\epsilon_{0}:=\mu^{3}$ we obtain

$$
\|u-w\|_{L^{\infty}(\Omega)} \leq \epsilon_{0}
$$

where $w$ is an optimal transport solution for the quadratic cost between $\Omega$ and $\Omega^{\prime}$, and with densities $f$ and $g$, respectively.

Step 1: Since $x_{0} \in \Omega_{\delta_{0}}$ and $\mu+\epsilon_{0} \leq \mu_{1}$, Proposition 4.5 gives

$$
T_{\mu-\epsilon_{0}}(w, \bar{x}) \subset S_{\mu}\left(u, x_{0}\right):=S_{\mu}\left(u, x_{0}, y_{0}\right) \subset T_{\mu+\epsilon_{0}}(w, \bar{x}) \subset \Omega_{\delta_{1}}
$$

where $\bar{x}$ is the minimum point of $\bar{w}:=w+\frac{1}{2}\left|x-y_{0}\right|^{2}$. Moreover as $\Omega^{\prime}$ is convex and $f, g$ are Hölder continuous we know from Caffarelli's $\mathcal{C}^{2, \alpha}$ estimate in [7] that $w$ is regular in the interior. It then follows by a similar argument yielding (4.20) that $B_{C_{1} \sqrt{\mu-\epsilon_{0}}}(\bar{x}) \subset$ $T_{\mu-\epsilon_{0}}(w, \bar{x})$ and $T_{\mu+\epsilon_{0}}(w, \bar{x}) \subset B_{C_{2} \sqrt{\mu+\epsilon_{0}}}(\bar{x})$. Thus $B_{C_{1} \sqrt{\mu-\epsilon_{0}}}(\bar{x}) \subset S_{\mu}\left(u, x_{0}\right) \subset B_{C_{2} \sqrt{\mu+\epsilon_{0}}}(\bar{x})$. This together with the facts $\left|\bar{x}-x_{0}\right| \leq C \sqrt{\epsilon_{0}}$ (see (6.30) for a proof of this fact) and $\epsilon_{0}$ is very small compared to $\mu$ gives

$$
B_{C \sqrt{\mu-\epsilon_{0}}}\left(x_{0}\right) \subset S_{\mu}\left(u, x_{0}\right) \subset B_{C^{\prime} \sqrt{\mu+\epsilon_{0}}}\left(x_{0}\right) .
$$

For simplicity and without loss of generality we assume that $x_{0}=0$. Therefore, $B_{1}(0) \subset$ $\mu^{\frac{-1}{2}} S_{\mu}\left(u, x_{0}\right) \subset B_{n}(0)$. Let $\Theta(x):=x-\nabla c^{*}(-\nabla u(x))$ be the optimal map corresponding to $u$ and let $\Theta^{*}(y):=y-\nabla c^{*}\left(-\nabla u^{*}(y)\right)$ be the dual optimal map which is known to be the inverse of $\Theta$. Then since $w$ is regular, we know from Lemma 3.4 and the proof of Corollary 4.7 that

$$
\begin{equation*}
\Theta B_{R}\left(x_{0}\right) \subset B_{C\left(R+\frac{1}{\lambda}+\frac{\epsilon_{0}}{R}\right)}\left(y_{0}\right) \quad \forall R \leq C \sqrt{\mu} . \tag{5.2}
\end{equation*}
$$

Likewise as $\left\|u^{*}-w^{*}\right\|_{L^{\infty}\left(\Omega^{\prime}\right)} \leq \epsilon_{0}+\frac{C}{\lambda}$ and the dual potential $w^{*}$ for the quadratic cost is regular thanks to the convexity of $\Omega$, we also have

$$
\begin{equation*}
\Theta^{*} B_{R}\left(y_{0}\right) \subset B_{C\left(R+\frac{1}{\lambda}+\frac{\epsilon_{0}}{R}\right)}\left(x_{0}\right) \quad \text { or } \quad B_{R}\left(y_{0}\right) \subset \Theta B_{C\left(R+\frac{1}{\lambda}+\frac{\epsilon_{0}}{R}\right)}\left(x_{0}\right) \quad \forall R \leq C \sqrt{\mu} . \tag{5.3}
\end{equation*}
$$

Let $\Omega_{1}:=\mu^{\frac{-1}{2}} \Theta^{*}\left[\Theta\left(S_{\mu}\left(u, x_{0}\right)\right)\right], \Omega_{1}^{\prime}:=\mu^{\frac{-1}{2}}\left[\Theta\left(S_{\mu}\left(u, x_{0}\right)\right)\right]$ and

$$
\begin{equation*}
u_{1}(x):=\frac{1}{\mu} u\left(\mu^{\frac{1}{2}} x\right), \quad x \in \mu^{\frac{-1}{2}} \Omega \supset \Omega_{1} . \tag{5.4}
\end{equation*}
$$

Here we are using the notation $[A]$ for the convex hull of the set $A$. The advantage of this choice of domains lies in the fact that $\Omega_{1}^{\prime}$ is convex. On the other hand, by Lemma 5.3, $u_{1} \mid \Omega_{1}$ is the optimal transport solution between $\Omega_{1}$ and $\Omega_{1}^{\prime}$ with densities $f_{1}(x):=f\left(\mu^{\frac{1}{2}} x\right)$, $g_{1}(y):=g\left(\mu^{\frac{1}{2}} y\right)$ and with cost

$$
c_{1}(x):=\frac{1}{\mu} c\left(\mu^{\frac{1}{2}} x\right)=\frac{1}{2}|x|^{2}+\frac{1}{\lambda_{1}} E_{1}(x) \quad \text { with } \quad \frac{1}{\lambda_{1}}:=\frac{\mu^{1 / 2}}{\lambda} .
$$

Note that all the dilations of regions or variables in $\Omega$ are with respect to $x_{0}=0$ and in $\Omega^{\prime}$ are always with respect to $y_{0}$. Moreover by using (5.2) and (5.3) we obtain

$$
B_{1}(0) \subset \mu^{\frac{-1}{2}} S_{\mu}\left(u, x_{0}\right) \subset \Omega_{1} \subset \mu^{\frac{-1}{2}} \Theta^{*} B_{C^{\prime} \sqrt{\mu+\epsilon_{0}}}\left(y_{0}\right) \subset \mu^{\frac{-1}{2}} B_{C \sqrt{\mu+\epsilon_{0}}}\left(x_{0}\right) \subset B_{n}(0)
$$

and

$$
B_{1}\left(y_{0}\right) \subset \mu^{\frac{-1}{2}} B_{C \sqrt{\mu-\epsilon_{0}}}\left(y_{0}\right) \subset \mu^{\frac{-1}{2}} \Theta S_{\mu}\left(u, x_{0}\right) \subset \Omega_{1}^{\prime} \subset \mu^{\frac{-1}{2}} B_{C^{\prime} \sqrt{\mu+\epsilon_{0}}}\left(y_{0}\right) \subset B_{n}\left(y_{0}\right)
$$

Hence we can apply Proposition 4.2 to get

$$
\begin{equation*}
\left\|u_{1}-w_{1}\right\|_{L^{\infty}\left(B_{1}\right)} \leq \epsilon_{0} \tag{5.5}
\end{equation*}
$$

where $w_{1}$ is the optimal transport solution for the quadratic cost between $\Omega_{1}$ and $\Omega_{1}^{\prime}$, and with densities $f_{1}$ and $g_{1}$.

Step 2: Let $\Theta_{1}(x):=x-\nabla c_{1}^{*}\left(-\nabla u_{1}(x)\right)=\mu^{\frac{-1}{2}} \Theta\left(\mu^{\frac{1}{2}} x\right)$ be the optimal map corresponding to $u_{1}$ and let $\Theta_{1}^{*}(y):=y-\nabla c_{1}^{*}\left(-\nabla u_{1}^{*}(y)\right)=\mu^{\frac{-1}{2}} \Theta^{*}\left(\mu^{\frac{1}{2}} y\right)$ be its dual optimal map. As the target domain $\Omega_{1}^{\prime}$ is convex and $f_{1}, g_{1}$ are Hölder continuous with $\mathcal{C}^{\alpha}$ norms are controlled by those of $f$ and $g$, we know from Caffarelli's $\mathcal{C}^{2, \alpha}$ estimate in [7] for quadratic cost that $w_{1}$ is regular in the interior. Therefore, by using (5.5) and arguing as in Step 1 we get

$$
\begin{equation*}
\Theta_{1} B_{R}\left(x_{0}\right) \subset B_{C\left(R+\frac{1}{\lambda_{1}}+\frac{\epsilon_{0}}{R}\right)}\left(y_{0}\right) \quad \forall R \leq C \sqrt{\mu} \tag{5.6}
\end{equation*}
$$

We would like to have (5.3) for $\Theta_{1}^{*}$ but this is not immediate since we do not know whether $w_{1}^{*}$ is regular due to the nonconvexity of $\Omega_{1}$. Observe that we do not really need $\Theta_{1}^{*}$ to be well behaved in every interior subregion of $\Omega_{1}^{\prime}$, but only in the subregion $B_{C \sqrt{\mu}}\left(y_{0}\right) \subset \Omega_{1}^{\prime}$. To obtain this regularity we proceed as follows.

Let $\tilde{\Omega}_{1}:=\mu^{\frac{-1}{2}}\left[S_{\mu}\left(u, x_{0}\right)\right]$ and $\tilde{\Omega}_{1}^{\prime}:=\mu^{\frac{-1}{2}} \Theta\left[S_{\mu}\left(u, x_{0}\right)\right]$. Then $\left.u_{1}\right|_{\tilde{\Omega}_{1}}$ is the optimal transport solution between $\tilde{\Omega}_{1}$ and $\tilde{\Omega}_{1}^{\prime}$ with cost $c_{1}$ and with densities $f_{1}$ and $g_{1}$. As above, by using (5.2) and (5.3) we also have $B_{1}(0) \subset \tilde{\Omega}_{1} \subset B_{n}(0)$ and $B_{1}\left(y_{0}\right) \subset \tilde{\Omega}_{1}^{\prime} \subset B_{n}\left(y_{0}\right)$. Note that even though we can not compare $\tilde{\Omega}_{1}^{\prime}$ to $\Omega_{1}^{\prime}$, we know that $\tilde{\Omega}_{1}^{\prime}$ contains the big ball $B_{1}\left(y_{0}\right)$ centered at $y_{0}$ and as mentioned above this is enough for our purpose. Hence we can apply Proposition 4.2 to get $\left\|u_{1}-\tilde{w}_{1}\right\|_{L^{\infty}\left(B_{1}\right)} \leq \epsilon_{0}$ implying

$$
\left\|u_{1}^{*}-\tilde{w}_{1}^{*}\right\|_{L^{\infty}\left(B_{1}\right)} \leq \epsilon_{0}+\frac{C}{\lambda_{1}} .
$$

Here $\tilde{w}_{1}^{*}$ is the optimal transport solution for the quadratic cost between $\tilde{\Omega}_{1}^{\prime}$ and $\tilde{\Omega}_{1}$, and with densities $g_{1}$ and $f_{1}$. Moreover as $\tilde{\Omega}_{1}$ is convex, we know that $\tilde{w}_{1}^{*}$ is regular in the interior of $\tilde{\Omega}_{1}^{\prime}$. Therefore, we also have

$$
\begin{equation*}
\Theta_{1}^{*} B_{R}\left(y_{0}\right) \subset B_{C\left(R+\frac{1}{\lambda_{1}}+\frac{\epsilon_{0}}{R}\right)}\left(x_{0}\right) \quad \text { or } \quad B_{R}\left(y_{0}\right) \subset \Theta_{1} B_{C\left(R+\frac{1}{\lambda_{1}}+\frac{\epsilon_{0}}{R}\right)}\left(x_{0}\right) \tag{5.7}
\end{equation*}
$$

From (5.6) and (5.7), we see that $\Theta_{1}$ and $\Theta_{1}^{*}$ enjoy the same property as $\Theta$ and $\Theta^{*}$ given by (5.2) and (5.3). This allows us to carry the next level of the rescaling process. Indeed, let $\Omega_{2}:=\mu^{-1} \Theta^{*}\left[\Theta\left(S_{\mu^{2}}\left(u, x_{0}\right)\right)\right], \Omega_{2}^{\prime}:=\mu^{-1}\left[\Theta\left(S_{\mu^{2}}\left(u, x_{0}\right)\right)\right]$ and

$$
u_{2}(x):=\frac{1}{\mu^{2}} u(\mu x)=\frac{1}{\mu} u_{1}\left(\mu^{\frac{1}{2}} x\right) \quad \text { for } \quad x \in \mu^{-1} \Omega \supset \mu^{\frac{-1}{2}} \Omega_{1} \supset \Omega_{2} .
$$

Then by Lemma 5.3, $\left.u_{2}\right|_{\Omega_{2}}$ is the optimal transport solution between $\Omega_{2}$ and $\Omega_{2}^{\prime}$ with cost $c_{2}(x):=\frac{1}{\mu^{2}} c(\mu x)$, and with densities $f_{2}(x):=f(\mu x)$ and $g_{2}(y):=g(\mu y)$. Notice that $\Omega_{2}=\mu^{\frac{-1}{2}} \Theta_{1}^{*}\left[\Theta_{1}\left(S_{\mu}\left(u_{1}, x_{0}\right)\right)\right]$ and $\Omega_{2}^{\prime}=\mu^{\frac{-1}{2}}\left[\Theta_{1}\left(S_{\mu}\left(u_{1}, x_{0}\right)\right)\right]$. Hence by arguing as in Step 1 and using (5.6) and (5.7) we obtain

$$
\left\|u_{2}-w_{2}\right\|_{L^{\infty}\left(B_{1}\right)} \leq \epsilon_{0}
$$

with $w_{2}$ is an optimal transport solution for the quadratic cost between $\Omega_{2}$ and $\Omega_{2}^{\prime}$, and with densities $f_{2}$ and $g_{2}$.

Step $k$ : Repeating the above process we get

$$
\left\|u_{k}-w_{k}\right\|_{L^{\infty}\left(B_{1}\right)} \leq \epsilon_{0},
$$

where $u_{k}(x):=\frac{1}{\mu^{k}} u\left(\mu^{\frac{k}{2}} x\right)$ and $w_{k}$ is the optimal transport solution for the quadratic cost between $\Omega_{k}:=\mu^{\frac{-k}{2}} \Theta^{*}\left[\Theta\left(S_{\mu^{k}}\left(u, x_{0}\right)\right)\right]$ and $\Omega_{k}^{\prime}:=\mu^{\frac{-k}{2}}\left[\Theta\left(S_{\mu^{k}}\left(u, x_{0}\right)\right)\right]$, and with densities $f_{k}(x):=f\left(\mu^{\frac{k}{2}} x\right)$ and $g_{k}(y):=g\left(\mu^{\frac{k}{2}} y\right)$. Since $0<\rho \leq \mu_{0}=\mu^{2}$, we can pick the positive integer number $k$ such that $\mu^{k+2}<\rho \leq \mu^{k+1}$. Thus

$$
\mu^{\frac{k}{2}} T_{\frac{\rho}{\mu^{k}}-\epsilon_{0}}\left(w_{k}, \bar{x}_{k}\right) \subset S_{\rho}\left(u, x_{0}\right)=\mu^{\frac{k}{2}} S_{\frac{\rho}{\mu^{k}}}\left(u_{k}, x_{0}\right) \subset \mu^{\frac{k}{2}} T_{\frac{\rho}{\mu^{k}}+\epsilon_{0}}\left(w_{k}, \bar{x}_{k}\right) .
$$

Here $\bar{x}_{k}$ denotes the minimum point of $\bar{w}_{k}:=w_{k}+\frac{1}{2}\left|x-y_{0}\right|^{2}$. As the target domain $\Omega_{k}^{\prime}$ is convex and $f_{k}, g_{k}$ are Hölder continuous with $\mathcal{C}^{\alpha}$ norms are controlled by those of $f$ and $g$, we know that $B_{C_{1} \sqrt{\frac{\rho}{\mu^{k}}-\epsilon_{0}}}\left(\bar{x}_{k}\right) \subset T_{\frac{\rho}{\mu^{k}}-\epsilon_{0}}\left(w_{k}, \bar{x}_{k}\right)$ and $T_{\frac{\rho}{\mu^{k}}+\epsilon_{0}}\left(w_{k}, \bar{x}_{k}\right) \subset B_{C_{2} \sqrt{\frac{\rho}{\mu^{k}}+\epsilon_{0}}}\left(\bar{x}_{k}\right)$. It follows from these and the fact $\left|\bar{x}_{k}-x_{0}\right| \leq C \sqrt{\epsilon_{0}}$ that

$$
S_{\rho}\left(u, x_{0}\right) \subset \mu^{\frac{k}{2}} B_{C \sqrt{\frac{\rho}{\mu^{k}}+\epsilon_{0}}}\left(x_{0}\right) \subset B_{C \sqrt{\rho+\mu^{k} \epsilon_{0}}}\left(x_{0}\right) \subset B_{C \sqrt{\rho+\rho \epsilon_{0} \mu^{2}}}\left(x_{0}\right) \subset B_{C \sqrt{\rho}}\left(x_{0}\right) .
$$

Similarly, we also have $B_{C^{\prime} \sqrt{\rho}}\left(x_{0}\right) \subset S_{\rho}\left(u, x_{0}\right)$ yielding, (5.1) as desired.

