

# Cap 4. Ondículas ortonormales (Wavelets)

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## 4.1. INTRODUCCIÓN Y DEFINICIONES

En el ejercicio 1.5.4 se ha probado que  $\{f_{n,k} : n, k \in \mathbb{Z}\}$  es base o.n. de  $L^2(\mathbb{R})$ , donde  $f_{n,k}(x) = e^{2\pi i n x} \chi_{[0,1]}(x-k)$ .

Def 4.1.1.  $f: \mathbb{R} \rightarrow \mathbb{C}$ . El operador traslación por  $y_0 \in \mathbb{R}$  es  $(T_{y_0} f)(x) = f(x - y_0)$ . El operador modulación por  $\omega_0 \in \mathbb{R}$  es  $(M_{\omega_0} f)(x) = e^{2\pi i \omega_0 x} f(x)$ .

Ejercicio 4.1.1 Prueba que las operaciones  $T_{y_0}$  y  $M_{\omega_0}$  son lineales y unitarias en  $L^2(\mathbb{R})$  y si  $f \in L^1(\mathbb{R})$  se tiene

$$F(T_{y_0} f)(\omega) = M_{-y_0} Ff(\omega)$$

$$F(M_{\omega_0} f)(\omega) = T_{\omega_0} Ff(\omega)$$

donde  $F$  denota la transformada de Fourier.

$$\begin{aligned} \text{S/ } \|T_{y_0} f\|_2^2 &= \int_{-\infty}^{\infty} |T_{y_0} f(x)|^2 dx = \int_{-\infty}^{\infty} |f(x - y_0)|^2 dx \stackrel{x - y_0 = y}{=} \\ &= \int_{-\infty}^{\infty} |f(y)|^2 dy = \|f\|_2^2. \end{aligned}$$

$$\begin{aligned} \|M_{\omega_0} f\|_2^2 &= \int_{-\infty}^{\infty} |M_{\omega_0} f(x)|^2 dx = \int_{-\infty}^{\infty} |e^{2\pi i \omega_0 x} f(x)|^2 dx \\ &= \int_{-\infty}^{\infty} |f(x)|^2 dx = \|f\|_2^2. \end{aligned}$$

$$F(T_{y_0} f)(\omega) = \int_{-\infty}^{\infty} T_{y_0} f(x) e^{-2\pi i \omega x} dx =$$

(2)

$$\begin{aligned}
 &= \int_{-\infty}^{\infty} f(x-y_0) e^{-2\pi i \omega x} dx \stackrel{x-y_0=y}{=} \int_{-\infty}^{\infty} f(y) e^{-2\pi i \omega (y+y_0)} dy \\
 &= e^{-2\pi i \omega y_0} \underbrace{\int_{-\infty}^{\infty} f(y) e^{-2\pi i \omega y} dy}_{Ff(\omega)} = M_{-y_0} Ff(\omega)
 \end{aligned}$$

$$\begin{aligned}
 \bullet F(M_{\omega_0} f)(\omega) &= \int_{-\infty}^{\infty} (M_{\omega_0} f)(x) e^{-2\pi i \omega x} dx = \\
 &= \int_{-\infty}^{\infty} e^{2\pi i \omega_0 x} f(x) e^{-2\pi i \omega x} dx = \int_{-\infty}^{\infty} f(x) e^{-2\pi i (\omega - \omega_0) x} dx \\
 &= Ff(\omega - \omega_0) = T_{\omega_0} Ff(\omega). \quad \blacksquare
 \end{aligned}$$

Def. 4.1.2.  $f: \mathbb{R} \rightarrow \mathbb{C}$ . El operador dilatación por  $a > 0$  es  $D_a f(x) = a^{1/2} f(ax)$ .