In the above proof, we have used the following simple lemma:
Lemma 5.3. Let $u_{1}$ be defined by (5.4). Then $u_{1}$ is the optimal transport solution between $\Omega_{1}:=\mu^{\frac{-1}{2}} \Theta^{*}\left[\Theta\left(S_{\mu}\left(u, x_{0}\right)\right)\right]$ and $\Omega_{1}^{\prime}:=\mu^{\frac{-1}{2}}\left[\Theta\left(S_{\mu}\left(u, x_{0}\right)\right)\right]$ with cost $c_{1}(x):=\frac{1}{\mu} c\left(\mu^{\frac{1}{2}} x\right)$, and with densities $f_{1}(x):=f\left(\mu^{\frac{1}{2}} x\right)$ and $g_{1}(y):=g\left(\mu^{\frac{1}{2}} y\right)$. The similar statement also holds for the function $u_{k}$.

Proof. Let $\Theta_{1}(x):=x-\nabla c_{1}^{*}\left(-\nabla u_{1}(x)\right)$. Then as $u_{1}$ is $c_{1}$-convex, it suffices to show that $\Theta_{1 \#} f_{1} \chi_{\Omega_{1}} d x=g_{1} \chi_{\Omega_{1}^{\prime}} d y$. Now since $c_{1}^{*}(x)=\frac{1}{\mu} c^{*}\left(\mu^{\frac{1}{2}} x\right)$, we obtain

$$
\Theta_{1}(x)=x-\mu^{\frac{-1}{2}} \nabla c^{*}\left(-\nabla u\left(\mu^{\frac{1}{2}} x\right)\right)=\mu^{\frac{-1}{2}}\left[\mu^{\frac{1}{2}} x-\nabla c^{*}\left(-\nabla u\left(\mu^{\frac{1}{2}} x\right)\right)\right]=\mu^{\frac{-1}{2}} \Theta\left(\mu^{\frac{1}{2}} x\right) .
$$

It follows that for any function $h \in \mathcal{C}\left(\Omega_{1}^{\prime}\right)$, by letting $\tilde{h}(x):=h\left(\mu^{\frac{-1}{2}} x\right)$ we have

$$
\begin{aligned}
\int_{\Omega_{1}} h\left(\Theta_{1}(x)\right) f_{1}(x) d x & =\int_{\mu^{\frac{-1}{2}} \Theta^{*}\left[\Theta\left(S_{\mu}\left(u, x_{0}\right)\right)\right]} h\left(\mu^{\frac{-1}{2}} \Theta\left(\mu^{\frac{1}{2}} x\right)\right) f\left(\mu^{\frac{1}{2}} x\right) d x \\
& =\mu^{\frac{-n}{2}} \int_{\Theta^{*}\left[\Theta\left(S_{\mu}\left(u, x_{0}\right)\right)\right]} \tilde{h}(\Theta(z)) f(z) d z \\
& =\mu^{\frac{-n}{2}} \int_{\left[\Theta\left(S_{\mu}\left(u, x_{0}\right)\right)\right]} \tilde{h}\left(y^{\prime}\right) g\left(y^{\prime}\right) d y^{\prime} \\
& =\int_{\mu^{\frac{-1}{2}}\left[\Theta\left(S_{\mu}\left(u, x_{0}\right)\right)\right]} \tilde{h}\left(\mu^{\frac{1}{2}} y\right) g\left(\mu^{\frac{1}{2}} y\right) d y=\int_{\Omega_{1}^{\prime}} h(y) g_{1}(y) d y .
\end{aligned}
$$

Thus $\Theta_{1 \#} f_{1} \chi_{\Omega_{1}} d x=g_{1} \chi_{\Omega_{1}^{\prime}} d y$ and the proof is completed.

## 6 Comparison Principles

We consider now the main ingredients in the proof of Theorem 1.2. Our strategy is to compare $\tilde{u}$ with the solution $\tilde{w}$ of a standard Monge-Ampère equation with given Dirichlet data at the boundary. However, the difficulty is to obtain an explicit estimate for the difference between $\tilde{u}$ and $\tilde{w}$. This task will be achieved by constructing some suitable barriers for the operator $M_{\lambda}$ in a certain class of admissible functions. We remind that

$$
M_{\lambda}[u]:=g(x+\nabla u+G(\nabla u)) \operatorname{det}\left(I+D^{2} u+H(\nabla u)\left(D^{2} u\right)\right) \text {, }
$$

where $G$ and $H$ in particular satisfy (2.11).
To motivate our choice of admissible functions, let us seek a pointwise condition guaranteeing that a function $\phi \in \mathcal{C}^{2}(\Omega)$ is $c_{\lambda}$-convex in $\Omega$. Note that $\phi$ is $c_{\lambda}$-convex at $x_{0}$ if and only if the function $h_{0}(x):=\phi(x)+c_{\lambda}\left(x-y_{0}\right)$ has a minimum point in $\Omega$ at $x_{0}$ or equivalently $h_{0}$ has a supporting plane in $\Omega$ at $x_{0}$, where $y_{0}:=x_{0}-\nabla c_{\lambda}^{*}\left(-\nabla \phi\left(x_{0}\right)\right)$. But for $\Omega \subset B_{n}$, we have

$$
\begin{aligned}
D^{2} h_{0}(x) & =I+D^{2} \phi(x)+\frac{1}{\lambda} D^{2} E\left(x-y_{0}\right) \geq I+D^{2} \phi(x)-\frac{C}{\lambda} I\left|x-y_{0}\right| \\
& \geq I+D^{2} \phi(x)-\frac{C}{\lambda} I\left(\left|x_{0}+\nabla \phi\left(x_{0}\right)\right|+2 n\right) \quad \forall x \in \Omega .
\end{aligned}
$$

Thus if $\Omega \subset B_{n}$ is convex and

$$
\inf _{\Omega}\left(I+D^{2} \phi\right)>\frac{C}{\lambda} I\left(\|i d+\nabla \phi\|_{L^{\infty}(\Omega)}+2 n\right),
$$

then for any $x_{0} \in \Omega$, the function $h_{0}$ has a supporting plane in $\Omega$ at $x_{0}$. This implies that $\phi$ is $c_{\lambda}$-convex in $\Omega$ and in particular $I+D^{2} c_{\lambda}^{*}(-\nabla \phi(x)) D^{2} \phi(x) \geq 0$ in $\Omega$. We will choose the class of admissible functions to be those satisfying the above inequality.

Definition 6.1. Let $\Omega \subset B_{n}$ be an open set. A function $\phi$ is said to be in the admissible class for the operator $M_{\lambda}$ in $\Omega$ if $\phi \in \mathcal{C}^{2}([\Omega])$ and

$$
\begin{equation*}
\inf _{[\Omega]}\left(I+D^{2} \phi\right)>\frac{C}{\lambda} I\left(\|i d+\nabla \phi\|_{L^{\infty}([\Omega])}+2 n\right) \tag{6.1}
\end{equation*}
$$

Remark 6.2. Let $u \in \mathcal{C}(\Omega)$ be a $c_{\lambda}$ - convex function. Then if $\phi$ is in the admissible class for $M_{\lambda}$ in $\Omega$ and $u-\phi$ has a maximum or minimum at $x_{0} \in \Omega$, we have $x_{0}-\nabla c_{\lambda}^{*}\left(-\nabla \phi\left(x_{0}\right)\right) \in$ $\partial_{c_{\lambda}} u\left(x_{0}\right)$.

Since Alexandroff weak solutions are expected to be stable with respect to a small perturbation, $u$ should be close to a Monge-Ampère solution if [ $\Omega$ ] is close to $\Omega$ (that is, $\Omega$ is almost convex) and $\lambda$ is large enough. To quantify this closeness precisely, we will need the following property for Alexandroff weak solutions:

Proposition 6.3. Let $\Omega \subset B_{n}$ be an open set and $f \in \mathcal{C}(\Omega), g \in \mathcal{C}\left(\Omega^{\prime}\right)$. Let $u \in \mathcal{C}(\bar{\Omega})$ be a $c_{\lambda}$-convex function satisfying $\partial_{c_{\lambda}} u(\Omega) \subset \Omega^{\prime}$ and

$$
\int_{\partial_{c_{\lambda}} u(E)} g(y) d y=\int_{E} f(x) d x \quad \text { for all Borel sets } \quad E \subset \Omega .
$$

Suppose that $\phi \in \mathcal{C}^{2}([\Omega]) \cap \mathcal{C}(\bar{\Omega})$ is a function in the admissible class for the operator $M_{\lambda}$ in $\Omega$. Then:
(i) If $M_{\lambda}[\phi]<f$ at any point $x \in \Omega$ satisfying $x-\nabla c_{\lambda}^{*}(-\nabla \phi(x)) \in \partial_{c_{\lambda}} u(\Omega)$, then

$$
\max _{\bar{\Omega}}\{u-\phi\}=\max _{\partial \Omega}\{u-\phi\} .
$$

(ii) If $M_{\lambda}[\phi]>f$ at any point $x \in \Omega$ satisfying $x-\nabla c_{\lambda}^{*}(-\nabla \phi(x)) \in \partial_{c_{\lambda}} u(\Omega)$, then

$$
\min _{\bar{\Omega}}\{u-\phi\}=\min _{\partial \Omega}\{u-\phi\} .
$$

Proof. (i) Suppose by contradiction that $\max _{\bar{\Omega}}\{u-\phi\}>\max _{\partial \Omega}\{u-\phi\}$. Then there exists $\tilde{x} \in \Omega$ such that $u(\tilde{x})-\phi(\tilde{x})=\max _{\bar{\Omega}}\{u-\phi\}$, and so $u-\phi$ has a maximum in $\bar{\Omega}$ at $\tilde{x}$. Notice that $\tilde{x}-\nabla c_{\lambda}^{*}(-\nabla \phi(\tilde{x})) \in \partial_{c_{\lambda}} u(\tilde{x}) \subset \partial_{c_{\lambda}} u(\Omega)$ by the remark after Definition 6.1. We can assume further that $u(\tilde{x})=\phi(\tilde{x})$ and $u(x)<\phi(x)$ for all $x \in \bar{\Omega} \backslash\{\tilde{x}\}$. This can be achieved by adding $u(\tilde{x})-\phi(\tilde{x})+r|x-\tilde{x}|^{2}$ to $\phi$ and later taking $r \rightarrow 0^{+}$. Indeed, the only thing we need to check is the function $\tilde{\phi}(x):=\phi(x)+u(\tilde{x})-\phi(\tilde{x})+r|x-\tilde{x}|^{2}$ also belongs to the admissible class for $M_{\lambda}$ in $\Omega$. But this is true because

$$
\begin{align*}
\inf _{[\Omega]}\left(I+D^{2} \tilde{\phi}\right)=\inf _{[\Omega]}\left(I+D^{2} \phi\right)+2 r I & >\frac{C}{\lambda} I\left(\|i d+\nabla \phi\|_{L^{\infty}([\Omega])}+2 n\right)+2 r I \\
& \geq \frac{C}{\lambda} I\left(\|i d+\nabla \tilde{\phi}\|_{L^{\infty}([\Omega])}+2 n\right)+2 r\left(1-\frac{2 n C}{\lambda}\right) I \\
& \geq \frac{C}{\lambda} I\left(\|i d+\nabla \tilde{\phi}\|_{L^{\infty}([\Omega])}+2 n\right) . \tag{6.2}
\end{align*}
$$

Let $m:=\min _{\delta / 2 \leq|x-\tilde{x}| \leq \delta}\{\phi(x)-u(x)\}>0$, where $\delta>0$ is a small number. Fix $0<\epsilon<m$ and set

$$
E_{\epsilon}:=\left\{x \in B_{\delta}(\tilde{x}): u(x)+\epsilon>\phi(x)\right\} .
$$

It is easy to see that $E_{\epsilon} \subset B_{\delta / 2}(\tilde{x})$, because if $\delta / 2 \leq|x-\tilde{x}|<\delta$, then $\phi(x)-u(x) \geq m$, so that $x \notin E_{\epsilon}$. Note also that $E_{\epsilon} \downarrow\{\tilde{x}\}$ as $\epsilon \rightarrow 0^{+}$, and $u+\epsilon=\phi$ on $\partial E_{\epsilon}$.

Since $u$ is an Alexandroff weak solution in $\Omega$, we have

$$
\begin{equation*}
\int_{E_{\epsilon}} f(x) d x=\int_{\partial_{c} u\left(E_{\epsilon}\right)} g(y) d y=\int_{\partial_{c}(u+\epsilon)\left(E_{\epsilon}\right)} g(y) d y \tag{6.3}
\end{equation*}
$$

But $\partial_{c}(u+\epsilon)\left(E_{\epsilon}\right) \subset \partial_{c}\left(\phi, E_{\epsilon}\right)\left(E_{\epsilon}\right)=\partial_{c} \phi\left(E_{\epsilon}\right)$ by Lemma 3.3 and the fact $\phi$ is a smooth $c$-convex function in $\Omega$. Thus the chain of equalities (6.3) continues as follows:
$(6.3) \leq \int_{\partial_{c} \phi\left(E_{\epsilon}\right)} g(y) d y=w_{c}(g, \phi)\left(E_{\epsilon}\right)=\int_{E_{\epsilon}} g\left(x-D c^{*}(-\nabla \phi)\right) \operatorname{det}\left(I+D^{2} c^{*}(-\nabla \phi) D^{2} \phi\right) d x$,
where the last equality is due to the representation (2.5). Next, because $f$ and $g$ are continuous, we can let $\epsilon \rightarrow 0$ to obtain

$$
M_{\lambda}[\phi] \geq f \quad \text { at } \quad \tilde{x},
$$

which is a contradiction with the hypothesis in (i).
(ii) The second statement follows in a similar manner. Suppose by contradiction that $\min _{\bar{\Omega}}\{u-\phi\}<\min _{\partial \Omega}\{u-\phi\}$. Then there exists $\tilde{x} \in \Omega$ such that $u-\phi$ has a minimum in $\bar{\Omega}$ at $\tilde{x}$. Notice that $\tilde{x}-\nabla c_{\lambda}^{*}(-\nabla \phi(\tilde{x})) \in \partial_{c_{\lambda}} u(\tilde{x})$ by the remark after Definition 6.1. We can assume in addition that $u(\tilde{x})=\phi(\tilde{x})$ and $u(x)>\phi(x)$ for all $x$ in $\bar{\Omega} \backslash\{\tilde{x}\}$. This can be
achieved by adding $u(\tilde{x})-\phi(\tilde{x})-r|x-\tilde{x}|^{2}$ to $\phi$ and later taking $r \rightarrow 0^{+}$. Indeed, we only need to check that the function $\tilde{\phi}(x):=\phi(x)+u(\tilde{x})-\phi(\tilde{x})-r|x-\tilde{x}|^{2}$ is in the admissible class for $M_{\lambda}$ in $\Omega$ for all $r>0$ sufficiently small. But as done in (6.2), we have

$$
\begin{aligned}
& \inf _{[\Omega]}\left(I+D^{2} \tilde{\phi}\right)-\frac{C}{\lambda} I\left(\|i d+\nabla \tilde{\phi}\|_{L^{\infty}([\Omega])}+2 n\right) \\
& \geq \inf _{[\Omega]}\left(I+D^{2} \phi\right)-\frac{C}{\lambda} I\left(\|i d+\nabla \phi\|_{L^{\infty}([\Omega])}+2 n\right)-2 r\left(1+\frac{2 n C}{\lambda}\right) I \\
& >0 \text { for all } r>0 \text { small enough. }
\end{aligned}
$$

Let $m:=\min _{\bar{\Omega} \backslash B_{\delta}(\tilde{x})}\{u(x)-\phi(x)\}>0$, where $\delta>0$ is small. Fix $0<\epsilon<m$ and set

$$
E_{\epsilon}:=\{x \in \Omega: u(x)-\epsilon<\phi(x)\} .
$$

It is easy to see that $E_{\epsilon} \subset B_{\delta}(\tilde{x}) \subset \Omega$, because if $|x-\tilde{x}| \geq \delta$, then $u(x)-\phi(x) \geq m$, so that $x \notin E_{\epsilon}$. Note also that $E_{\epsilon} \downarrow\{\tilde{x}\}$ as $\epsilon \rightarrow 0^{+}$and $u-\epsilon=\phi$ on $\partial E_{\epsilon}$.