Ejercicio 4.1.2. Probar que  $D_a$  es un operador lineal y unitario en  $L^2(\mathbb{R})$  y si  $f \in L^1(\mathbb{R})$

$$F(D_a f)(\omega) = D_{1/a} Ff(\omega)$$

$$\begin{aligned}
 \text{si} \\
 \|D_a f\|_2^2 &= \int_{-\infty}^{\infty} |D_a f(x)|^2 dx = \int_{-\infty}^{\infty} |a^{1/2} f(ax)|^2 dx \\
 &= \int_{-\infty}^{\infty} a |f(ax)|^2 dx \stackrel{ax=y}{=} \int_{-\infty}^{\infty} a |f(y)|^2 \frac{dy}{a} = \int_{-\infty}^{\infty} |f(y)|^2 dy = \|f\|_2^2
 \end{aligned}$$

$$\begin{aligned}
 \bullet F(D_a f)(\omega) &= \int_{-\infty}^{\infty} D_a f(x) e^{-2\pi i \omega x} dx = \int_{-\infty}^{\infty} a^{1/2} f(ax) e^{-2\pi i \omega x} dx \\
 &\stackrel{ax=y}{=} a^{1/2} \int_{-\infty}^{\infty} f(y) e^{-2\pi i \omega \frac{y}{a}} \frac{dy}{a} = \frac{1}{a^{1/2}} \int_{-\infty}^{\infty} f(y) e^{-2\pi i \frac{\omega}{a} y} dy \\
 &= a^{-1/2} Ff(a^{-1} \omega) = D_{a^{-1}} Ff(\omega) = D_{1/a} Ff(\omega). \quad \blacksquare
 \end{aligned}$$

$\psi = \text{psi}$

③

Def 4.1.3 Una función  $\psi \in L^2(\mathbb{R})$  es una ondícula ortogonal (orthogonal wavelet) para  $L^2(\mathbb{R})$  si el conjunto

$$\{\psi_{j,k}(x) = D_{2^j} T_k \psi(x) = 2^{j/2} T_k \psi(2^j x) = 2^{j/2} \psi(2^j x - k) : j, k \in \mathbb{Z}\}$$

es base o.n. de  $L^2(\mathbb{R})$ .

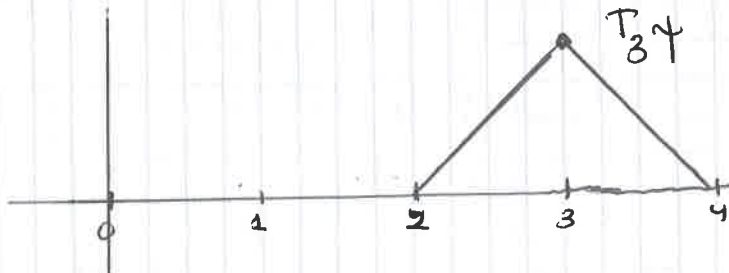
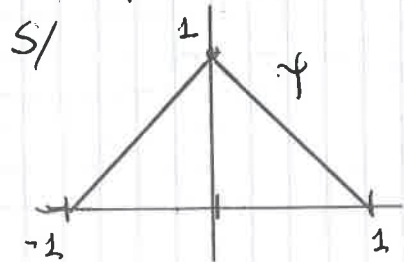
Ejercicio 4.1.3. Comprueban que  $D_{2^j} T_k \psi \neq T_k D_{2^j} \psi$

s/  $D_{2^j} T_k \psi(x) = 2^{j/2} \psi(2^j x - k)$ , pero

$$T_k D_{2^j} \psi(x) = D_{2^j} \psi(x - k) = 2^{j/2} \psi(2^j(x - k)) = 2^{j/2} \psi(2^j x - 2^j k)$$

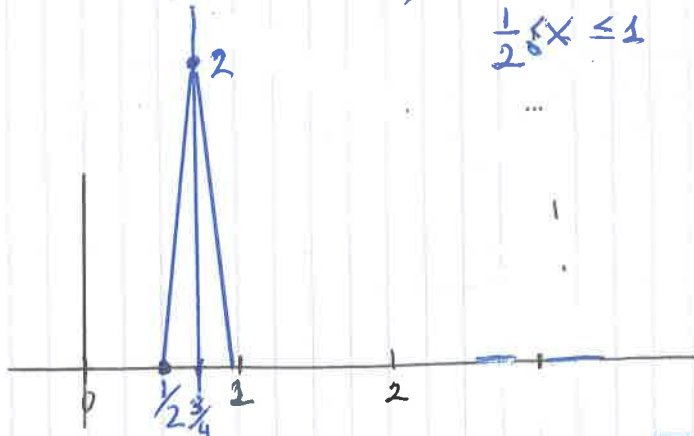
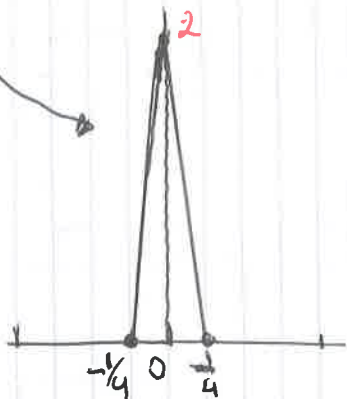
Ejercicio 4.1.4 Sea  $\psi(x) = \begin{cases} x+1 & \text{si } -1 \leq x \leq 0 \\ 1-x & \text{si } 0 < x \leq 1 \\ 0 & \text{resto} \end{cases}$