Now it follows from the assumption for $u$ that $\int_{\partial_{c} u\left(E_{\epsilon}\right)} g(y) d y=\int_{E_{\epsilon}} f(x) d x$. But as a consequence of the proof of Lemma 3.3 and the fact $\phi(x) \leq u(x)-\epsilon$ for $x \in \Omega \backslash E_{\epsilon}$, we get $\partial_{c} \phi\left(E_{\epsilon}\right) \subset \partial_{c} u\left(E_{\epsilon}\right)$. Therefore we obtain

$$
\int_{E_{\epsilon}} g\left(x-D c^{*}(-\nabla \phi)\right) \operatorname{det}\left(I+D^{2} c^{*}(-\nabla \phi) D^{2} \phi\right) d x=\int_{\partial_{c} \phi\left(E_{\epsilon}\right)} g(y) d y \leq \int_{E_{\epsilon}} f(x) d x
$$

where the first equality is due to (2.5) and the fact $\phi$ is $c$-convex in $\Omega$. Next, we let $\epsilon \rightarrow 0$ to obtain $M_{\lambda}[\phi] \leq f$ at $\tilde{x}$, which is a contradiction with the hypothesis in (ii).

### 6.1 An approximation result

Our objective here is to establish the main step in the proof of Theorem 1.2, which allows us to compare weak solutions coming from two different equations. This approximation result is proved by constructing very careful barriers to which we apply the previous comparison principle:

Theorem 6.4. Let $B_{1} \subset \Omega \subset B_{n}$ be a domain in $\mathbb{R}^{n}$, not necessarily convex, but such that

$$
\begin{equation*}
\operatorname{dist}(x, \partial[\Omega]) \leq \frac{d}{\lambda} \quad \forall x \in \partial \Omega \tag{6.4}
\end{equation*}
$$

Assume that $f \in \mathcal{C}(\Omega)$ and $g \in \mathcal{C}\left(\Omega^{\prime}\right)$. Let $u \in \mathcal{C}(\bar{\Omega})$ be a $c_{\lambda}$-convex function such that $\partial_{c_{\lambda}} u(\Omega) \subset \Omega^{\prime}$ and $u$ is an Alexandroff weak solution of the problem

$$
\left\{\begin{align*}
M_{\lambda}[u] & =f & & \text { in } \Omega  \tag{6.5}\\
u(x) & =-c_{\lambda}(x) & & \text { on } \partial \Omega
\end{align*}\right.
$$

where $M_{\lambda}$ is the operator defined in (2.10) satisfying (2.11). Consider also the $c_{q}$-convex solution $w \in \mathcal{C}^{3}(\Omega) \cap \mathcal{C}(\bar{\Omega})$ of

$$
\left\{\begin{align*}
\operatorname{det}\left(I+D^{2} w\right) & =1 & & \text { in } \quad V:=\mathcal{N}_{\left(\frac{d}{\lambda}\right)^{\beta}}[\Omega]  \tag{6.6}\\
w & =-\frac{1}{2}|x|^{2} & & \text { on } \quad \partial V,
\end{align*}\right.
$$

where $0<\beta \leq 1$. Assume also that

1. For some $\theta>0$ and $0<\alpha<1$,

$$
\begin{equation*}
|f(x)-1| \leq \theta \quad \text { in } \quad \Omega \quad \text { and } \quad|g(y)-1| \leq \theta|y|^{\alpha} \quad \text { in } \quad \partial_{c_{\lambda}} u(\Omega) . \tag{6.7}
\end{equation*}
$$

2. For a universal constant $C_{n}$ given precisely in the proof,

$$
\begin{equation*}
C_{n}\left(\frac{d}{\lambda}\right)^{-\beta\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)}\left[\theta+\frac{1}{\lambda}\left(\frac{d}{\lambda}\right)^{\frac{n-2}{n}(\alpha-1) \beta}\right] \leq 1 \tag{6.8}
\end{equation*}
$$

Then we have

$$
\begin{equation*}
-C\left(\frac{d}{\lambda}\right)^{\beta(n-1)^{2}} \leq u-w \leq C\left(\frac{d}{\lambda}\right)^{\frac{2 \beta}{n}} \quad \text { in } \Omega \tag{6.9}
\end{equation*}
$$

In case that $[\Omega]$ (equivalently $V$ ) not strictly convex, we can modify $V$ slightly so that $V$ is strictly convex and still $V \approx \mathcal{N}_{\left(\frac{d}{\lambda}\right)^{\beta}}[\Omega]$. Now we note that $w$ exists because $\bar{w}=$ $\frac{1}{2}|x|^{2}+w$ is the solution of a standard Monge-Ampère equation with smooth right hand side and homogeneous Dirichlet boundary data, and thus, $\bar{w} \in \mathcal{C}^{3}(V)$ and it is strictly convex. However, the estimates on $\bar{w}$ (in particular, Pogorelov estimates for the second derivatives) are not uniform up to the boundary. To handle this issue, we will use Proposition 6.5 and Proposition 6.6 below concerning the behavior of $\bar{w}$ near the boundary when we do not have extra information on the domain.

The following is a barrier argument by Caffarelli [6]:
Proposition 6.5. Let $\Omega_{1} \subset B_{n}$ be a convex domain in $\mathbb{R}^{n}$ and $\bar{w} \in \mathcal{C}\left(\bar{\Omega}_{1}\right)$ be a generalized solution of the Monge-Ampère equation

$$
\left\{\begin{array}{rlll}
\operatorname{det}\left(D^{2} \bar{w}\right) & \leq 1 & \text { in } & \Omega_{1}, \\
\bar{w} & =0 & \text { on } & \partial \Omega_{1} .
\end{array}\right.
$$

Then

$$
\bar{w}(x) \geq-C(n) \operatorname{dist}\left(x, \partial \Omega_{1}\right)^{\frac{2}{n}} \quad \text { if } \quad n \geq 3
$$

and

$$
\bar{w}(x) \geq-C(\rho) \operatorname{dist}\left(x, \partial \Omega_{1}\right)^{\rho} \quad \text { for any } \rho \in(0,1) \text { if } \quad n=2 .
$$

Pogorelov type estimates for the second derivatives of solutions for a Monge-Ampère equation hold. Although it is well known that in the interior of the domain, the Hessian matrix of the solution is positive definite, the constants deteriorate at the boundary. The next lemma by Caffarelli-Li [3] quantifies this deterioration.

Proposition 6.6. Let $\Omega_{1} \subset \mathbb{R}^{n}$ be an open convex subset satisfying $B_{1} \subset \Omega_{1} \subset B_{n}$, and let $\bar{w} \in \mathcal{C}^{2}\left(\Omega_{1}\right) \cap \mathcal{C}\left(\bar{\Omega}_{1}\right)$ be a convex solution of

$$
\left\{\begin{array}{rlll}
\operatorname{det}\left(D^{2} \bar{w}\right) & =1 & \text { in } \quad \Omega_{1}, \\
\bar{w} & =0 & \text { on } & \partial \Omega_{1} .
\end{array}\right.
$$

Then

$$
\left|D^{2} \bar{w}(x)\right| \leq C^{\prime}(n) \operatorname{dist}\left(x, \partial \Omega_{1}\right)^{1-n} \quad \text { for } \quad n \geq 3
$$

and

$$
\left|D^{2} \bar{w}(x)\right| \leq C^{\prime}(\rho) \operatorname{dist}\left(x, \partial \Omega_{1}\right)^{1-\frac{2}{\rho}} \quad \text { for } \quad n=2
$$

where $\rho$ is any real number in $(0,1)$.

Proof of Theorem 6.4: Without loss of generality, assume that $n \geq 3$. The proof reduces to finding suitable smooth sub and supersolutions of the equation $M_{\lambda}[u]=f$, and using the comparison principle for Alexandroff solutions given in Proposition 6.3 in the domain $\Omega$.

Let $w$ be the solution of problem (6.6). We will show that the functions

$$
w^{-}=w+\frac{\delta}{2}|x|^{2} \quad \text { and } \quad w^{+}=w-\frac{\delta}{2}|x|^{2}
$$

are a sub and a supersolution, respectively, of problem (6.5), for a suitable $\delta>0$ which will be determined later. In the following, we denote $\bar{w}:=\frac{1}{2}|x|^{2}+w$ and $\bar{m}:=\|\nabla \bar{w}\|_{L^{\infty}([\Omega])}$. Observe that $\bar{w}$ is the unique convex function satisfying

$$
\left\{\begin{array}{rlll}
\operatorname{det} D^{2} \bar{w} & =1 & \text { in } & V, \\
\bar{w} & =0 & \text { on } & \partial V .
\end{array}\right.
$$

By the hypothesis for $f$ and $g$, we have that

$$
\begin{equation*}
|f(x)-1| \leq \theta \tag{6.10}
\end{equation*}
$$

and

$$
\begin{equation*}
|g(x+\nabla w \pm \delta x+G(\nabla w \pm \delta x))-1| \leq \theta(2 \bar{m}+2 n)^{\alpha} . \tag{6.11}
\end{equation*}
$$

We note that to obtain the estimate (6.11), we have also used the fact $C \bar{m} / \lambda \leq 1$ which is a consequence of the estimate for $\bar{m}$ given in the Step 3 below and our assumption (6.8).
Step 1. Construction of a supersolution. We show first that $w^{+}$is in the admissible class for the operator $M_{\lambda}$ in $\Omega$ (see Definition 6.1) and that the matrix

$$
A^{+}:=I+D^{2} c_{\lambda}^{*}\left(-\nabla w^{+}\right) D^{2} w^{+}=I+D^{2} w-\delta I+\left(D^{2} w-\delta I\right) H(\nabla w-\delta x)
$$

is positive definite. Indeed, we can write

$$
\begin{aligned}
A^{+} & =D^{2} \bar{w}[I+H]-\delta I-(1+\delta) H \\
& \geq \frac{1}{2} D^{2} \bar{w}-\delta I-2|H| I \\
& =\left[\frac{1}{4} D^{2} \bar{w}-\delta I\right]+\left[\frac{1}{4} D^{2} \bar{w}-2|H| I\right] .
\end{aligned}
$$

Moreover, it follows from the hypothesis (2.11) for $H$ that

$$
\begin{equation*}
|H| \leq \frac{C(\bar{m}+2 n)}{\lambda} \leq \frac{\delta}{4} \tag{6.12}
\end{equation*}
$$

for a suitable choice of $\delta$. Therefore,

$$
\begin{equation*}
A^{+} \geq 2\left[\frac{1}{4} D^{2} \bar{w}-\delta I\right] \tag{6.13}
\end{equation*}
$$

Since $\operatorname{det}\left(D^{2} \bar{w}\right)=1$, an upper bound for the eigenvalues of $D^{2} \bar{w}$, given in Proposition 6.6, implies a lower bound. In fact,

$$
\begin{equation*}
D^{2} \bar{w}(x) \geq\left(\frac{1}{C_{n}^{\prime} \operatorname{dist}(x, \partial V)^{1-n}}\right)^{n-1} I \geq \frac{1}{C_{n}^{\prime}}\left(\frac{d}{\lambda}\right)^{\beta(n-1)^{2}} I>4 \delta I \quad \text { for all } x \in[\Omega] \tag{6.14}
\end{equation*}
$$

For this choice of $\delta$, looking at (6.13), we obtain that $A^{+}$is positive definite in $\Omega$. Furthermore, (6.14) together with the second inequality in (6.12) gives

$$
\inf _{[\Omega]} D^{2} \bar{w}>\delta I+\frac{C}{\lambda} I\left(\bar{m}+\delta\|i d\|_{L^{\infty}([\Omega])}+2 n\right)
$$

implying that $w^{+}$satisfies (6.1). Therefore, $w^{+}$is in the admissible class for $M_{\lambda}$ in $\Omega$.
Now we check that $w^{+}$is indeed a supersolution. It remains to show that

$$
\begin{equation*}
M_{\lambda}\left[w^{+}\right]<f \tag{6.15}
\end{equation*}
$$

for

$$
M_{\lambda}\left[w^{+}\right]=g(x+\nabla w-\delta x+G(\nabla w-\delta x)) \operatorname{det}\left(A^{+}\right) .
$$

We compute

$$
\begin{equation*}
(I+H)\left(I+D^{2} w\right)=A^{+}+\frac{\delta}{2} I+\left[\frac{\delta}{2} I+(1+\delta) H\right] . \tag{6.16}
\end{equation*}
$$

Because $|H| \leq \frac{\delta}{4}$ by the choice of $\delta$ from (6.12), we have that the term $\frac{\delta}{2} I+(1+\delta) H$ is positive definite. On the other hand, given any two positive definite matrices $A_{1}, A_{2}$,

$$
\operatorname{det}^{\frac{1}{n}}\left(A_{1}+A_{2}\right) \geq \operatorname{det}^{\frac{1}{n}}\left(A_{1}\right)+\operatorname{det}^{\frac{1}{n}}\left(A_{2}\right) .
$$

Consequently, from (6.16) we obtain that

$$
1+\frac{\delta}{4} \geq \operatorname{det}^{\frac{1}{n}}(I+H) \geq \operatorname{det}^{\frac{1}{n}}\left(A^{+}\right)+\frac{\delta}{2},
$$

yielding

$$
\operatorname{det}\left(A^{+}\right) \leq\left(1-\frac{\delta}{4}\right)^{n} \leq 1-\frac{n}{8} \delta
$$

Finally, using the bounds for $f, g$ from (6.10) and (6.11),

$$
\begin{aligned}
M_{\lambda}\left[w^{+}\right] & =g(x+\nabla w-\delta x+G(\nabla w-\delta x)) \operatorname{det}\left(A^{+}\right) \\
& \leq\left(1+\theta(2 \bar{m}+2 n)^{\alpha}\right)\left(1-\frac{n}{8} \delta\right)<1-\theta \leq f
\end{aligned}
$$

given a suitable choice of

$$
\begin{equation*}
\delta \geq \frac{16 \theta}{n}(2 \bar{m}+2 n)^{\alpha} . \tag{6.17}
\end{equation*}
$$

Then (6.15) follows as desired.
Step 2. Construction of a subsolution. We now check that $w^{-}$is a subsolution. As $w^{-}$is obviously in the admissible class for the operator $M_{\lambda}$ in $\Omega$, we only need to show $M_{\lambda}\left[w^{-}\right]>f$. We compute

$$
M_{\lambda}\left[w^{-}\right]=g(x+\nabla w+\delta x+G(\nabla w+\delta x)) \operatorname{det}\left(A^{-}\right)
$$

for

$$
A^{-}:=I+D^{2} w+\delta I+\left(D^{2} w+\delta I\right) H(\nabla w+\delta x)
$$

If we choose $\delta$ satisfy (6.12), then we know that the following matrices $D^{2} \bar{w}, \frac{\delta}{4} I+H$ and $\frac{\delta}{2} I-H+\delta H$ are positive definite because of the estimate $|H(\nabla w+\delta x)| \leq \frac{C(\bar{m}+2 n)}{\lambda} \leq \frac{\delta}{4}$. Consequently,

$$
A^{-}=\left(1-\frac{\delta}{4}\right) D^{2} \bar{w}+\left(\frac{\delta}{4} I+H\right) D^{2} \bar{w}+\frac{\delta}{2} I+\left[\frac{\delta}{2} I-H+\delta H\right] \geq\left(1-\frac{\delta}{4}\right) D^{2} \bar{w}+\frac{\delta}{2} I .
$$

Thus the matrix $A^{-}$is positive definite. Moreover by using the bounds for $f, g$ from (6.10) and (6.11) we obtain the desired conclusion

$$
M_{\lambda}\left[w^{-}\right] \geq\left(1-\theta(2 \bar{m}+2 n)^{\alpha}\right)\left(1+\frac{\delta}{4}\right)^{n} \geq\left(1-\theta(2 \bar{m}+2 n)^{\alpha}\right)\left(1+\frac{n \delta}{4}\right)>1+\theta \geq f
$$

where $\delta$ is chosen as in (6.17). Note that in the above estimate for $w^{-}$we actually do not need a uniform estimate of $D^{2} \bar{w}$ as in the case of $w^{+}$.

## Step 3. Choice of $\delta$.

We now check that all the conditions on $\delta$ given in (6.12), (6.14) and (6.17) are compatible. Let us take

$$
\begin{equation*}
\delta:=\frac{1}{4 C_{n}^{\prime}}\left(\frac{d}{\lambda}\right)^{\beta(n-1)^{2}} \tag{6.18}
\end{equation*}
$$

which guarantees (6.14). To see that this choice of $\delta$ also satisfies (6.12) and (6.17), we first observe that Proposition 6.5 implies

$$
\bar{m} \leq C_{n} \operatorname{dist}([\Omega], \partial V)^{\frac{2}{n}-1} \leq C_{n}\left(\frac{d}{\lambda}\right)^{\frac{2-n}{n} \beta}
$$

Therefore, both (6.12) and (6.17) hold provided that

$$
C_{n} \frac{1}{\lambda}\left(\frac{d}{\lambda}\right)^{\frac{2-n}{n} \beta} \leq\left(\frac{d}{\lambda}\right)^{\beta(n-1)^{2}} \quad \text { and } \quad C_{n} \theta\left(\frac{d}{\lambda}\right)^{\frac{2-n}{n} \alpha \beta} \leq\left(\frac{d}{\lambda}\right)^{\beta(n-1)^{2}}
$$

But these are true thanks to the hypothesis (6.8).
Step 4. Completion of the proof. The proof of the theorem is completed using the comparison principle for Alexandroff solutions of Proposition 6.3 and the initial hypothesis on $u$, which guarantee that

$$
\begin{equation*}
(u-w)(x) \leq \max _{\partial \Omega}\left\{u-w^{+}\right\}-\frac{\delta}{2}|x|^{2}, \quad \text { for all } \quad x \in \Omega, \tag{6.19}
\end{equation*}
$$

and

$$
\begin{equation*}
\min _{\partial \Omega}\left\{u-w^{-}\right\}+\frac{\delta}{2}|x|^{2} \leq(u-w)(x), \quad \text { for all } \quad x \in \Omega \tag{6.20}
\end{equation*}
$$

Let us estimate the boundary values for $u-w^{+}$and $u-w^{-}$. Recall that $\bar{w}(x)=\frac{1}{2}|x|^{2}+$ $w(x)$ and $u(x)=-c_{\lambda}(x)=-\frac{1}{2}|x|^{2}-\frac{1}{\lambda} E(x)$ on $\partial \Omega$. Hence

$$
u(x)-w^{+}(x)=u(x)-w(x)+\frac{\delta}{2}|x|^{2}=-\bar{w}(x)-\frac{1}{\lambda} E(x)+\frac{\delta}{2}|x|^{2} \quad \text { on } \quad \partial \Omega .
$$

We are going to use Proposition 6.5 to estimate the behavior of $\bar{w}$ near $\partial V$. In particular, for $x \in \partial \Omega$,

$$
\bar{w}(x) \geq-C \operatorname{dist}(x, \partial V)^{\frac{2}{n}} \geq-C\left(\frac{d}{\lambda}+\left(\frac{d}{\lambda}\right)^{\beta}\right)^{\frac{2}{n}} \geq-C\left(\frac{d}{\lambda}\right)^{\frac{2 \beta}{n}}
$$

where we have used the assumption (6.4) and the definition of $V$ in the second inequality. It follows that

$$
\begin{equation*}
u(x)-w^{+}(x) \leq C\left(\frac{d}{\lambda}\right)^{\frac{2 \beta}{n}}+\frac{C}{\lambda}+C \delta \leq C\left(\frac{d}{\lambda}\right)^{\frac{2 \beta}{n}} \quad \text { for } x \in \partial \Omega \tag{6.21}
\end{equation*}
$$

by our choice of $\delta$ in (6.18). From (6.19) and (6.21) we obtain the second inequality in (6.9). To estimate the values of $u-w^{-}$on $\partial \Omega$, observe that $\bar{w} \leq 0$. Hence, we have for $x \in \partial \Omega$,

$$
u(x)-w^{-}(x)=u(x)-w(x)-\frac{\delta}{2}|x|^{2}=-\bar{w}(x)-\frac{1}{\lambda} E(x)-\frac{\delta}{2}|x|^{2} \geq-\frac{C}{\lambda}-C \frac{\delta}{2} \geq-C \delta .
$$

This together with (6.20) yields the first inequality in (6.9) and the theorem is proved.

Remark 6.7. In the proof of Theorem 6.4, it might happen that $x-\nabla c_{\lambda}^{*}\left(-\nabla w^{ \pm}(x)\right)$ does not belong to the domain of $g$ for some $x$ in $\Omega$. Therefore, $M_{\lambda}\left[w^{ \pm}\right]$is not well defined at those $x$. However we do not have any problem since to apply Proposition 6.3 in Step 4 above we only need to check $M_{\lambda}\left[w^{+}\right]<f$ and $M_{\lambda}\left[w^{-}\right]>f$ at points $x$ where $x-\nabla c_{\lambda}^{*}\left(-\nabla w^{ \pm}(x)\right) \in \partial_{c_{\lambda}} u(\Omega)$.

### 6.2 Sections are round

Here we compare sections $S_{\mu}\left(u, x_{0}\right)$ of $u$ to those of a solution of a Monge-Ampère equation. Let $[\Omega]$ be the convex hull of $\Omega$ and $w$ be the solution of

$$
\left\{\begin{align*}
\operatorname{det}\left(I+D^{2} w\right) & =1 & & \text { in }[\Omega],  \tag{6.22}\\
w & =-\frac{1}{2}\left|x-y_{0}\right|^{2} & & \text { on } \partial[\Omega],
\end{align*}\right.
$$

which is given by $\bar{w}=w+\frac{1}{2}\left|x-y_{0}\right|^{2}$ where $\bar{w}$ is the unique convex solution of $\operatorname{det} D^{2} \bar{w}=1$ in $[\Omega]$ and $\bar{w}=0$ on $\partial[\Omega]$.

Note that in Theorem 5.1 we already have a preliminary control of the sections of $u$. However, we would like to obtain a more precise estimate using the comparison principle from Theorem 6.4, that allows us to say that $w$ is a good approximation for $u$, and this approximation may be quantified explicitly. In the following three lemmas, we compare sections of $u$ and $w$ by assuming in advance that $\|u-w\|$ is small. We remind the reader that the $\mu$-section of $u$ was defined in (2.12) and (2.13), while the sections for $w$ are defined in (4.13) and (4.14).

Lemma 6.8. Let $B_{1} \subset \Omega \subset[\Omega] \subset B_{n}$ be a domain. Let $u \in \mathcal{C}(\bar{\Omega})$ be a $c_{\lambda}$-convex function and $y_{0} \in \partial_{c} u(\Omega)$. Consider $w$ be the solution of (6.22), and $\bar{u}, \bar{w}$ be defined as above. Assume that we have proved that

$$
\begin{equation*}
\|\bar{u}-\bar{w}\|_{L^{\infty}(\Omega)} \leq \epsilon \tag{6.23}
\end{equation*}
$$

for some $\epsilon>0$. There exist some dimensional constants $\delta_{n}>0, \mu_{0}>0, \tau_{0}>0$ and $C>0$ such that if

$$
\begin{equation*}
\operatorname{dist}(x, \partial[\Omega]) \leq \frac{\delta_{n}}{2} \text { for all } x \in \partial \Omega \tag{6.24}
\end{equation*}
$$

and $x_{0} \in \Omega$ is a minimum point of $\bar{u}$ in $\Omega$, then we have

$$
\begin{equation*}
T_{\mu-C \epsilon^{1 / 2}}\left(w, x_{0}\right) \subset S_{\mu}\left(u, x_{0}\right) \subset T_{\mu+C \epsilon^{1 / 2}}\left(w, x_{0}\right) \subset \Omega_{\frac{\delta_{n}}{4}} \tag{6.25}
\end{equation*}
$$

for all $\mu \leq \mu_{0}, \epsilon \leq \tau_{0} \mu^{2}$.
Proof. Recall that $\bar{w}$ is the convex function satisfying $\operatorname{det} D^{2} \bar{w}=1$ in $[\Omega]$ and $\bar{w}=0$ on $\partial[\Omega]$. Let $\bar{x} \in[\Omega]$ be the point where $\bar{w}$ attains its minimum value in [ $\Omega$ ]; it exists because $\bar{w}$ is a strictly convex function (see [6]). We would like to say that $\bar{x}$ and $x_{0}$ stay away from the boundary of $\Omega$. Indeed, [26, Proposition 3.2.3] gives an estimate for the value of $\bar{w}$ at its
minimum, $|\bar{w}(\bar{x})| \approx C_{n}, C_{n}$ dimensional constant, and from Alexandroff's maximum principle [26, Theorem 1.4.2], we get that $\operatorname{dist}(\bar{x}, \partial[\Omega]) \geq \delta_{n}, \delta_{n}$ universal constant. It then follows from the assumption (6.24) that $\operatorname{dist}(\bar{x}, \partial \Omega) \geq \frac{\delta_{n}}{2}$. Consequently, we obtain from (6.23) that $\bar{w}(\bar{x})-\epsilon \leq \bar{w}\left(x_{0}\right)-\epsilon \leq \bar{u}\left(x_{0}\right) \leq \bar{u}(\bar{x}) \leq \bar{w}(\bar{x})+\epsilon$, i.e.,

$$
\begin{equation*}
\left|\bar{u}\left(x_{0}\right)-\bar{w}(\bar{x})\right| \leq \epsilon . \tag{6.26}
\end{equation*}
$$

Since $\bar{w}\left(x_{0}\right) \leq \bar{w}(\bar{x})+2 \epsilon \leq-C_{n}$ if $\epsilon$ is small enough, we conclude from Alexandroff's maximum principle that $\operatorname{dist}\left(x_{0}, \partial[\Omega]\right) \geq \delta_{n}$. In particular, $\operatorname{dist}\left(x_{0}, \partial \Omega\right) \geq \frac{\delta_{n}}{2}$ by (6.24).

In $[\Omega]_{\delta_{n}}:=\left\{x \in[\Omega]: \operatorname{dist}(x, \partial[\Omega]) \geq \delta_{n}\right\}$, we can use Pogorelov's estimates [26, formula (4.2.6)], to get that

$$
\begin{equation*}
\frac{2}{C_{2}^{2}} I \leq D^{2} \bar{w}(x) \leq \frac{2}{C_{1}^{2}} I \quad \text { for all } \quad x \in[\Omega]_{\delta_{n}} . \tag{6.27}
\end{equation*}
$$

With these in mind, it is known (see [27], for instance) that

$$
\begin{equation*}
\left|\nabla \bar{w}\left(x_{0}\right)\right| \leq C \epsilon^{1 / 2} . \tag{6.28}
\end{equation*}
$$

We present the proof of (6.28) here for the sake of clarity. First of all, Taylor estimates give

$$
\begin{equation*}
\bar{w}\left(x_{0}\right)-\bar{u}\left(x_{0}\right)=\bar{w}(\bar{x})-\bar{u}\left(x_{0}\right)+\frac{1}{2}\left\langle D^{2} w(\xi)\left(x_{0}-\bar{x}\right), x_{0}-\bar{x}\right\rangle \tag{6.29}
\end{equation*}
$$

for some $\xi$ between $x_{0}$ and $\bar{x}$. Therefore, since $x_{0}, \bar{x} \in[\Omega]_{\delta_{n}}$, we can use (6.26) and Pogorelov estimates (6.27) in (6.29) to obtain

$$
\begin{equation*}
\left|\bar{x}-x_{0}\right| \leq C \epsilon^{1 / 2} \tag{6.30}
\end{equation*}
$$

On the other hand,

$$
D_{i} \bar{w}\left(x_{0}\right)-D_{i} \bar{w}(\bar{x})=-\int_{0}^{1} D\left(D_{i} \bar{w}\right)\left(x_{0}+t\left(\bar{x}-x_{0}\right)\right) \cdot\left(\bar{x}-x_{0}\right) d t
$$

thus using (6.30) we obtain the desired (6.28).
Now, we would like to say that the sections of $w$ do not approach the boundary of $\Omega$. Using Alexandroff principle again we know that $T_{\mu}(w, \bar{x}) \subset \Omega_{\frac{\delta_{n}}{4}}$ for $0<\mu<C_{n}$. Then, because of (6.23), (6.26) and (6.28), there exist $\tau_{0}, \mu_{0}$ universal constants such that

$$
T_{\mu}\left(w, x_{0}\right) \subset T_{C \epsilon^{1 / 2}+\mu}(w, \bar{x}) \subset \Omega_{\frac{\delta_{n}}{4}}
$$

for $\epsilon \leq \tau_{0} \mu^{2}$ and $\mu \leq \mu_{0}$.
Next, we compare the sections of $u$ and $w$. Let $x \in S_{\mu}\left(u, x_{0}\right)$. Then, using (2.13), (6.23) and (6.28), it follows,

$$
\begin{aligned}
\bar{w}(x) & \leq \bar{u}(x)+\epsilon \leq \bar{u}\left(x_{0}\right)+\mu+\epsilon \\
& \leq \bar{w}\left(x_{0}\right)+\left\langle\nabla \bar{w}\left(x_{0}\right), x-x_{0}\right\rangle+\mu+2 \epsilon-\left\langle\nabla \bar{w}\left(x_{0}\right), x-x_{0}\right\rangle \\
& \leq \bar{w}\left(x_{0}\right)+\left\langle\nabla \bar{w}\left(x_{0}\right), x-x_{0}\right\rangle+\mu+2 \epsilon+2 C \epsilon^{1 / 2} \\
& \leq \bar{w}\left(x_{0}\right)+\left\langle\nabla \bar{w}\left(x_{0}\right), x-x_{0}\right\rangle+\mu+C \epsilon^{1 / 2} .
\end{aligned}
$$

Hence, $x \in T_{\mu+C \epsilon^{1 / 2}}(\bar{w})$ and the second inclusion of (6.25) is proved. The other inclusion is very similar.

We can also show that $u$ is close to a quadratic polynomial:
Lemma 6.9. In the hypothesis of the previous lemma, there exist a positive definite matrix $M=A^{t} A$ and a vector $p \in \mathbb{R}^{n}$ satisfying

$$
\operatorname{det} M=1, \quad 0<C_{1} I \leq M \leq C_{2} I, \quad \text { and } \quad|p| \leq C \epsilon^{1 / 2}
$$

such that we have

$$
\begin{equation*}
B_{\left(1-C\left(\mu^{1 / 2}+\mu^{-1} \epsilon^{1 / 2}\right)\right) \sqrt{2}} \subset \mu^{-1 / 2} T S_{\mu}\left(u, x_{0}\right) \subset B_{\left(1+C\left(\mu^{1 / 2}+\mu^{-1} \epsilon^{1 / 2}\right)\right) \sqrt{2}} \tag{6.31}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\bar{u}(x)-\left(\bar{u}\left(x_{0}\right)+\left\langle p, x-x_{0}\right\rangle+\frac{1}{2}\left\langle M\left(x-x_{0}\right), x-x_{0}\right\rangle\right)\right| \leq C\left(\mu^{3 / 2}+\epsilon\right) \tag{6.32}
\end{equation*}
$$

for $x \in S_{\mu}\left(u, x_{0}\right)$, where $T x=A\left(x-x_{0}\right)$ and $\bar{u}(x):=u(x)+c\left(x-y_{0}\right)$. Or equivalently for $u$, $\left|u(x)-\left(u\left(x_{0}\right)+\left\langle p-x_{0}+y_{0}, x-x_{0}\right\rangle+\frac{1}{2}\left\langle(M-I)\left(x-x_{0}\right), x-x_{0}\right\rangle\right)\right| \leq C\left(\mu^{3 / 2}+\frac{\mu^{1 / 2}}{\lambda}+\epsilon\right)$.
Proof. The behavior of the sections of the function $w$ follows from the following three claims (that are proven in [27]).