Dibujar  $\psi$ ,  $T_3 \psi$ ,  $D_4 \psi$ ,  $D_4 T_3 \psi$  y  $T_3 D_4 \psi$



$$D_4 \psi(x) = 2\psi(4x) \\ -1 \leq 4x \leq 1 \Leftrightarrow -\frac{1}{4} \leq x \leq \frac{1}{4}$$

$$D_4 T_3 \psi(x) = 2\psi(4x-3) \\ -1 \leq 4x-3 \leq 1 ; 2 \leq 4x \leq 4 \\ \frac{1}{2} \leq x \leq 1$$

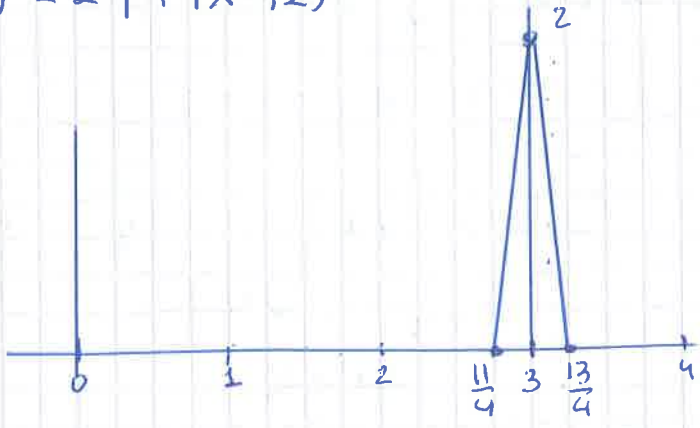


$$T_3 D_4 \psi(x) = 2 \psi(4(x-3)) = 2 \psi(4x-12)$$

$$-1 \leq 4x-12 \leq 1$$

$$11 \leq 4x \leq 13$$

$$\frac{11}{4} \leq x \leq \frac{13}{4}$$



Ejercicio 4.1.5 Probar

(a)  $\|\psi_{j,k}\|_2 = \|\psi\|_2$

(b)  $\mathcal{F}(\psi_{j,k})(\omega) = 2^{-j/2} e^{-2\pi i k \frac{\omega}{2^j}} \mathcal{F}\psi(\frac{\omega}{2^j})$

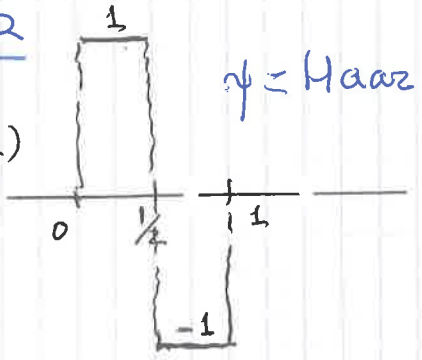
S/ (a)  $\|\psi_{j,k}\|_2 = \|D_{2^j} T_k \psi\|_2 \stackrel{\text{Ej 4.1.2}}{=} \|T_k \psi\|_2 \stackrel{\text{Ej 4.1.1}}{=} \|\psi\|_2$

(b)  $\mathcal{F}(\psi_{j,k})(\omega) = \mathcal{F}(D_{2^j} T_k \psi)(\omega) = D_{2^j} \mathcal{F}(T_k \psi)(\omega) =$   
 $= D_{2^j} M_{-k} \mathcal{F}\psi(\omega) = 2^{-j/2} (M_{-k} \mathcal{F}\psi)(2^{-j} \omega) =$   
 $= 2^{-j/2} e^{-2\pi i k \frac{\omega}{2^j}} \mathcal{F}\psi(\frac{\omega}{2^j})$