1. There exist $\tau_{0}$ and $\mu_{0}, \delta_{n}$, dimensional constants, such that, for $\epsilon \leq \tau_{0} \mu^{2}$ and $\mu \leq \mu_{0}$,

$$
T_{\mu}\left(w, x_{0}\right) \subset[\Omega]_{\delta_{n}} .
$$

2. There exists $\mu_{n}$ such that if $\mu \leq \mu_{n}$ and $\gamma \leq \frac{3}{4} \mu$, then

$$
\begin{equation*}
\partial T_{\mu+\gamma}\left(w, x_{0}\right) \subset \mathcal{N}_{\frac{C \gamma}{\sqrt{\mu}}}\left(\partial T_{\mu}\left(w, x_{0}\right)\right) \quad \text { and } \quad \partial T_{\mu-\gamma}\left(w, x_{0}\right) \subset \mathcal{N}_{\frac{C \gamma}{\sqrt{\mu}}}\left(\partial T_{\mu}\left(w, x_{0}\right)\right) \tag{6.33}
\end{equation*}
$$

3. Let $M=D^{2} \bar{w}\left(x_{0}\right)$ and $E=\left\{x: \frac{1}{2}\left\langle D^{2} \bar{w}\left(x_{0}\right)\left(x-x_{0}\right), x-x_{0}\right\rangle \leq 1\right\}$. We compare $T_{\mu}\left(w, x_{0}\right)$ with ellipsoids and claim that

$$
\begin{equation*}
\partial T_{\mu}\left(w, x_{0}\right) \subset \mathcal{N}_{C \mu}\left(\partial \mu^{1 / 2} E\right) \tag{6.34}
\end{equation*}
$$

for some structural constant $C$ and all $0<\mu \leq \mu_{n}$. Here the dilation is with respect to the point $x_{0}$.

We set $\mu_{0}:=\min \left\{\mu_{0}, \mu_{n}\right\}$, and let $\mu \leq \mu_{0}, \epsilon \leq \tau_{0} \mu^{2}$. From (6.25) and (6.33), we obtain

$$
\begin{equation*}
\partial S_{\mu}\left(u, x_{0}\right) \subset \mathcal{N}_{C \mu^{-1 / 2} \epsilon^{1 / 2}}\left(\partial T_{\mu}\left(w, x_{0}\right)\right) \tag{6.35}
\end{equation*}
$$

which together with (6.34) gives

$$
\begin{equation*}
\partial S_{\mu}\left(u, x_{0}\right) \subset \mathcal{N}_{C\left(\mu+\mu^{-1 / 2} \epsilon^{1 / 2}\right)}\left(\partial \mu^{1 / 2} E\right) \tag{6.36}
\end{equation*}
$$

We also notice that

$$
\begin{equation*}
\partial \mu^{1 / 2} E=x_{0}+A^{-1}\left(\partial B_{\sqrt{2 \mu}}\right) \tag{6.37}
\end{equation*}
$$

where $M=A^{t} A$. If we let $T x=A\left(x-x_{0}\right)$, then we have that

$$
\begin{equation*}
T\left(\mathcal{N}_{\delta}\left(\partial \mu^{1 / 2} E\right)\right) \subset B_{\sqrt{2 \mu}+\|A\| \delta} \tag{6.38}
\end{equation*}
$$

and

$$
T\left(\mathcal{N}_{\delta}\left(\partial \mu^{1 / 2} E\right)\right) \subset\left(B_{\sqrt{2 \mu}-\|A\| \delta}\right)^{c}
$$

Hence taking $\delta=C\left(\mu+\mu^{-1 / 2} \epsilon^{1 / 2}\right)$ from (6.36) we get (6.31).
And for the second part of the lemma, using (6.31), the fact that $\|A\|$ is bounded, and that $\epsilon \leq \tau_{0} \mu^{2}$, we obtain $S_{\mu}\left(u, x_{0}\right) \subset B_{C \mu^{1 / 2}}\left(x_{0}\right)$. Then, taking $p=\nabla \bar{w}\left(x_{0}\right)$ yields

$$
\begin{aligned}
& \left|\bar{u}(x)-\left(\bar{u}\left(x_{0}\right)+\left\langle p, x-x_{0}\right\rangle+\frac{1}{2}\left\langle M\left(x-x_{0}\right), x-x_{0}\right\rangle\right)\right| \\
& \quad \leq|\bar{u}(x)-\bar{w}(x)|+\left|\bar{w}(x)-\bar{w}\left(x_{0}\right)-\left\langle p, x-x_{0}\right\rangle-\frac{1}{2}\left\langle M\left(x-x_{0}\right), x-x_{0}\right\rangle\right|+\left|\bar{w}\left(x_{0}\right)-\bar{u}\left(x_{0}\right)\right| \\
& \quad \leq 2 \epsilon+\left\|D^{3} \bar{w}\right\|_{L^{\infty}\left(B_{\left.C \mu^{1 / 2}\right)}\right.}\left|x-x_{0}\right|^{3} \leq C\left(\epsilon+\mu^{3 / 2}\right),
\end{aligned}
$$

as desired. Notice that $|p| \leq C \epsilon^{1 / 2}$ by (6.28).
If the boundary of the convex hull of $\Omega$ is close to that of $B_{\sqrt{2}}$, then we can get better estimates for $M$ and $\partial S_{\mu}$ :

Lemma 6.10. In addition to the hypothesis of Lemma 6.8, suppose that $\Omega$ is a domain satisfying

$$
B_{(1-\sigma) \sqrt{2}} \subset[\Omega] \subset B_{(1+\sigma) \sqrt{2}}
$$

for some $0<\sigma \leq 1$. Then, there exist $\mu_{0}>0, \tau_{0}>0$ which are independent of $\sigma$, a positive definite matrix $M=A^{t} A$, and $p \in \mathbb{R}^{n}$ with

$$
\operatorname{det} M=1, \quad(1-C \sigma)|\xi|^{2} \leq\langle M \xi, \xi\rangle \leq(1+C \sigma)|\xi|^{2}, \quad|p| \leq C \epsilon^{1 / 2} \text { and }\left|p-x_{0}\right| \leq C \sigma \quad \text { (6.39) }
$$

such that for $0<\mu \leq \mu_{0}$ and $\epsilon \leq \tau_{0} \mu^{2}$ we have

$$
\begin{equation*}
B_{\left(1-C\left(\sigma \mu^{1 / 2}+\mu^{-1} \epsilon^{1 / 2}\right)\right) \sqrt{2}} \subset \mu^{-1 / 2} T S_{\mu}\left(u, x_{0}\right) \subset B_{\left(1+C\left(\sigma \mu^{1 / 2}+\mu^{-1} \epsilon^{1 / 2}\right)\right) \sqrt{2}}, \tag{6.40}
\end{equation*}
$$

and in $S_{\mu}\left(u, x_{0}\right)$,

$$
\begin{equation*}
\left|\bar{u}(x)-\left(\bar{u}\left(x_{0}\right)+\left\langle p, x-x_{0}\right\rangle+\frac{1}{2}\left\langle M\left(x-x_{0}\right), x-x_{0}\right\rangle\right)\right| \leq C\left(\sigma \mu^{3 / 2}+\epsilon\right), \tag{6.41}
\end{equation*}
$$

where $T x=A\left(x-x_{0}\right)$ and $\bar{u}(x):=u(x)+c\left(x-y_{0}\right)$. Or equivalently,
$\left|u(x)-\left(u\left(x_{0}\right)+\left\langle p-x_{0}+y_{0}, x-x_{0}\right\rangle+\frac{1}{2}\left\langle(M-I)\left(x-x_{0}\right), x-x_{0}\right\rangle\right)\right| \leq C\left(\sigma \mu^{3 / 2}+\frac{\mu^{1 / 2}}{\lambda}+\epsilon\right)$ in $S_{\mu}\left(u, x_{0}\right)$.
Proof. It is similar to that of Lemma 6.9. The only difference is that since now $\partial[\Omega] \subset$ $\mathcal{N}_{\sigma \sqrt{2}}\left(\partial B_{\sqrt{2}}\right)$, we will get the following improvement of (6.34):

$$
\begin{equation*}
\partial T_{\mu}\left(w, x_{0}\right) \subset \mathcal{N}_{C \sigma \mu}\left(\partial \mu^{1 / 2} E\right) \tag{6.42}
\end{equation*}
$$

To do this, let $P(x)=\frac{1}{2}|x|^{2}-1$, then $\operatorname{det} D^{2} P(x)=1$. Since $\partial[\Omega] \subset \mathcal{N}_{\sigma \sqrt{2}}\left(\partial B_{\sqrt{2}}\right)$, we have $P-3 \sigma \leq 0 \leq P+3 \sigma$ on $\partial[\Omega]$ and by the comparison principle we get $\|\bar{w}-P\|_{L^{\infty}([\Omega])} \leq 3 \sigma$. Since $\bar{w}$ is smooth, $v:=\bar{w}-P$ satisfies the following uniformly elliptic linear equation

$$
\operatorname{trace}\left(D(x) D^{2} v\right)=0
$$

where $D(x)=\int_{0}^{1}\left(t D^{2} \bar{w}(x)+(1-t) I\right)^{-1} \operatorname{det}\left(t D^{2} \bar{w}(x)+(1-t) I\right) d t$. Therefore, by interior Schauder estimates

$$
\|\bar{w}-P\|_{C_{\mathrm{loc}}^{2}} \leq C\|\bar{w}-P\|_{L_{\mathrm{loc}}^{\infty}} \leq C \sigma
$$

In particular, $\left|p-x_{0}\right|=\left|D \bar{w}\left(x_{0}\right)-D P\left(x_{0}\right)\right| \leq C \sigma$ and $\left|D_{i j} \bar{w}\left(x_{0}\right)-\delta_{i j}\right| \leq C \sigma$ for $1 \leq i, j \leq n$. By differentiating det $D^{2} \bar{w}(x)=1$, we obtain that the function $v=D(\bar{w}-P)$ satisfies the linearized equation

$$
\operatorname{trace}\left(\left(D^{2} \bar{w}(x)\right)^{-1} D^{2} v\right)=0
$$

Again by the interior Schauder estimates $\left(\mathcal{C}_{\text {loc }}^{2}\right.$ estimates) we get

$$
\begin{equation*}
\left\|D^{3} \bar{w}\right\|_{L_{\text {loc }}^{\infty}}=\left\|D^{3}(\bar{w}-P)\right\|_{L_{\text {loc }}^{\infty}} \leq C\|D(\bar{w}-P)\|_{L_{\text {loc }}^{\infty}} \leq C^{*} \sigma \tag{6.43}
\end{equation*}
$$

The rest of the proof is standard and follows as in [27, Lemma 1.2].

## 7 Interior $\mathcal{C}^{2, \alpha}$ estimates

Here we present the proof of Theorem 1.2 . Let $u$ be a solution of (1.1), or equivalently (thanks to the results in Subsection 2.2), $u$ is an Alexandroff weak solution of

$$
\left\{\begin{align*}
M_{\lambda}[u] & =f \quad \text { in } \Omega  \tag{7.1}\\
\partial_{c_{\lambda}} u(\Omega) & =\Omega^{\prime}
\end{align*}\right.
$$

for some $\lambda>0$. Fix $\delta_{0}>0$, and let

$$
\Omega_{\delta_{0}}:=\left\{x \in \Omega: \operatorname{dist}(x, \partial \Omega)>\delta_{0}\right\}
$$

We will find $\lambda_{0}$ and $r>0$ depending on $\delta_{0}$ and $\alpha$ such that for all $\lambda \geq \lambda_{0}$, and for every point $x_{0} \in \Omega_{\delta_{0}}$ it is possible to construct a second order polynomial in $x$, call it $P_{0}$, such that

$$
\begin{equation*}
\left|u(x)-P_{0}(x)\right| \leq C\left|x-x_{0}\right|^{\alpha^{\prime}} \quad \forall x \in B_{r}\left(x_{0}\right) \tag{7.2}
\end{equation*}
$$

for some $\alpha^{\prime} \in(0,1)$ depending only on $n, \alpha$ and chosen later, and $C$ depending only on $\lambda_{0}$, $n, \Lambda_{1}, \Lambda_{2}, \alpha, \delta_{0},\|f\|_{\mathcal{C}^{\alpha}\left(\overline{\Omega_{\delta_{1}}}\right)},\|g\|_{\mathcal{C}^{\alpha}\left(\overline{\Omega_{\delta_{1}}^{\prime}}\right)}, \Omega$ and $\Omega^{\prime}$. The constant $\delta_{1}>0$ will be made precise in the proof and depends only on $\delta_{0}$ and universal constants.

The proof goes by induction. In the first step, we will compare the function $u$ with the solution of an optimal transport problem with quadratic cost. In the successive steps, once we have localized the problem, the comparison will be with respect to a solution of a Monge-Ampère equation with homogeneous Dirichlet boundary data.

### 7.1 Localization

We first localize the problem in a section so that we have control of the boundary data. Fix any $x_{0} \in \Omega_{\delta_{0}}$ and $y_{0} \in \partial_{c} u\left(x_{0}\right)$. The notion of sections for $u$ at $x_{0}$ was defined in (2.12), or equivalently for $\bar{u}(x):=u(x)+c\left(x-y_{0}\right)$, as in (2.13). Note that $x_{0}$ is a minimum point for $\bar{u}$ in $\Omega$.

We make now one further simplification. It is enough to assume that $x_{0}=y_{0}$, i.e., that $p_{0}:=-\nabla c\left(x_{0}-y_{0}\right)=0$. This can be achieved by subtracting a linear function to the potential $u$.

This first step in the induction argument was considered in Section 4. In fact, Proposition 4.1 shows that $u$ is actually close in the $L^{\infty}$ norm to a strictly convex function $w$. It
also implies that the sections of $u$ stay away from the boundary $\partial \Omega$. This is done in Proposition 4.5 and Corollary 4.6. These imply, in particular, that there exists a positive constant $\mu_{1}=\mu_{1}\left(\delta_{0}, n, \Lambda_{1}, \Lambda_{2}, \Omega, \Omega^{\prime}\right)$ satisfying: for all $h \leq \mu_{1}$ we can find $\lambda_{0}>0$ depending on $h$ such that for all $\lambda \geq \lambda_{0}$, if $u$ is the solution of (7.1) for the cost $c_{\lambda}$, then

$$
S_{h}\left(u, x_{0}\right):=S_{h}\left(u, x_{0}, y_{0}\right) \subset \Omega_{c \delta_{0}^{\theta}} .
$$

As we assume

$$
\begin{equation*}
\partial_{c}^{\mathrm{loc}} u(x)=\partial_{c} u(x) \quad \text { for all } \quad x \in \Omega_{c \delta_{0}^{\theta}}, \tag{7.3}
\end{equation*}
$$

the restriction of $u$ to any subregion of $\Omega_{c \delta_{0}^{\theta}}$ is still an Alexandroff weak solution in that subregion of the same equation $M_{\lambda}[u]=f$. Next look at the image set. Note that Proposition 4.4 gives that $\partial_{c} u\left(\Omega_{c \delta_{0}^{\theta}}\right) \subset \Omega_{\delta_{1}}^{\prime}$ for some small $\delta_{1}>0$ depending only on $\delta_{0}$ and universal constants. Since $\Omega^{\prime}$ is strongly convex by the assumption, so is $\Omega_{\delta_{1}}^{\prime}$ and hence as in Lemma 2.4 we conclude that $\Omega_{\delta_{1}}^{\prime}$ is $c_{\lambda}^{*}$-convex with respect to $\Omega$. It then follows from [35] that $\partial_{c} u\left(\Omega_{c \delta_{0}^{\theta}}\right) \subset \overline{\Omega_{\delta_{1}}^{\prime}}$ which allows to restrict the dependence on $g$ in the arguments to $\|g\|_{\mathcal{C}^{\alpha}\left(\overline{\Omega_{\delta_{1}}^{\prime}}\right)}$, as stated in the theorem.

Next, Corollary 4.6 and its subsequent remark (see formulas (4.18) and (4.19)) show that $S_{h}\left(u, x_{0}\right) \subset B_{C \sqrt{h}}\left(x_{0}\right)$ and there exists an affine transformation $T x=A x+b$ with $C_{1} \leq|\operatorname{det} A|^{\frac{2}{n}} h \leq C_{2},\|A\| \leq C h^{\frac{-1}{2}}$ and $\left\|A^{-1}\right\| \leq C h^{\frac{1}{2}}$ such that

$$
B_{K}(0) \subset \tilde{\Omega}:=T S_{h}\left(u, x_{0}\right) \subset B_{n}(0) \quad \text { and } \quad \operatorname{dist}(x, \partial[\tilde{\Omega}]) \leq \frac{C h^{\frac{-1}{2}}}{\lambda} \quad \text { for all } \quad x \in \partial \tilde{\Omega}
$$

Let us consider the function

$$
\tilde{u}(x):=\frac{1}{h}\left[u\left(T^{-1} x\right)-u\left(x_{0}\right)-c\left(x_{0}-y_{0}\right)-h\right], \quad \text { for } \quad x \in \tilde{\Omega} .
$$

Then $\tilde{u}$ is an Alexandroff weak solution to the Monge-Ampère equation:

$$
\left\{\begin{align*}
\tilde{g}\left(x-\nabla \tilde{c}^{*}(-\nabla \tilde{u})\right) \operatorname{det}\left(I+D^{2} \tilde{c}^{*}(-\nabla \tilde{u}) D^{2} \tilde{u}\right) & =\tilde{f}(x) & & \text { in } \tilde{\Omega},  \tag{7.4}\\
\tilde{u}(x) & =-\tilde{c}\left(x-T y_{0}\right) & & \text { on } \partial \tilde{\Omega}
\end{align*}\right.
$$

for $\tilde{f}(x):=f\left(T^{-1} x\right), \tilde{g}(y):=g\left(T^{-1} y\right)$ and

$$
\tilde{c}(x):=\frac{1}{h} c\left(T^{-1} x\right)=\frac{1}{2}|x|^{2}+\frac{1}{\tilde{\lambda}} \tilde{E}(x) \quad \text { with } \quad \tilde{\lambda}:=h^{\frac{-1}{2}} \lambda .
$$

Equivalently, $\tilde{c}^{*}(z)=\frac{1}{h} c^{*}\left(h T^{t} z\right)$ by the Legendre transform.
We will use the comparison principle from Theorem 6.4 to get a better $L^{\infty}$ approximation for $\tilde{u}$. In order to apply this theorem, we first need to verify the assumptions. Without loss of generality, we can assume that $x_{0}=y_{0}=0$ and $f\left(x_{0}\right)=g\left(y_{0}\right)=1$. We then have

$$
|\tilde{f}(x)-1|=\left|f\left(T^{-1} x\right)-f(0)\right| \leq\left|A^{-1} x\right|^{\alpha} \leq\left\|A^{-1}\right\|^{\alpha}|x|^{\alpha} \leq C h^{\frac{\alpha}{2}}|x|^{\alpha} \quad \text { for all } x \in \tilde{\Omega} .
$$

In particular, $\|\tilde{f}-1\|_{L^{\infty}(\tilde{\Omega})} \leq C h^{\frac{\alpha}{2}}$. On the other hand for $y \in \partial_{\tilde{c}} \tilde{u}(\tilde{\Omega})=T \partial_{c} u\left(S_{h}\left(u, x_{0}\right)\right)$, we write $y=T x$ for $x \in \partial_{c} u\left(S_{h}\left(u, x_{0}\right)\right)$ and thus

$$
|\tilde{g}(y)-1|=|g(x)-g(0)| \leq C|x-0|^{\alpha}=C\left|A^{-1} y\right|^{\alpha} \leq C h^{\frac{\alpha}{2}}|y|^{\alpha} .
$$

Therefore, we would like to apply Theorem 6.4 with $d=C h^{-1}, \lambda \rightsquigarrow \tilde{\lambda}, \theta \rightsquigarrow \tilde{\theta}:=C h^{\frac{\alpha}{2}}$ and $\beta>0$ satisfying $\left(\frac{d}{\hat{\lambda}}\right)^{\beta}=\tilde{\theta}^{n \hat{\alpha}}$, where $\hat{\alpha}=\hat{\alpha}(n, \alpha)$ is chosen arbitrary as long as $0<\hat{\alpha}<$ $\frac{\alpha}{n\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)}$. Notice that $\beta \leq 1$ because we can assume without loss of generality that $\lambda \geq \lambda_{0} \geq \frac{C}{h^{\frac{n \alpha+2}{2}}}$ ensuring that $\frac{d}{\bar{\lambda}} \leq \tilde{\theta}^{\frac{n \hat{\alpha}}{\alpha}}$. Observe also that condition (6.8) is satisfied since

$$
\begin{aligned}
& C_{n}\left(\frac{d}{\tilde{\lambda}}\right)^{-\beta\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)}\left[\tilde{\theta}+\frac{1}{\tilde{\lambda}}\left(\frac{d}{\tilde{\lambda}}\right)^{\frac{n-2}{n}(\alpha-1) \beta}\right] \\
& \leq C_{n} \tilde{\theta}^{-\frac{n \hat{\alpha}}{\alpha}\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)}\left[\tilde{\theta}+\frac{1}{d} \tilde{\theta}^{\frac{n \hat{\alpha}}{\alpha}} \tilde{\theta}^{\hat{\alpha}}(n-2)(\alpha-1)\right.
\end{aligned} \leq 2 C_{n} \tilde{\theta}^{1-\frac{n \hat{\alpha}}{\alpha}\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)} \leq 1 .
$$

Thus we conclude from Theorem 6.4 that

$$
\begin{equation*}
\|\tilde{u}-w\|_{L^{\infty}(\tilde{\Omega})} \leq C\left(\frac{d}{\tilde{\lambda}}\right)^{\frac{2 \beta}{n}}=C \tilde{\theta}^{\frac{2 \hat{\alpha}}{\alpha}}=C h^{\hat{\alpha}}=: \epsilon_{0}, \tag{7.5}
\end{equation*}
$$

where $w$ is the $c_{q}$-convex solution of

$$
\left\{\begin{aligned}
\operatorname{det}\left(I+D^{2} w\right) & =1 & & \text { in } V:=\mathcal{N}_{\tilde{\theta} \tilde{n} \hat{\alpha}}[\Omega], \\
w & =-\frac{1}{2}\left|x-T y_{0}\right|^{2} & & \text { on } \quad \partial V .
\end{aligned}\right.
$$

(Note that $T y_{0} \in \partial_{\tilde{c}} \tilde{u}\left(T x_{0}\right)$ ). The above estimate holds for any $0<h \leq \mu_{1}$. We now fix $h \leq \mu_{1}$ (and hence $\lambda_{0}$ is fixed) such that $\epsilon_{0}=\tau_{0} \mu^{3}$, where $0<\mu \leq \min \left\{\mu_{0}, 1 / 2\right\}$ is a universal constant whose precise value will be determined later and $\mu_{0}, \tau_{0}>0$ are the constants given in Lemma 6.8. In order to prove $u$ is $\mathcal{C}^{2, \alpha^{\prime}}$ at $x_{0}$, it is enough to prove that its normalization $\tilde{u}$ defined in $\tilde{\Omega}$ is $\mathcal{C}^{2, \alpha^{\prime}}$ at $T x_{0}$. However for simplicity, from now on we will abuse notations by writing $u, c, \lambda, \Omega, f, g, x_{0}$ and $y_{0}$ for $\tilde{u}, \tilde{c}, \tilde{\lambda}, \tilde{\Omega}, \tilde{f}, \tilde{g}, T x_{0}$ and $T y_{0}$ respectively. In particular, we have

$$
\begin{aligned}
& B_{K}(0) \subset \Omega \subset B_{n}(0) \quad \text { and } \quad \frac{1}{\lambda} \leq h^{\frac{n \tilde{\alpha}+3}{2}}, \\
& |f(x)-1| \leq \tilde{\theta}|x|^{\alpha} \quad \forall x \in \Omega \quad \text { and } \quad|g(y)-1| \leq \tilde{\theta}|y|^{\alpha} \quad \forall y \in \partial_{c} u(\Omega) .
\end{aligned}
$$

Define

$$
\bar{u}(x):=u(x)+c\left(x-y_{0}\right), \quad \text { for } \quad x \in \Omega .
$$

### 7.2 Induction

To show that $u$ (equivalently $\bar{u}$ ) is $\mathcal{C}^{2, \alpha^{\prime}}$ at the point $x_{0}$, we will compare smaller and smaller sections $S_{\mu}\left(u, x_{0}\right)$ of $u$ with those of a solution $w$ of a Monge-Ampère equation with Dirichlet boundary data. The precise inductive procedure will be carried out next.
First induction step. As a consequence of (7.5) and Lemma 6.9 we can find a map $T_{1} x=A_{1}\left(x-x_{0}\right)$, where $A_{1}$ is a unitary matrix, such that $C_{1}|x|^{2} \leq\left|A_{1} x\right|^{2} \leq C_{2}|x|^{2}$ and

$$
\begin{equation*}
B_{\left(1-\sigma_{1}\right) \sqrt{2}} \subset \mu^{-1 / 2} T_{1} S_{\mu}\left(u, x_{0}\right) \subset B_{\left(1+\sigma_{1}\right) \sqrt{2}}, \quad \text { with } \quad \sigma_{1}:=C\left(\mu^{1 / 2}+\mu^{-1} \epsilon_{0}^{1 / 2}\right) . \tag{7.6}
\end{equation*}
$$

Moreover,

$$
\left|\bar{u}(x)-\left(\bar{u}\left(x_{0}\right)+\left\langle p_{1}, x-x_{0}\right\rangle+\frac{1}{2}\left\langle M_{1}\left(x-x_{0}\right), x-x_{0}\right\rangle\right)\right| \leq C\left(\mu^{3 / 2}+\epsilon_{0}\right) \quad \text { in } \quad S_{\mu}\left(u, x_{0}\right),
$$

where $M_{1}:=A_{1}^{t} A_{1}$ and $p_{1}$ is some vector satisfying $\left|p_{1}\right| \leq C \epsilon_{0}^{1 / 2}$.

Now consider the rescaled function

$$
u_{1}(x):=\frac{1}{\mu}\left[u\left(\mu^{1 / 2} T_{1}^{-1} x\right)-u\left(x_{0}\right)-c\left(x_{0}-y_{0}\right)-\mu\right], \quad x \in \Omega_{1}:=\mu^{-1 / 2} T_{1} S_{\mu}\left(u, x_{0}\right) .
$$

Then, thanks to the assumption (7.3), $u_{1}$ is an Alexandroff weak solution of the following Monge-Ampère equation:

$$
\left\{\begin{align*}
g_{1}\left(x-\nabla c_{1}^{*}\left(-\nabla u_{1}\right)\right) \operatorname{det}\left(I+D^{2} c_{1}^{*}\left(-\nabla u_{1}\right) D^{2} u_{1}\right) & =f_{1}(x) & & \text { in } \Omega_{1},  \tag{7.7}\\
u_{1}(x) & =-c_{1}\left(x-\mu^{-\frac{1}{2}} T_{1} y_{0}\right) & & \text { on } \partial \Omega_{1},
\end{align*}\right.
$$

for

$$
f_{1}(x):=f\left(\mu^{1 / 2} T_{1}^{-1} x\right), \quad g_{1}(y):=g\left(\mu^{1 / 2} T_{1}^{-1} y\right)
$$

and

$$
c_{1}(x):=\frac{1}{\mu} c\left(\mu^{1 / 2} T_{1}^{-1} x\right)=\frac{1}{2}|x|^{2}+\frac{1}{\lambda_{1}} E_{1}(x) \quad \text { with } \quad \frac{1}{\lambda_{1}}:=\frac{\mu^{1 / 2}}{\lambda} .
$$

The definition of $c_{1}$ implies that $c_{1}^{*}(z)=\frac{1}{\mu} c^{*}\left(\mu^{1 / 2} T_{1}^{t} z\right)$. It is clear from (7.7) that the boundary values of $u_{1}$ are controlled. Notice also that we can apply Proposition 3.2 to say,

$$
\operatorname{dist}\left(x, \partial\left[S_{\mu}\left(u, x_{0}\right)\right]\right) \leq \frac{C}{\lambda} \quad \text { for all } \quad x \in \partial S_{\mu}\left(u, x_{0}\right)
$$

and after rescaling,

$$
\operatorname{dist}\left(x, \partial\left[\Omega_{1}\right]\right) \leq \mu^{\frac{-1}{2}}\left\|A_{1}\right\| \frac{C}{\lambda} \leq \frac{C \sqrt{C_{2}} \mu^{-1}}{\lambda_{1}} \quad \forall x \in \partial \Omega_{1} .
$$

We have

$$
\begin{equation*}
\left|f_{1}(x)-1\right|=\left|f\left(\mu^{\frac{1}{2}} T_{1}^{-1} x\right)-1\right| \leq \tilde{\theta} \mu^{\frac{\alpha}{2}}\left|A_{1}^{-1} x\right|^{\alpha} \leq \tilde{\theta}\left(C_{1}^{-1} \mu\right)^{\frac{\alpha}{2}}|x|^{\alpha}=\theta_{1}|x|^{\alpha} \quad \forall x \in \Omega_{1}, \tag{7.8}
\end{equation*}
$$

where $\theta_{1}:=\tilde{\theta}\left(C_{1}^{-1} \mu\right)^{\frac{\alpha}{2}} \leq \tilde{\theta}$. Also as $\mu^{\frac{1}{2}} T_{1}^{-1} \partial_{c_{1}} u_{1}\left(\Omega_{1}\right)=\partial_{c} u\left(S_{\mu}\left(u, x_{0}\right)\right)$,

$$
\begin{equation*}
\left|g_{1}(y)-1\right|=\left|g\left(\mu^{\frac{1}{2}} T_{1}^{-1} y\right)-1\right| \leq \tilde{\theta} \mu^{\frac{\alpha}{2}}\left|A_{1}^{-1} y\right|^{\alpha} \leq \theta_{1}|y|^{\alpha} \quad \text { for all } y \in \partial_{c_{1}} u_{1}\left(\Omega_{1}\right) . \tag{7.9}
\end{equation*}
$$

Therefore, we would like to apply Theorem 6.4 with $d=C \sqrt{C_{2}} \mu^{-1}, \lambda \rightsquigarrow \lambda_{1}, \theta \rightsquigarrow \theta_{1}$ and $\beta>0$ satisfying $\left(\frac{d}{\lambda_{1}}\right)^{\beta}=\theta_{1}^{\frac{n \hat{\alpha}}{\alpha}}$. Notice that $\beta \leq 1$ because our condition on $\hat{\alpha}$ yields

$$
\begin{equation*}
\frac{d}{\lambda_{1}}=\frac{C \sqrt{C_{2}} \mu^{1 / 2} \mu^{-1}}{\lambda} \leq C \sqrt{C_{2}} \mu^{1 / 2} \mu^{-1} h^{\frac{n \hat{\alpha}+2}{2}} \leq C \mu^{\left(\frac{3}{\alpha}-\frac{1+n \hat{\alpha}}{2}\right)} \theta_{1}^{\frac{n \hat{\alpha}}{\alpha}} \leq \theta_{1}^{\frac{n \hat{\alpha}}{\alpha}} \tag{7.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{\lambda_{1}}=\frac{\mu^{\frac{1}{2}}}{\lambda} \leq C h^{\frac{n \hat{\alpha}+1}{2}}(h \mu)^{\frac{(1-\alpha)}{2}[1-(n-2) \hat{\alpha}]} \theta_{1} \theta_{1}^{\frac{\hat{\alpha}}{\alpha}(n-2)(1-\alpha)} \leq \theta_{1} \theta_{1}^{\frac{\hat{\alpha}}{\alpha}(n-2)(1-\alpha)} \tag{7.11}
\end{equation*}
$$

Observe also that condition (6.8) is satisfied since

$$
\begin{aligned}
& C_{n}\left(\frac{d}{\lambda_{1}}\right)^{-\beta\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)}\left[\theta_{1}+\frac{1}{\lambda_{1}}\left(\frac{d}{\lambda_{1}}\right)^{\frac{n-2}{n}(\alpha-1) \beta}\right] \\
& \leq C_{n} \theta_{1}^{-\frac{n \hat{\alpha}}{\alpha}\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)}\left[\theta_{1}+\theta_{1} \theta_{1}^{\frac{\hat{\alpha}}{\alpha}(n-2)(1-\alpha)} \theta_{1}^{\frac{\hat{\alpha}}{\alpha}(n-2)(\alpha-1)}\right] \\
& \leq 2 C_{n} \theta_{1}^{1-\frac{n \hat{\alpha}}{\alpha}\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)} \leq 2 C_{n} \tilde{\theta}^{1-\frac{n \hat{\alpha}}{\alpha}\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)} \leq 1 .
\end{aligned}
$$