## 4.2. LAS ONDÍCULAS DE HAAR Y DE SHANNON

### 4.2.1. LA ONDÍCULA DE HAAR

$$\psi(x) = \begin{cases} 1 & \text{si } 0 \leq x \leq 1/2 \\ -1 & \text{si } 1/2 < x \leq 1 \\ 0 & \text{resto} \end{cases} \quad (4.2.1)$$



$$\|\psi\|_2^2 = \int_0^{1/2} 1^2 dx + \int_{1/2}^1 (-1)^2 dx = \int_0^1 1 dx = 1$$

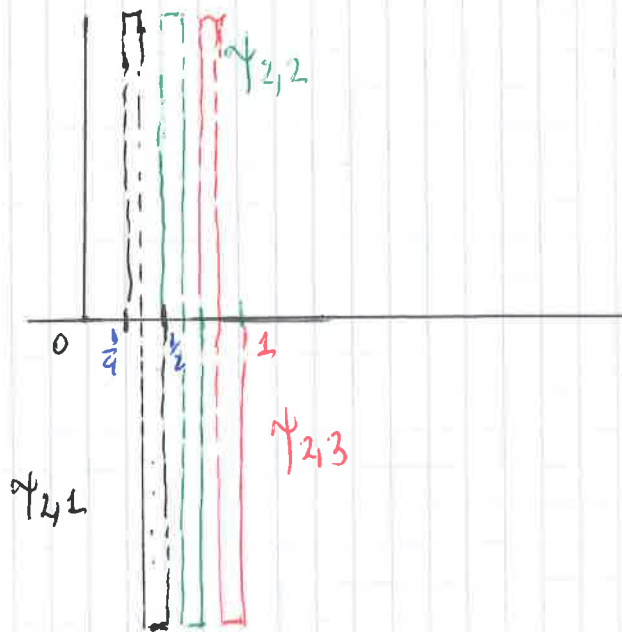
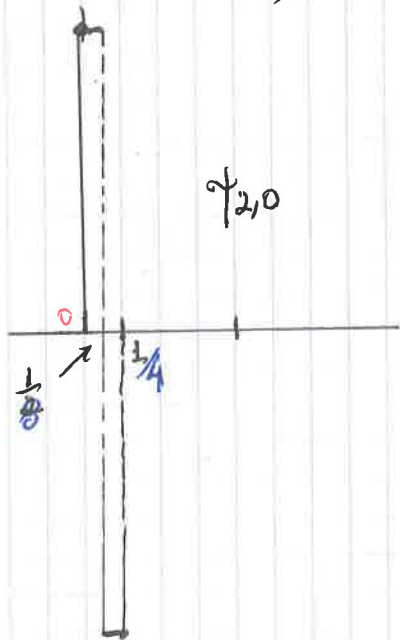
Por el ejercicio 4.1.5,  $\|\psi_{j,k}\|_2 = 1 \quad \forall j, k \in \mathbb{Z}$ .

Observar :  $\int_{-\infty}^{\infty} \psi(x) dx = \int_0^{\frac{1}{2}} 1 dx + \int_{\frac{1}{2}}^1 (-1) dx = \frac{1}{2} - \frac{1}{2} = 0$  (5)

Ejercicio 4.2.1. Para la ondícula de Haar, dibujar

$\psi_{2,0}$  ;  $\psi_{2,1}$  ;  $\psi_{2,2}$  y  $\psi_{2,3}$

S/  $\psi_{2,0}(x) = D_{2^2} T_0 \psi(x) = 2\psi(4x)$  ;  $\psi_{2,1}(x) = D_{2^2} T_1(x) = 2\psi(4x-1)$   
 $0 \leq 4x \leq 1$  ;  $0 \leq x \leq \frac{1}{4}$        $0 \leq 4x-1 \leq 1$  ;  $\frac{1}{4} \leq x \leq \frac{1}{2}$



Proposición 4.2.1. Con  $\psi$  dada por (4.2.1), el conjunto  $\{\psi_{j,k} : j,k \in \mathbb{Z}\}$  es un sistema o.n. en  $L^2(\mathbb{R})$ .