We have checked all the hypotheses necessary to apply Theorem 6.4. Thus Theorem 6.4 gives

$$
\begin{equation*}
\left\|u_{1}-w_{1}\right\|_{L^{\infty}\left(\Omega_{1}\right)} \leq C\left(\frac{d}{\lambda_{1}}\right)^{\frac{2 \beta}{n}}=C \theta_{1}^{\frac{2 \hat{\alpha}}{\alpha}}=: \epsilon_{1} \tag{7.12}
\end{equation*}
$$

where $w_{1}$ is the $c_{q}$-convex solution of

$$
\left\{\begin{aligned}
\operatorname{det}\left(I+D^{2} w_{1}\right) & =1 & & \text { in } \quad V:=\mathcal{N}_{\theta_{1}^{\frac{n \hat{\alpha}}{\alpha}}}\left[\Omega_{1}\right] \\
w_{1} & =-\frac{1}{2}\left|x-\mu^{\frac{-1}{2}} T_{1} y_{0}\right|^{2} & & \text { on } \partial V .
\end{aligned}\right.
$$

Then as $\epsilon_{1} \leq \epsilon_{0} \leq \tau_{0} \mu^{2}$, Lemma 6.10 (now centered at $x_{0}=0$ ) shows the existence of a positive definite matrix $M=A^{t} A$ and a vector $p \in \mathbb{R}^{n}$ satisfying

$$
\operatorname{det} M=1, \quad\left(1-C \sigma_{1}\right) I \leq M \leq\left(1+C \sigma_{1}\right) I, \quad|p| \leq C \sigma_{1}
$$

such that

$$
\begin{equation*}
B_{\left(1-\sigma_{2}\right) \sqrt{2}} \subset \mu^{-1 / 2} A S_{\mu}\left(u_{1}, 0\right) \subset B_{\left(1+\sigma_{2}\right) \sqrt{2}}, \quad \text { with } \quad \sigma_{2}:=C\left(\sigma_{1} \mu^{1 / 2}+\mu^{-1} \epsilon_{1}^{1 / 2}\right) \tag{7.13}
\end{equation*}
$$

Moreover

$$
\begin{equation*}
\left|\bar{u}_{1}(x)-\left(\bar{u}_{1}(0)+\langle p, x\rangle+\frac{1}{2}\langle M x, x\rangle\right)\right| \leq C\left(\sigma_{1} \mu^{3 / 2}+\epsilon_{1}\right) \quad \text { in } \quad S_{\mu}\left(u_{1}, 0\right) \tag{7.14}
\end{equation*}
$$

where

$$
\bar{u}_{1}(x):=u_{1}(x)+c_{1}\left(x-\mu^{\frac{-1}{2}} T_{1} y_{0}\right)=\frac{1}{\mu}\left[\bar{u}\left(\mu^{1 / 2} T_{1}^{-1} x\right)-u\left(x_{0}\right)-c\left(x_{0}-y_{0}\right)-\mu\right]
$$

Let $M_{2}:=A_{1}^{t} M A_{1}, A_{1}$ is computed in the previous step. Since $M_{2}-M_{1}=A_{1}^{t}(M-I) A_{1}$, we have that $\left\|M_{2}-M_{1}\right\| \leq\left\|A_{1}\right\|^{2}\|M-I\| \leq C \sigma_{1}$. Also, because $M=A^{t} A$, we have that $\left(1-C \sigma_{1}\right)|x|^{2} \leq|A x|^{2} \leq\left(1+C \sigma_{1}\right)|x|^{2}$. If we let $A_{2}:=A A_{1}$, then $M_{2}=A_{2}^{t} A_{2}$ and therefore,

$$
C\left(1-C \sigma_{1}\right)|x|^{2} \leq\left(1-C \sigma_{1}\right)\left|A_{1} x\right|^{2} \leq\left|A A_{1} x\right|^{2} \leq\left(1+C \sigma_{1}\right)\left|A_{1} x\right|^{2} \leq C\left(1+C \sigma_{1}\right)|x|^{2}
$$

that is $\left\|A_{2}\right\|^{2} \leq C\left(1+C \sigma_{1}\right)$ and $\left\|A_{2}^{-1}\right\|^{2} \leq 1 / C\left(1-C \sigma_{1}\right)$.
On the other hand, $S_{\mu}\left(u_{1}, 0\right)=\mu^{-1 / 2} T_{1} S_{\mu^{2}}\left(u, x_{0}\right)$. Hence it follows from (7.13) that

$$
\begin{equation*}
B_{\left(1-\sigma_{2}\right) \sqrt{2}} \subset \mu^{-1} T_{2} S_{\mu^{2}}\left(u, x_{0}\right) \subset B_{\left(1+\sigma_{2}\right) \sqrt{2}} \tag{7.15}
\end{equation*}
$$

for $T_{2} x:=A_{2}\left(x-x_{0}\right)$. Moreover, from (7.14), letting $x=\mu^{-1 / 2} T_{1} z$ we obtain

$$
\left|\bar{u}(z)-\left(\bar{u}\left(x_{0}\right)+\left\langle p_{2}, z\right\rangle+\frac{1}{2}\left\langle M_{2} z, z\right\rangle\right)\right| \leq C \mu\left(\sigma_{1} \mu^{3 / 2}+\epsilon_{1}\right) \quad \text { for all } \quad z \in S_{\mu^{2}}\left(u, x_{0}\right)
$$

where $p_{2}:=\mu^{1 / 2} A_{1}^{t} p$ and so $\left|p_{2}\right| \leq C \mu^{1 / 2} \sigma_{1}$.
We are now in good shape to iterate the induction process, that will be explained in the following paragraph. The crucial observation is that problem (7.7) is better than the original one in the sense that the rescaled cost $c_{1}$ is closer to the quadratic one by a factor $\mu^{1 / 2}$.

The general induction step. It is clear from above step that how the induction process continues. The main point is the following: while the gradient of $u_{k}$ deteriorates by a factor of $\mu^{k / 2}$ (bad!) with respect to the gradient of the original $u$, on the other hand the $\operatorname{cost} c_{\lambda_{k}}$ is
closer to quadratic by the same factor (good!). The strategy is to find the right compensation between the two facts.

Let us explain the main details of the $k$-th step given the $(k-1)$-th step. From the ( $k-1$ )-th step, we have

$$
\begin{equation*}
\left\|u_{k-1}-w_{k-1}\right\|_{L^{\infty}\left(\Omega_{k-1}\right)} \leq C \theta_{k-1}^{\frac{2 \hat{\alpha}}{\alpha}}=: \epsilon_{k-1}, \tag{7.16}
\end{equation*}
$$

where $\Omega_{k-1}:=\mu^{-\frac{k-1}{2}} T_{k-1} S_{\mu^{k-1}}\left(u, x_{0}\right)$,

$$
u_{k-1}(x):=\frac{1}{\mu^{k-1}}\left[u\left(\mu^{\frac{k-1}{2}} T_{k-1}^{-1} x\right)-u\left(x_{0}\right)-c\left(x_{0}-y_{0}\right)-\mu^{k-1}\right]
$$

and $w_{k-1}$ is the $c_{q}$-convex solution of

$$
\left\{\begin{aligned}
\operatorname{det}\left(I+D^{2} w_{k-1}\right) & =1 & & \text { in } \quad V:=\mathcal{N}_{\theta_{k-1}^{\frac{n \hat{\alpha}}{\alpha}}}\left[\Omega_{k-1}\right], \\
w_{k-1} & =-\frac{1}{2}\left|x-\mu^{-\frac{k-1}{2}} T_{k-1} y_{0}\right|^{2} & & \text { on } \quad \partial V .
\end{aligned}\right.
$$

Also there is a positive definite matrix $M_{k}=A_{k}^{t} A_{k}$ such that for $T_{k} x:=A_{k}\left(x-x_{0}\right)$,

$$
B_{\left(1-\sigma_{k}\right) \sqrt{2}} \subset \mu^{-\frac{k}{2}} T_{k} S_{\mu^{k}}\left(u, x_{0}\right) \subset B_{\left(1+\sigma_{k}\right) \sqrt{2}}
$$

To continue the process, we consider the rescaled function

$$
u_{k}(x):=\frac{1}{\mu^{k}}\left[u\left(\mu^{\frac{k}{2}} T_{k}^{-1} x\right)-u\left(x_{0}\right)-c\left(x_{0}-y_{0}\right)-\mu^{k}\right], \quad x \in \Omega_{k}:=\mu^{-\frac{k}{2}} T_{k} S_{\mu^{k}}\left(u, x_{0}\right) .
$$

Note that $u_{k}$ satisfies a Monge-Ampère equation in the weak sense of Alexandroff with a new $\operatorname{cost} c_{k}$, i.e.,

$$
\left\{\begin{aligned}
g_{k}\left(x-\nabla c_{k}^{*}\left(-\nabla u_{k}\right)\right) \operatorname{det}\left(I+D^{2} c_{k}^{*}\left(-\nabla u_{k}\right) D^{2} u_{k}\right) & =f_{k}(x) & & \text { in } \Omega_{k}, \\
u_{k}(x) & =-c_{k}\left(x-\mu^{-\frac{k}{2}} T_{k} y_{0}\right) & & \text { on } \partial \Omega_{k},
\end{aligned}\right.
$$

for

$$
f_{k}(x):=f_{k-1}\left(\mu^{\frac{1}{2}} A^{-1} x\right)=f\left(\mu^{\frac{k}{2}} T_{k}^{-1} x\right), \quad g_{k}(y):=g_{k-1}\left(\mu^{\frac{1}{2}} A^{-1} y\right)=g\left(\mu^{\frac{k}{2}} T_{k}^{-1} y\right)
$$

and

$$
c_{k}(x):=\frac{1}{\mu} c_{k-1}\left(\mu^{\frac{1}{2}} A^{-1} x\right)=\frac{1}{\mu^{k}} c\left(\mu^{\frac{k}{2}} T_{k}^{-1} x\right) \quad \text { or equivalently } \quad c_{k}^{*}(z)=\frac{1}{\mu^{k}} c^{*}\left(\mu^{\frac{k}{2}} T_{k}^{t} z\right),
$$

where the matrix $A$ is given precisely in the $(k-1)$-th step which in particular satisfy $A_{k}=A A_{k-1},\|A\| \leq \sqrt{C_{2}}$ and $\left\|A^{-1}\right\| \leq C_{1}^{-1 / 2}$. Because $\left\|A^{-1}\right\|$ is universally bounded, we can write

$$
c_{k}(x)=\frac{1}{2}|x|^{2}+\frac{1}{\lambda_{k}} E_{k}(x) \quad \text { with } \quad \frac{1}{\lambda_{k}}:=\frac{\mu^{1 / 2}}{\lambda_{k-1}}=\frac{\mu^{k / 2}}{\lambda}
$$

for some

$$
\left|E_{k}(x)\right| \leq C|x|^{3}, \quad\left|\nabla E_{k}(x)\right| \leq C|x|^{2}, \quad\left|D^{2} E_{k}(x)\right| \leq C|x| .
$$

We write

$$
M_{\lambda_{k}}\left[u_{k}\right]=f_{k}
$$

for

$$
M_{\lambda_{k}}\left[u_{k}\right]:=g_{k}\left(x+\nabla u_{k}+G_{k}\left(\nabla u_{k}\right)\right) \operatorname{det}\left[I+D^{2} u_{k}+\left(D^{2} u_{k}\right) H_{k}\left(\nabla u_{k}\right)\right]
$$

where $H_{k}$ and $G_{k}$ satisfy for some universal constant $C$,

$$
\left\|H_{k}\left(\nabla u_{k}\right)\right\|_{L^{\infty}} \leq \frac{C}{\lambda_{k}}\left\|\nabla u_{k}\right\|_{L^{\infty}\left(\Omega_{k}\right)}, \quad\left\|G_{k}\left(\nabla u_{k}\right)\right\|_{L^{\infty}} \leq \frac{C}{\lambda_{k}}\left\|\nabla u_{k}\right\|_{L^{\infty}\left(\Omega_{k}\right)}^{2} .
$$

First observe that Proposition 3.2 and the fact $\Omega_{k}=\mu^{-1 / 2} A S_{\mu}\left(u_{k-1}, 0\right)$ give

$$
\operatorname{dist}\left(x, \partial\left[\Omega_{k}\right]\right) \leq \mu^{\frac{-1}{2}}\|A\| \frac{C}{\lambda_{k-1}} \leq \frac{C \sqrt{C_{2}} \mu^{-1}}{\lambda_{k}} \quad \forall x \in \partial \Omega_{k} .
$$

Next since $\mu^{\frac{1}{2}} A^{-1} \Omega_{k}=\Omega_{k-1}$, we obtain from the $(k-1)$-th step that

$$
\left|f_{k}(x)-1\right|=\left|f_{k-1}\left(\mu^{\frac{1}{2}} A^{-1} x\right)-1\right| \leq \theta_{k-1} \mu^{\frac{\alpha}{2}}\left|A^{-1} x\right|^{\alpha} \leq \theta_{k}|x|^{\alpha} \quad \text { for all } x \in \Omega_{k},
$$

where $\theta_{k}:=\theta_{k-1}\left(C_{1}^{-1} \mu\right)^{\frac{\alpha}{2}} \leq \theta_{k-1}$. Similarly as $\mu^{\frac{1}{2}} A^{-1} \partial_{c_{k}} u_{k}\left(\Omega_{k}\right)=\partial_{c_{k-1}} u_{k-1}\left(S_{\mu}\left(u_{k-1}, 0\right)\right)$,

$$
\left|g_{k}(y)-1\right|=\left|g_{k-1}\left(\mu^{\frac{1}{2}} A^{-1} y\right)-1\right| \leq \theta_{k-1} \mu^{\frac{\alpha}{2}}\left|A^{-1} y\right|^{\alpha} \leq \theta_{k}|y|^{\alpha} \quad \text { for all } y \in \partial_{c_{k}} u_{k}\left(\Omega_{k}\right)
$$

We apply Theorem 6.4 with $d=C \sqrt{C_{2}} \mu^{-1}, \lambda \rightsquigarrow \lambda_{k}, \theta \rightsquigarrow \theta_{k}$ and $\beta>0$ satisfying $\left(\frac{d}{\lambda_{k}}\right)^{\beta}=$ $\theta_{k}^{\frac{n \hat{\alpha}}{\alpha}}$. Notice that the constant $d$ is the same for all steps. Also $\beta \leq 1$ because it follows from the ( $k-1$ )-th step that

$$
\frac{d}{\lambda_{k}}=\mu^{\frac{1}{2}} \frac{d}{\lambda_{k-1}} \leq \mu^{\frac{1}{2}} \theta_{k-1}^{\frac{n \hat{\alpha}}{\alpha}}=C_{1}^{\frac{n \hat{\alpha}}{2}} \mu^{\frac{1-n \hat{\alpha}}{2}} \theta_{k}^{\frac{n \hat{\alpha}}{\alpha}} \leq \theta_{k}^{\frac{n \hat{\alpha}}{\alpha}}
$$

and

$$
\begin{aligned}
\frac{1}{\lambda_{k}}=\frac{\mu^{\frac{1}{2}}}{\lambda_{k-1}} & \leq \mu^{\frac{1}{2}} \theta_{k-1} \theta_{k-1}^{\frac{\hat{\alpha}}{\alpha}(n-2)(1-\alpha)}=C_{1}^{\frac{1}{2}(\alpha+\hat{\alpha}(n-2)(1-\alpha))} \mu^{\frac{(1-\alpha)}{2}[1-(n-2) \hat{\alpha}]} \theta_{k} \theta_{k}^{\frac{\hat{\alpha}}{\alpha}(n-2)(1-\alpha)} \\
& \leq \theta_{k} \theta_{k}^{\frac{\hat{\alpha}}{\alpha}(n-2)(1-\alpha)} .
\end{aligned}
$$

These also imply that condition (6.8) is satisfied because

$$
\begin{aligned}
& C_{n}\left(\frac{d}{\lambda_{k}}\right)^{-\beta\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)}\left[\theta_{k}+\frac{1}{\lambda_{k}}\left(\frac{d}{\lambda_{k}}\right)^{\frac{n-2}{n}(\alpha-1) \beta}\right] \\
& \leq 2 C_{n} \theta_{k}^{1-\frac{n \hat{\alpha}}{\alpha}\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)} \leq 2 C_{n} \theta_{k-1}^{1-\frac{n \hat{\alpha}}{\alpha}\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)} \leq 1 .
\end{aligned}
$$

Now apply Theorem 6.4 and continue the same procedure as done in the first induction step. Observe that we always have

$$
\epsilon_{k}:=C\left(\frac{d}{\lambda_{k}}\right)^{\frac{2 \beta}{n}}=C \theta_{k}^{\frac{2 \hat{\alpha}}{\alpha}} .
$$

Conclusion. The above induction process allows to construct a sequence of matrices $M_{k}=$ $A_{k}^{t} A_{k}$, transformations $T_{k}$, vectors $p_{k} \in \mathbb{R}^{n}$, and polynomials $P_{k}(x), k \geq 1$, such that

1. $\operatorname{det} M_{k}=1, \quad\left\|M_{k+1}-M_{k}\right\| \leq C\left(1+C \sigma_{1}\right) \cdots\left(1+C \sigma_{k-1}\right) \sigma_{k}$,
2. $\left\|A_{k}\right\|^{2} \leq C\left(1+C \sigma_{1}\right) \cdots\left(1+C \sigma_{k-1}\right), \quad\left\|A_{k}^{-1}\right\|^{2} \leq C\left[\left(1-C \sigma_{1}\right) \cdots\left(1-C \sigma_{k-1}\right)\right]^{-1}$,
3. $B_{\left(1-\sigma_{k}\right) \sqrt{2}} \subset \mu^{-k / 2} T_{k} S_{\mu^{k}}\left(u, x_{0}\right) \subset B_{\left(1+\sigma_{k}\right) \sqrt{2}}$,
4. $\left\|\bar{u}(x)-P_{k}(x)\right\|_{L^{\infty}\left(S_{\mu^{k}}\left(u, x_{0}\right)\right)} \leq C \mu^{k-1}\left(\sigma_{k-1} \mu^{3 / 2}+\epsilon_{k-1}\right)$,
5. $\left|p_{k}\right| \leq C \mu^{\frac{k-1}{2}} \sqrt{\left(1+C \sigma_{1}\right) \cdots\left(1+C \sigma_{k-2}\right)} \sigma_{k-1}$,
where
6. $P_{k}(x):=\bar{u}\left(x_{0}\right)+\left\langle p_{k}, x-x_{0}\right\rangle+\frac{1}{2}\left\langle M_{k}\left(x-x_{0}\right), x-x_{0}\right\rangle$,
7. $\epsilon_{0}=\tau_{0} \mu^{3}$ and $\epsilon_{k}=C \theta_{k}^{\frac{2 \hat{\alpha}}{\alpha}}=\epsilon_{0}\left(C_{1}^{-1} \mu\right)^{k \hat{\alpha}}$ for $k \geq 1$,
8. $\sigma_{0}=1$ and $\sigma_{k}=C\left(\sigma_{k-1} \mu^{1 / 2}+\mu^{-1} \epsilon_{k-1}^{1 / 2}\right)$ for $k \geq 1$.

Let us estimate $\sigma_{k}$. First choose $\mu$ small enough such that $\sigma_{1}=C\left(1+\tau_{0}^{\frac{1}{2}}\right) \mu^{\frac{1}{2}}<1$. Then as the sequence $\left\{\epsilon_{k}\right\}_{k=1}^{\infty}$ is decreasing to 0 , it is clear that $\left\{\sigma_{k}\right\}_{k=1}^{\infty}$ is decreasing to 0 as well.

By induction we see that for all $k \geq 1$

$$
\begin{aligned}
\sigma_{k} & =(C \sqrt{\mu})^{k}+\frac{C}{\mu}(C \sqrt{\mu})^{k-1} \sum_{i=0}^{k-1}(C \sqrt{\mu})^{-i} \epsilon_{i}^{1 / 2} \\
& =(C \sqrt{\mu})^{k}+\frac{C \epsilon_{0}^{1 / 2}}{\mu}(C \sqrt{\mu})^{k-1} \sum_{i=0}^{k-1}\left[C^{-1} C_{1}^{\frac{-\hat{\alpha}}{2}} \mu^{\frac{\hat{\alpha}-1}{2}}\right]^{i} \\
& \leq(C \sqrt{\mu})^{k}+\frac{C \epsilon_{0}^{1 / 2}}{\mu}(C \sqrt{\mu})^{k-1} \frac{C^{-k} C_{1}^{\frac{-k \hat{\alpha}}{2}} \mu^{\frac{k(\hat{\alpha}-1)}{2}}}{C^{-1} C_{1}^{\frac{-\hat{\alpha}}{2}} \mu^{\frac{\hat{\alpha}-1}{2}}} \leq \hat{C}\left(C_{1}^{-1} \mu\right)^{\frac{k \hat{\alpha}}{2}} .
\end{aligned}
$$

Thus $\sum_{i=1}^{\infty} \sigma_{i}$ is convergent and hence $\prod_{i=1}^{\infty}\left(1+C \sigma_{i}\right)$ and $\prod_{i=1}^{\infty}\left(1-C \sigma_{i}\right)$ converge to some positive numbers since $C \sigma_{i} \leq 1 / 2$ for all $i$ large. Consequently, we obtain the following estimates for all $k \geq 1$ :

$$
\left\|A_{k}\right\| \leq C, \quad\left\|A_{k}^{-1}\right\| \leq C, \quad\left|p_{k}\right| \leq C \mu^{\frac{k-1}{2}}\left(C_{1}^{-1} \mu\right)^{\frac{(k-1) \hat{\alpha}}{2}} \text { and }\left\|M_{k+1}-M_{k}\right\| \leq C\left(C_{1}^{-1} \mu\right)^{\frac{k \hat{\alpha}}{2}} .
$$

It follows that $M_{k}$ is a Cauchy sequence and hence $M_{k} \rightarrow M$ for some matrix $M$ and $p_{k} \rightarrow 0$ as $k \rightarrow \infty$. So if we let $P_{0}(x):=\bar{u}\left(x_{0}\right)+\left\langle M\left(x-x_{0}\right),\left(x-x_{0}\right)\right\rangle$, then $P_{k}(x) \rightarrow P_{0}(x)$ pointwise in $\mathbb{R}^{n}$. Also the above estimates together with the fact $T_{k} x_{0}=0$ yield $B_{C \mu^{k / 2}}\left(x_{0}\right) \subset$ $S_{\mu^{k}}\left(u, x_{0}\right)$ for all $k \geq 1$. Therefore, for any $l=1,2, \ldots$ and for any $x \in B_{C \mu^{l / 2}}\left(x_{0}\right) \subset S_{\mu^{l}}\left(u, x_{0}\right)$, we obtain

$$
\begin{aligned}
\mid \bar{u}(x) & -P_{0}(x)\left|\leq\left|\bar{u}(x)-P_{l}(x)\right|+\left|P_{l}(x)-P_{0}(x)\right|\right. \\
& \leq C \mu^{l-1}\left(\sigma_{l-1} \mu^{3 / 2}+\epsilon_{l-1}\right)+\left\|p_{l}\right\|\left|x-x_{0}\right|+\sum_{k=l}^{\infty}\left|\left\langle\left(M_{k}-M_{k+1}\right)\left(x-x_{0}\right),\left(x-x_{0}\right)\right\rangle\right| \\
& \leq C^{\prime} \mu^{l-1}\left(C_{1}^{-1} \mu\right)^{\frac{(l-1) \hat{\alpha}}{2}}+C^{\prime} \mu^{l} \sum_{k=l}^{\infty}\left(C_{1}^{-1} \mu\right)^{\frac{k \hat{\alpha}}{2}} \leq C^{\prime \prime} \mu^{l}\left(C_{1}^{-1} \mu\right)^{\frac{l \hat{\alpha}}{2}} \leq C^{\prime \prime} \mu^{\frac{l}{2}\left(2+\alpha^{\prime}\right)},
\end{aligned}
$$

where the last inequality holds if we choose $0<\alpha^{\prime}<\hat{\alpha}$ (note that $C_{1}<1$ ). This yields (7.2) as desired and we conclude that $u$ is $\mathcal{C}^{2, \alpha^{\prime}}$ at $x_{0}$ for any $0<\alpha^{\prime}<\frac{\alpha}{n\left((n-1)^{2}+\frac{n-2}{n} \alpha\right)}$. The proof of Theorem 1.2 is completed.

## References

[1] L. Ambrosio, N. Gigli, and G. Savaré. Gradient flows in metric spaces and in the space of probability measures. Lectures in Mathematics ETH Zürich. Birkhäuser Verlag, Basel, 2005.
[2] Y. Brenier. Polar factorization and monotone rearrangement of vector-valued functions. Comm. Pure Appl. Math., 44(4):375-417, 1991.
[3] L. Caffarelli and Y. Y. Li. A Liouville theorem for solutions of the Monge-Ampère equation with periodic data. Ann. Inst. H. Poincaré Anal. Non Linéaire, 21(1):97-120, 2004.
[4] L. A. Caffarelli. Interior a priori estimates for solutions of fully nonlinear equations. Ann. of Math. (2), 130(1):189-213, 1989.
[5] L. A. Caffarelli. Interior $W^{2, p}$ estimates for solutions of the Monge-Ampère equation. Ann. of Math. (2), 131(1):135-150, 1990.
[6] L. A. Caffarelli. A localization property of viscosity solutions to the Monge-Ampère equation and their strict convexity. Ann. of Math. (2), 131(1):129-134, 1990.
[7] L. A. Caffarelli. The regularity of mappings with a convex potential. J. Amer. Math. Soc., 5(1):99-104, 1992.
[8] L. A. Caffarelli. Boundary regularity of maps with convex potentials. Comm. Pure Appl. Math., 45(9):1141-1151, 1992.
[9] L. A. Caffarelli. Allocation maps with general cost functions. In Partial differential equations and applications, volume 177 of Lecture Notes in Pure and Appl. Math., pages 29-35. Dekker, New York, 1996.
[10] L. A. Caffarelli. Boundary regularity of maps with convex potentials. II. Ann. of Math. (2), 144(3):453-496, 1996.
[11] L. A. Caffarelli. A priori estimates and the geometry of the Monge Ampère equation. In Nonlinear partial differential equations in differential geometry (Park City, UT, 1992), volume 2 of IAS/Park City Math. Ser., pages 5-63. Amer. Math. Soc., Providence, RI, 1996.
[12] L. A. Caffarelli and X. Cabré. Fully nonlinear elliptic equations, volume 43 of American Mathematical Society Colloquium Publications. American Mathematical Society, Providence, RI, 1995.
[13] L. A. Caffarelli, C. E. Gutiérrez, and Q. Huang. On the regularity of reflector antennas. Ann. of Math. (2), 167(1):299-323, 2008.
[14] G. De Philippis and A. Figalli. Partial regularity for optimal transport maps. Preprint, 2012.
[15] P. Delanoë. Gradient rearrangement for diffeomorphisms of a compact manifold. Differential Geom. Appl., 20(2):145-165, 2004.
[16] P. Delanoë and Y. Ge. Regularity of optimal transport on compact, locally nearly spherical, manifolds. J. Reine Angew. Math., 646:65-115, 2010.
[17] A. Figalli, Y.-H. Kim, and R. J. McCann. Hölder Continuity and Injectivity of Optimal Maps. Arch. Ration. Mech. Anal., 209(3):747-795, 2013.
[18] A. Figalli, Y.-H. Kim, and R. J. McCann. Regularity of optimal transport maps on multiple products of spheres. J. Eur. Math. Soc. (JEMS), to appear.
[19] A. Figalli and G. Loeper. $C^{1}$ regularity of solutions of the Monge-Ampère equation for optimal transport in dimension two. Calc. Var. Partial Differential Equations, 35(4):537550, 2009.
[20] A. Figalli and L. Rifford. Continuity of optimal transport maps and convexity of injectivity domains on small deformations of $\mathbb{S}^{2}$. Comm. Pure Appl. Math. 62, no. 12, 1670-1706, 2009.
[21] A. Figalli, L. Rifford and C. Villani. On the Ma-Trudinger-Wang curvature on surfaces. Calc. Var. Partial Differential Equations 39, no. 3-4, 307-332, 2010.
[22] A. Figalli, L. Rifford, and C. Villani. Necessary and sufficient conditions for continuity of optimal transport maps on Riemannian manifolds. Tohoku Math. J. (2), 63(4):855-876, 2011.
[23] W. Gangbo and R. J. McCann. Optimal maps in Monge's mass transport problem. C. R. Acad. Sci. Paris Sér. I Math., 321(12):1653-1658, 1995.
[24] W. Gangbo and R. J. McCann. The geometry of optimal transportation. Acta Math., 177(2):113-161, 1996.
[25] D. Gilbarg and N.S. Trudinger. Elliptic Partial Differential Equations of Second Order. Springer-Verlag, New York, 2001.
[26] C. E. Gutiérrez. The Monge-Ampère equation. Progress in Nonlinear Differential Equations and their Applications, 44. Birkhäuser Boston Inc., Boston, MA, 2001.
[27] C. E. Gutiérrez, Q. Huang and T. V. Nguyen. Interior Hölder estimates for solutions of the Monge-Ampère equation. Preprint.
[28] C. E. Gutiérrez and T. V. Nguyen. On Monge-Ampère type equations arising in optimal transportation problems. Calc. Var. Partial Differential Equations, 28(3):275-316, 2007.
[29] Y.-H. Kim and R. J. McCann. Continuity, curvature, and the general covariance of optimal transportation. J. Eur. Math. Soc. (JEMS) 12, no. 4, 1009-1040, 2010.
[30] Y.-H. Kim and R. J. McCann. Towards the smoothness of optimal maps on Riemannian submersions and Riemannian products (of round spheres in particular). J. Reine Angew. Math., 664, 1-27, 2012.
[31] J. Liu. Hölder regularity of optimal mappings in optimal transportation. Calc. Var. Partial Differential Equations, 34, 435-451, 2009.
[32] J. Liu, N.S. Trudinger and X.-J. Wang. Interior $C^{2, \alpha}$ regularity for potential functions in optimal transportation. Comm. Partial Differential Equations 35, no. 1, 165-184, 2010.
[33] G. Loeper. On the regularity of solutions of optimal transportation problems. Acta Math., 202(2):241-283, 2009.
[34] G. Loeper. Regularity of optimal maps on the sphere: the quadratic cost and the reflector antenna. Arch. Ration. Mech. Anal. 199, no. 1, 269-289, 2011.
[35] X.-N. Ma, N. S. Trudinger, and X.-J. Wang. Regularity of potential functions of the optimal transportation problem. Arch. Ration. Mech. Anal., 177(2):151-183, 2005.
[36] R. Schneider. Convex bodies: the Brunn-Minkowski theory. Encyclopedia of Mathematics and its Applications, 44. Cambridge University Press, Cambridge, 1993.
[37] N. S. Trudinger and X.-J. Wang. On strict convexity and continuous differentiability of potential functions in optimal transportation. Arch. Ration. Mech. Anal., 192(3):403418, 2009.
[38] N. S. Trudinger and X.-J. Wang. On convexity notions in optimal transportation. Preprint, 2008.
[39] N. S. Trudinger and X.-J. Wang. On the second boundary value problem for MongeAmpère type equations and optimal transportation. Ann. Sc. Norm. Super. Pisa Cl. Sci. (5), 8(1):143-174, 2009.
[40] C. Villani. Topics in optimal transportation, volume 58 of Graduate Studies in Mathematics. American Mathematical Society, Providence, RI, 2003.
[41] C. Villani. Optimal transport, Old and new, volume 338 of Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences]. SpringerVerlag, Berlin, 2009.
[42] M. Warren. Regularity for a log-concave to log-concave mass transfer problem with near Euclidean cost. Comm. Anal. Geom. 19, no. 1, 191-208, 2011.


[^0]:    *Partially supported by National Science Foundation Grant DMS-0654267.
    ${ }^{\dagger}$ Supported by Spain Government projects MTM2008-06349-C03-01, MTM2011-27739-C04-01, GenCat 2009SGR345 and NSF grant DMS-0635607.
    $\ddagger$ Supported by NSF grant DMS-0901449.

